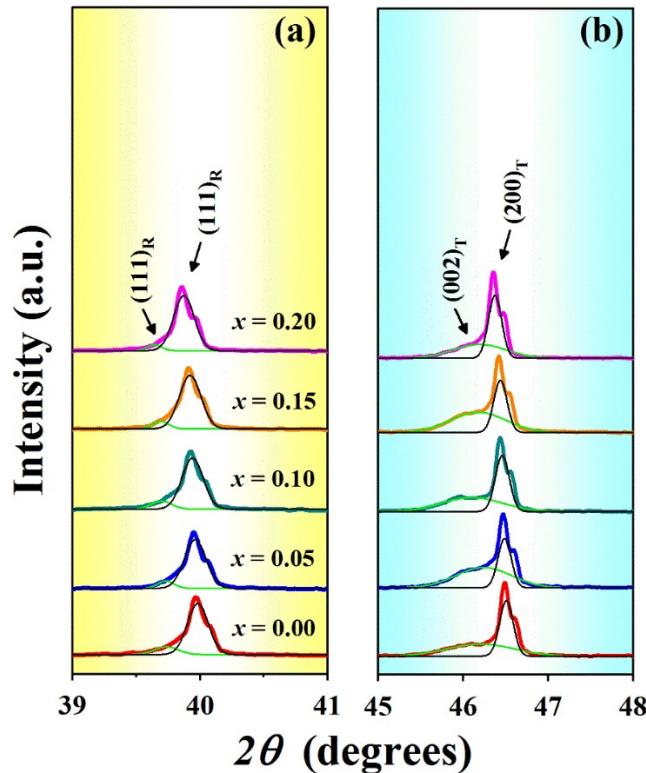


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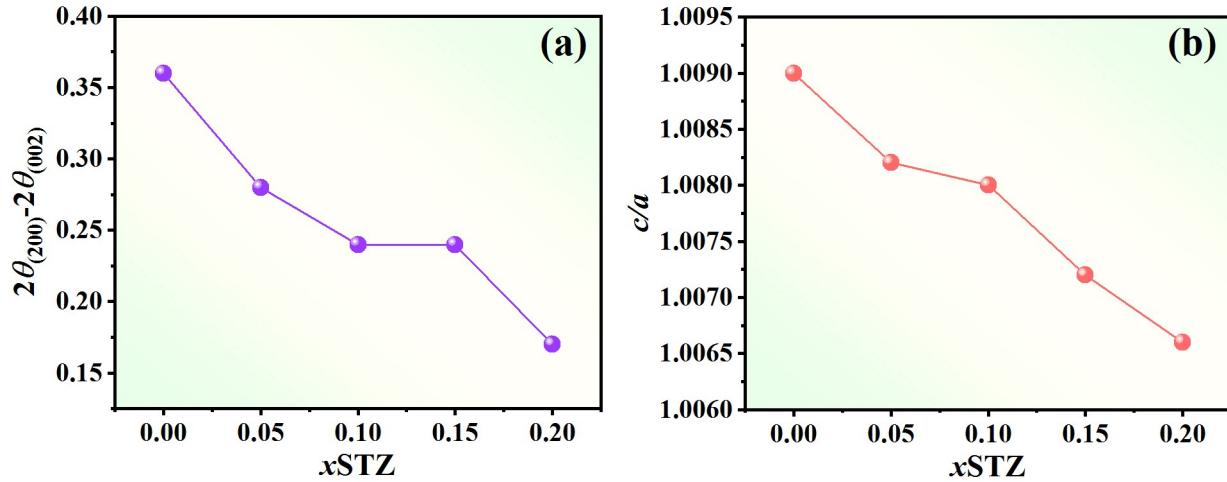
Supplementary information



2

3 **Fig. S1.** XRD patterns of the ceramics at (a)  $2\theta = 39 - 41^\circ$  and (b)  $2\theta = 45 - 48^\circ$ .

4 X-ray diffraction (XRD) patterns of the ceramics, recorded over the range  $2\theta = 20 - 70^\circ$ ,  
5 are presented in the main text. The observed diffraction peaks correspond well to the structure  
6 reported in JCPDS file No. 36-0340. A specific JCPDS file for SBT is not available because it is  
7 a solid solution rather than a single-phase compound; however, its crystal structure adopts a  
8 perovskite-type arrangement similar to its parent compound  $\text{SrTiO}_3$  (JCPDS file No. 35-0734).  
9 Similarly, no JCPDS file exists for STZ. XRD patterns of the ceramics at (a)  $2\theta = 39 - 41^\circ$  and (b)  
10  $2\theta = 45 - 48^\circ$  are shown in Fig. S1. All samples exhibited peak splitting near  $2\theta = 39 - 41^\circ$   
11 ( $(111)_R/(1\bar{1}\bar{1})_R$ ) and  $45 - 48^\circ$  ( $(200)_T/(002)_T$ ), as shown in Fig. S1(a) and (b), respectively. This  
12 indicates the coexistence of rhombohedral (R) and tetragonal (T) phases.



13

14

**Fig. S2.** (a)  $\Delta 2\theta_{200} = 2\theta_{(200)} - 2\theta_{(002)}$ , and (b) tetragonality ( $c/a$ ) versus of  $x$

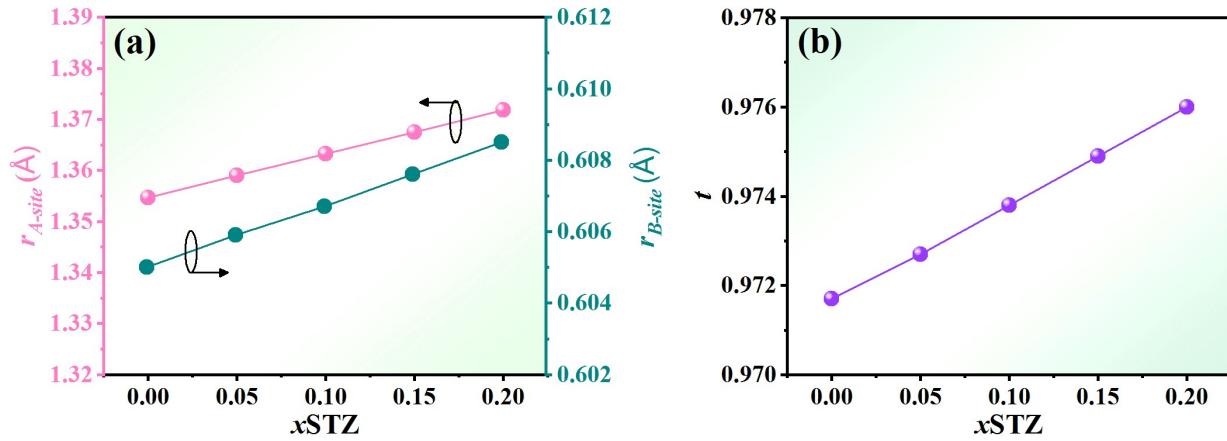
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16

The values of  $\Delta 2\theta_{200} = 2\theta_{(200)} - 2\theta_{(002)}$  and tetragonality ( $c/a$ ) as a function of  $x$  are presented

17 in Fig. S2. Both parameters decrease with increasing  $x$ , indicating a progressive shift toward a more cubic structure and reduced tetragonality.

18



19

20

**Fig. S3.** (a)  $r_{A\text{-site}}$  and  $r_{B\text{-site}}$  and (b)  $t$  factors as a function of  $x$

21

22

In this work, the average ionic radius of the A-and B-site ions ( $r_{A\text{-site}}$  and  $r_{B\text{-site}}$ ) of the

ceramics was calculated from the equation:

23

24 
$$r_{A-site} = x[r_{Sr^{2+}}] + (1-x)[0.35r_{Bi^{3+}} + 0.35r_{Na^+} + 0.21r_{Sr^{2+}} + 0.6r_{Bi^{3+}}] \quad (S1)$$

25

26 and

27

28 
$$r_{B-site} = x[0.85r_{Ti^{4+}} + 0.15r_{Zr^{4+}}] + (1-x)[0.7r_{Ti^{4+}} + 0.3r_{Ti^{4+}}] \quad (S2)$$

29

30 , where  $r_{Sr^{2+}}$ ,  $r_{Bi^{3+}}$ ,  $r_{Na^+}$ ,  $r_{Ti^{4+}}$ , and  $r_{Zr^{4+}}$  are ionic radii of Sr, Bi, Ti, and Zr, respectively. Fig. 31 S3 presents  $r_{A-site}$  and  $r_{B-site}$  values as a function of  $x$ . The  $r_{A-site}$  and  $r_{B-site}$  values increased with 32 increasing  $x$ . To determine the lattice distortion, a simple description of the geometric packing 33 within the perovskite structure can be characterized by the tolerance factor ( $t$ ), which is defined by 34 the following equation [S1,S2]:

35

36 
$$t = \frac{(r_{A-site} + r_O)}{\sqrt{2}(r_{B-site} + r_O)} \quad (S3)$$

37

38 , where  $r_A$ ,  $r_B$  and  $r_O$  are the ionic radii of A-site cations, B-site cations, and the oxygen anion, 39 respectively. In this work, Shannon ionic radii were used:  $Bi^{3+}$  (C.N. = 12) = 1.38 Å,  $Na^+$  (C.N. = 40 12) = 1.39 Å,  $Sr^{2+}$  (C.N. = 12) = 1.44 Å,  $Ti^{4+}$  (C.N. = 6) = 0.605 Å,  $Zr^{4+}$  (C.N. = 6) = 0.72 Å, and 41  $O^{2-}$  (C.N. = 6) = 1.40 Å. When  $t = 1$ , a perfect cubic cell is formed, and when  $t$  deviates from 1, 42 the perovskite cell gets deformed [S2]. The tolerance factor ( $t$ ) was calculated for each composition 43 and found to increase progressively with increasing  $x$ , ranging from 0.9717 to 0.9782. This range 44 lies well within the typical stability window for perovskite structures ( $t = 0.880 - 1.090$ ) [S3],

45 suggesting that all compositions maintain a structurally stable perovskite phase. This trend is  
 46 illustrated in Fig. S3(b). It should be noted that the c/a ratio and the tolerance factor exhibit  
 47 opposite trends. The tolerance factor increases with x because Sr and Zr enlarge the effective A-  
 48 and B-site ionic radii, shifting the structure toward a more cubic configuration. In contrast,  
 49 tetragonality decreases since the incorporation of  $\text{Sr}^{2+}$  and  $\text{Zr}^{4+}$  suppresses the off-centering of Bi  
 50 and Ti, thereby reducing lattice distortion symmetry. Meanwhile, tetragonality decreases since  $\text{Sr}^{2+}$   
 51 and  $\text{Zr}^{4+}$  reduce Bi- and Ti off-centering.

52

53 **Table S1.** Phase formation refinement results of the studied ceramics

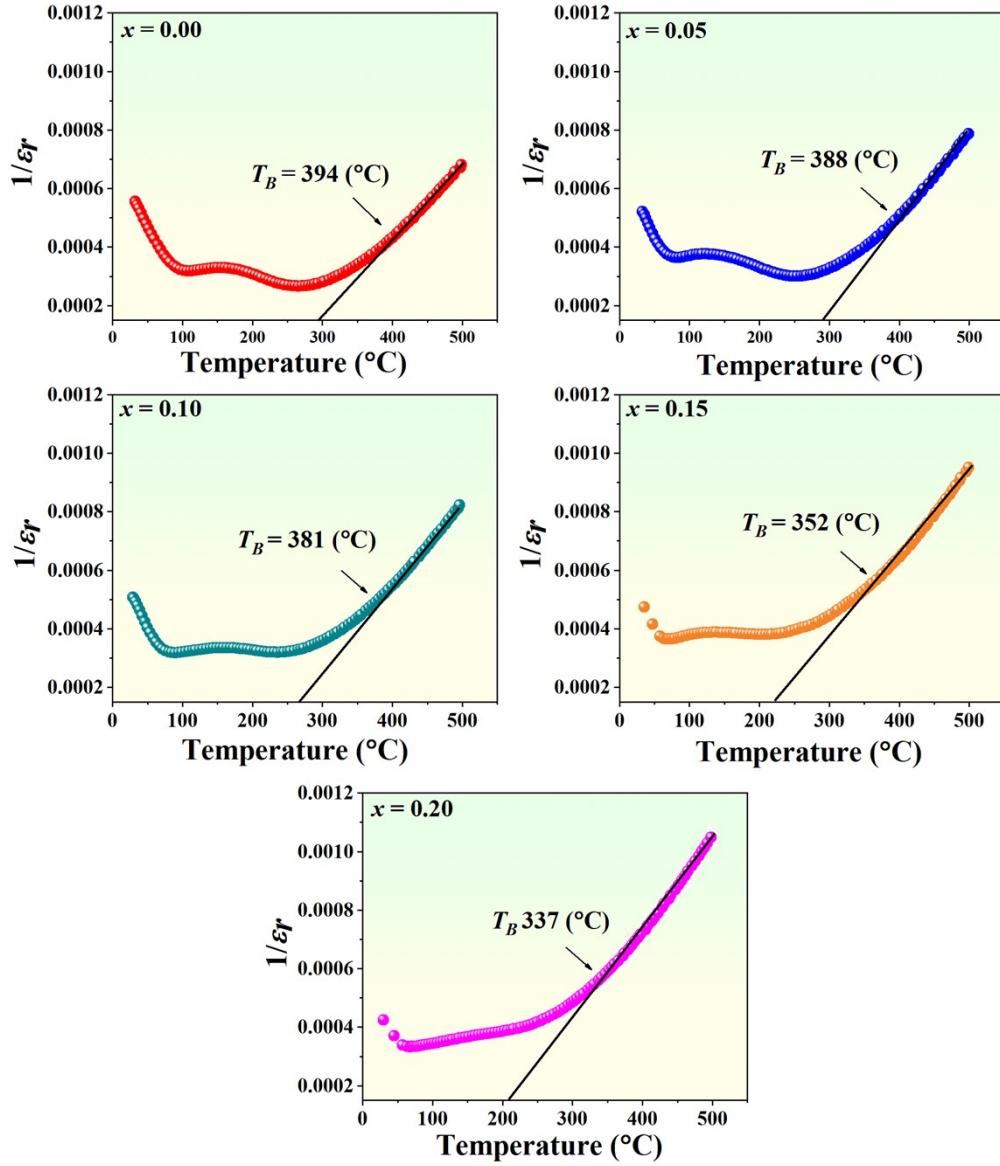
<i>x</i>	0.00		0.05		0.10		0.15		0.20	
<b>Space group</b>	<i>R3c</i>	<i>P4bm</i>								
<b>Phase frac. (%)</b>	41	59	58	42	64	36	66	34	78	22
$x^2$		1.91		1.41		1.99		1.37		1.98
$R_w$ (%)		11.62		9.80		13.18		9.60		11.68

54

55 Table S1 provides the Rietveld refinement data obtained using the GSAS-II software. The  
 56 corresponding results, illustrated in Fig. 2(b-f), confirm the phase evolution across the samples.  
 57 All compositions exhibited a coexistence of rhombohedral (*R3c*) and tetragonal (*P4bm*) phases.  
 58 As the STZ content increased, the fraction of the rhombohedral *R3c* phase gradually rose, while  
 59 the tetragonal *P4bm* phase diminished. This compositional shift is consistent with the observed  
 60 decrease in the *c/a* ratio.

61

62

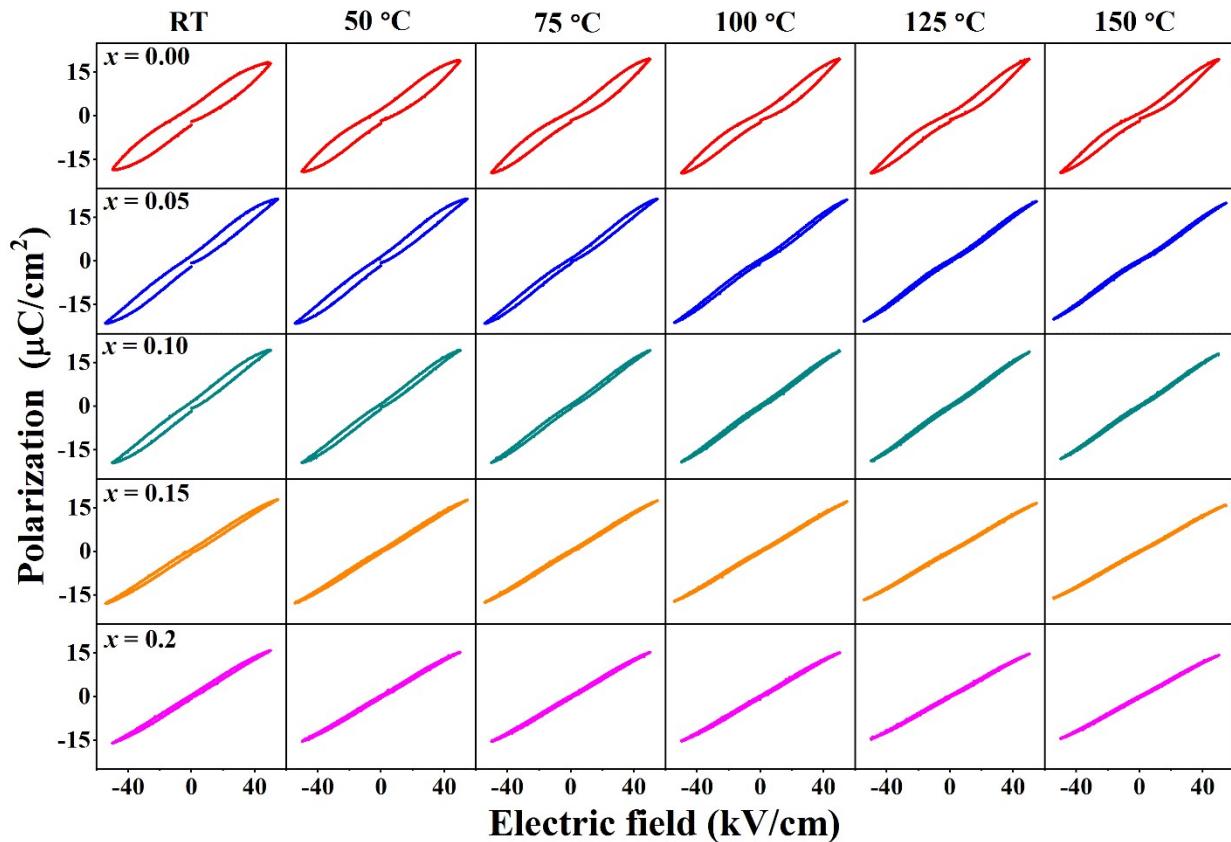


**Fig. S4.** Plots of  $1/\epsilon_r$  versus temperature of the samples.

66 The Burns temperature ( $T_B$ ) denotes the point at which the dielectric constant begins to  
 67 deviate from the Curie-Weiss law. In this work,  $T_B$  was determined from the  $1/\epsilon_r$  versus  $T$  plot, as  
 68 shown in Fig. S4. The  $T_B$  values for BNT-SBT- $x$ STZ ceramics were observed to decrease with

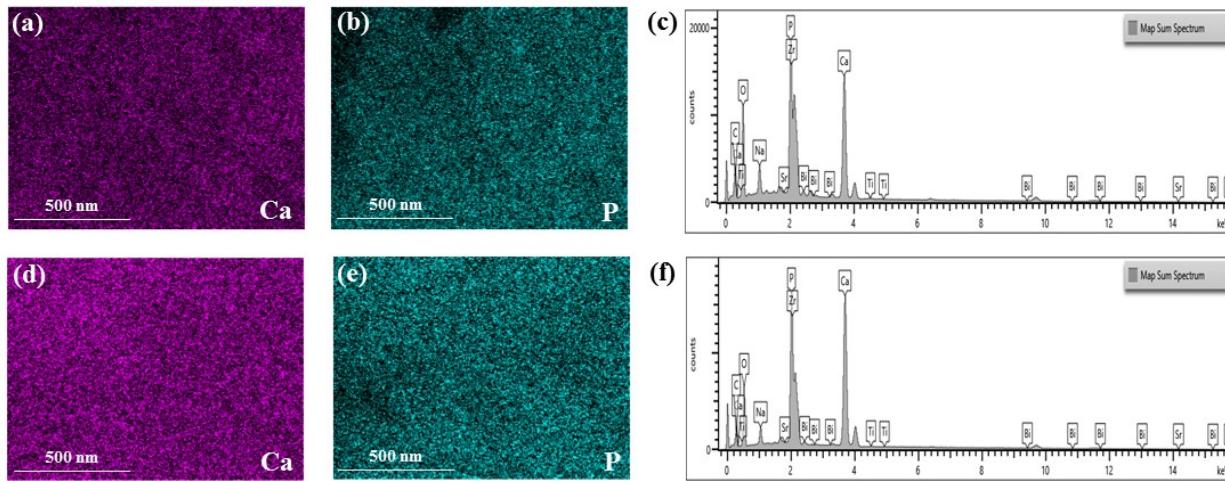
69 increasing STZ content, specifically: 394 °C, 388 °C, 381 °C, 352 °C, 337 °C, and 326 °C for  $x =$   
70 0.00, 0.05, 0.10, 0.15, and 0.20, respectively.

71



73 **Fig. S5.** Plots of polarization versus electric filed of the samples at various temperatures

74 Effects of temperature on the ferroelectric properties are presented in Fig. S5. Samples  
75 with compositions ranging from  $x = 0.00$  to  $x = 0.15$  exhibited slim and pinched  $P$ - $E$  hysteresis  
76 loops, which are characteristic of ergodic relaxor ferroelectric behavior. The temperature-  
77 dependent  $P$ - $E$  loops were measured from RT to 150°C under 50 kV/cm. It can be seen that both  
78  $P_{max}$  and  $P_r$  slightly increase with increasing temperature (Fig. 9). Furthermore, the  $x = 0.20$   
79 samples exhibited a very slim and barely visible pinched loop, indicating a further suppression of  
80 long-range ferroelectric order and a stronger tendency toward a relaxor state.



81

82 **Fig. S6.** EDS elemental mapping corresponding to Fig. 12(c) and (f), showing the apatite  
 83 layer formed on the samples after immersion in SBF for 30 days: (a-c) uncoated sample and (d-f)  
 84  $\beta$ -TCP–coated sample, respectively.

85 The EDS analysis results of the apatite layer formed after immersion in SBF for 30 days  
 86 are presented in Fig. S6, consistent with the surface morphology shown in Fig. 12(c and f). The  
 87 EDS data confirmed the formation of an apatite layer, as indicated by the presence of calcium (Ca)  
 88 and phosphorus (P), a key component of apatite, despite the fact that the original composition of  
 89 the studied samples (0.85[BNT–SBT]-0.15STZ) does not contain Ca or P. Furthermore, the Ca/P  
 90 ratio increased from 1.76 in uncoated samples to 2.05 in  $\beta$ -TCP-coated samples. Based on these  
 91 EDS results, it can be concluded that bone-like apatite layers developed on the sample surfaces  
 92 after immersion in SBF.

93

94

95

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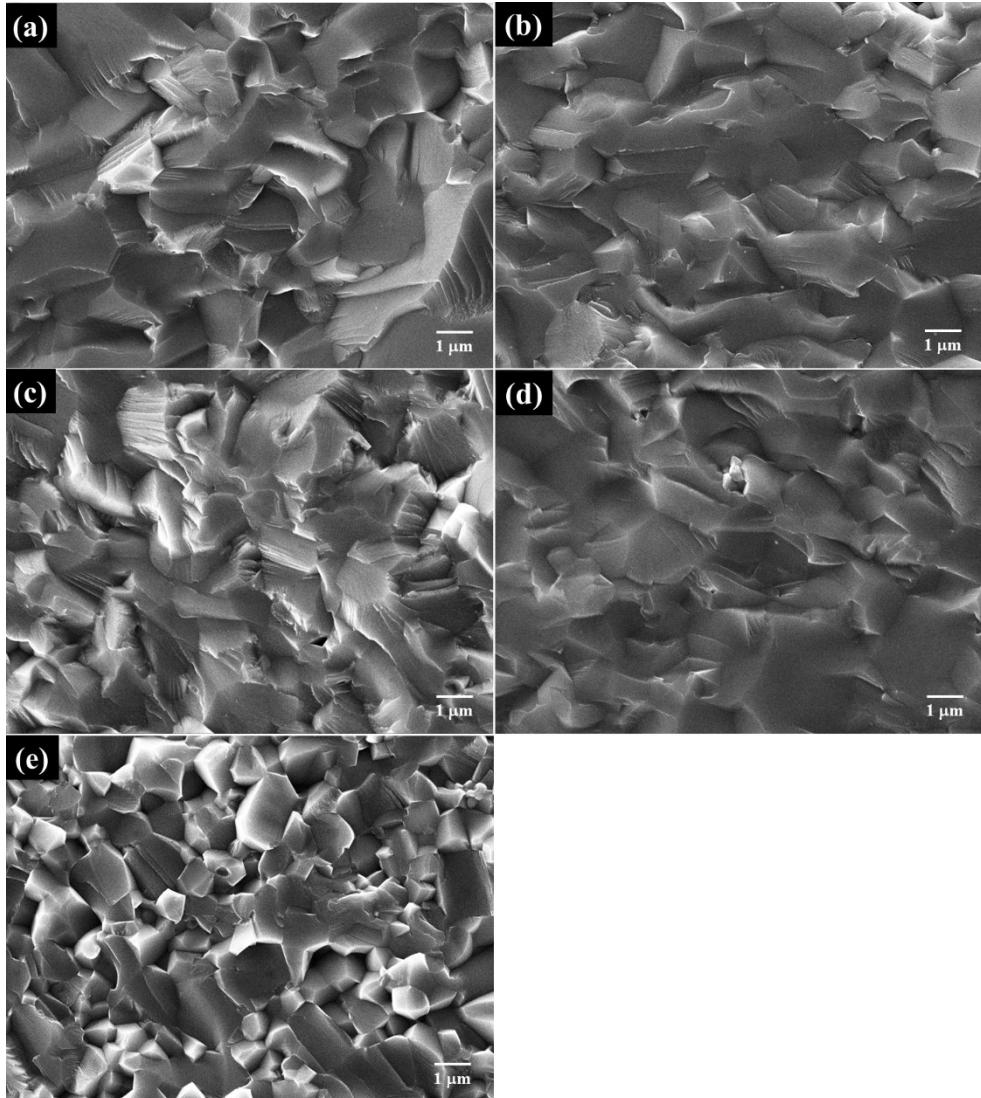
97 **Table S2.** Peak positions and full width at half maximum of Raman spectra of  $(1-x)$ [BNT-SBT]-  
98  $x$ STZ ceramics.

Sample	The mode of Raman	A-site vibration	Ti-O Band		TiO <sub>6</sub> Octahedra		A1+E	
		P1	P2	P3	P4	P5	P6	P7
$x = 0.00$	Intensity (a.u.)	17071	17840	12667	10328	8570	1348	3005
	Peak position (cm <sup>-1</sup> )	104	246	315	527	604	754	824
$x = 0.05$	FWHM (cm <sup>-1</sup> )	212	139	89	101	62	62	135
	Intensity (a.u.)	18645	17823	11642	10493	7602	1749	2521
$x = 0.05$	Peak position cm <sup>-1</sup> )	103	250	314	527	603	755	830
	FWHM (cm <sup>-1</sup> )	217	139	89	112	65	82	111
$x = 0.10$	Intensity (a.u.)	17327	18426	13118	10493	8800	1639	2848
	Peak position (cm <sup>-1</sup> )	106	250	314	526	603	755	830
$x = 0.10$	FWHM (cm <sup>-1</sup> )	204	134	94	99	69	76	112
	Intensity (a.u.)	19456	18694	12632	11048	8146	1808	2679
$x = 0.15$	Peak position (cm <sup>-1</sup> )	106	251	312	526	601	755	829
	FWHM (cm <sup>-1</sup> )	203	134	95	112	72	74	112
$x = 0.20$	Intensity (a.u.)	21025	19274	12653	11305	7623	1818	2582
	Peak position (cm <sup>-1</sup> )	105	252	311	525	599	756	831
	FWHM (cm <sup>-1</sup> )	202	131	98	114	74	78	108

100 The Raman spectra were deconvoluted using Lorentzian functions and the resulting  
101 modes were labeled P1 – P7. Table S2 summarizes the corresponding peak positions and full  
102 widths at half maximum (FWHM). The mode at  $\leq 200$  cm<sup>-1</sup> (P1) is assigned to A-site cation  
103 vibrations of the perovskite lattice (Bi, Na and Sr). The modes in the 200 – 400 cm<sup>-1</sup> region (P2  
104 and P3) arise from Ti-O/(Ti,Zr)-O vibrations. As seen in Fig. 3, the P2 peak position shows  
105 almost no dependence on  $x$ , remaining in the range 246 – 252 cm<sup>-1</sup> for all compositions, whereas  
106 the relative intensity of P3 gradually decreases with increasing STZ content. This reduction in P3  
107 intensity is consistent with a decrease in the tetragonal *P4bm* fraction, in agreement with the  
108 XRD results [S4]. The FWHM of P2 exhibited a slight narrowing, whereas the P3 band showed  
109 an overall broadening as  $x$  increased (Table S2), indicating that the tetragonal-like local TiO<sub>3</sub>  
110 octahedral environments became increasingly disordered, even as their fraction decreased within  
111 the phase-coexistence region. In the 430 – 700 cm<sup>-1</sup> range (P4 and P5), which corresponded to  
112 TiO<sub>6</sub> stretching vibrations in rhombohedral-like and tetragonal-like local structures, respectively,  
113 P5 displayed a clear increase in FWHM with increasing STZ content, whereas P4 showed only  
114 modest, non-monotonic changes. The progressive broadening of the P3 and P5 modes therefore  
115 reflects a high level of local structural disorder that strongly affects the dynamics of the  
116 TiO<sub>3</sub>/(Ti,Zr)O<sub>3</sub> octahedra. The high-frequency modes above 700 cm<sup>-1</sup> (P6 and P7) are attributed  
117 to overlapping A1(LO) and E(LO) bands [S5].

118

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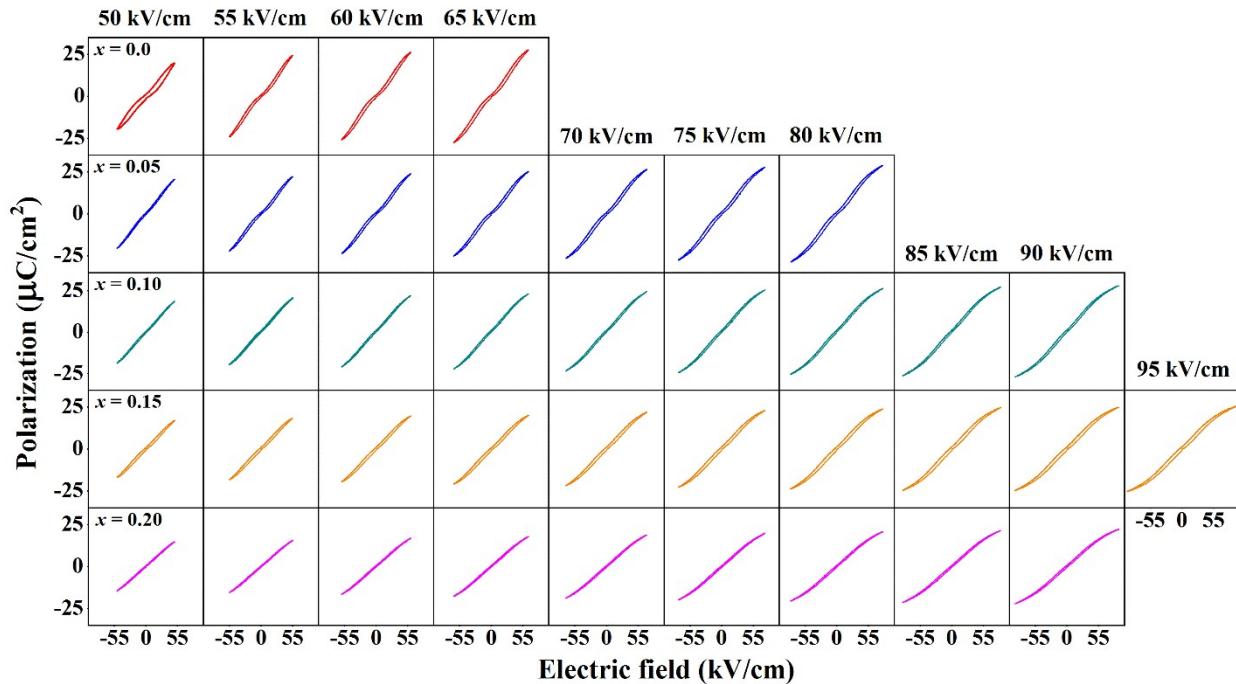


120

121 **Fig. S7.** SEM images of the cross-sections of the as-prepared  $(1-x)[\text{BNT-SBT}]-x\text{STZ}$

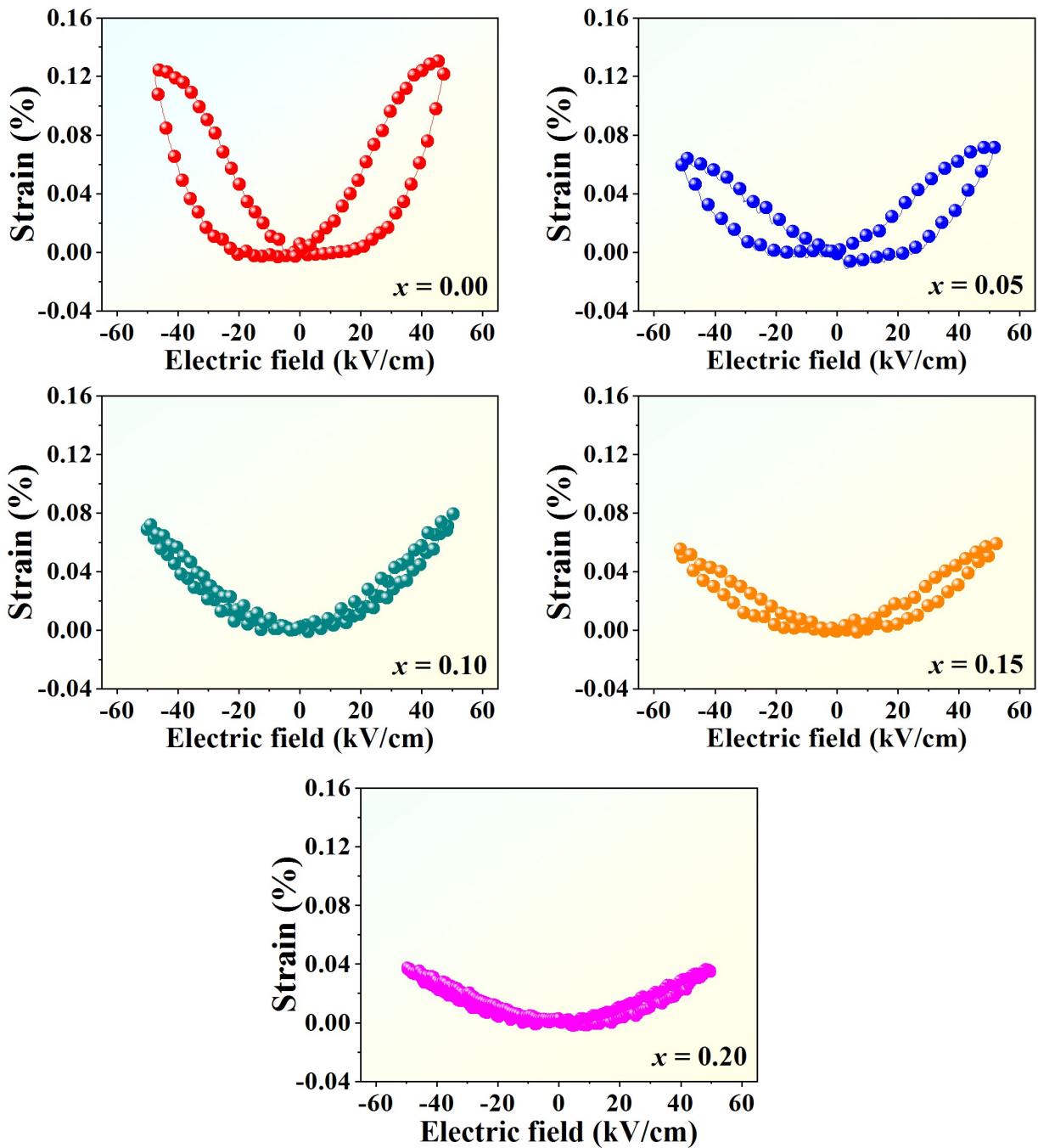
122 ceramics: (a)  $x = 0.00$ , (b)  $x = 0.05$ , (c)  $x = 0.10$ , (d)  $x = 0.15$  and (e)  $x = 0.20$

123 Fig. S7 shows that increasing the STZ content progressively refines the cross-sectional  
124 microstructure of the studied ceramics. The undoped sample exhibits course, irregular grains,  
125 whereas higher- $x$  compositions ( $x \geq 0.10$ ) develop smaller, more uniform, and more densely  
126 packed grains. This indicates that STZ addition effectively suppresses abnormal grain growth and  
127 promotes a fine, homogeneous microstructure.



130 **Fig. S8.** Polarization-electric field ( $P$ - $E$ ) loops of the samples measured at 125 °C under  
 131 different applied electric fields.

132 Fig. S8 illustrates how the  $P$ - $E$  loops of the samples ( $x = 0.00 - 0.20$ ) evolve with increasing  
 133 electric field at 125 °C. As the field rises from 50 to 95 kV/cm, all compositions exhibit higher  
 134 polarization, reflecting a stronger field-induced polarization response. Moreover, increasing the  
 135 STZ content (higher  $x$ ) results in progressively slimmer and more linear loops, indicating an  
 136 enhanced ergodic relaxor character relative to the lower  $x$  compositions.

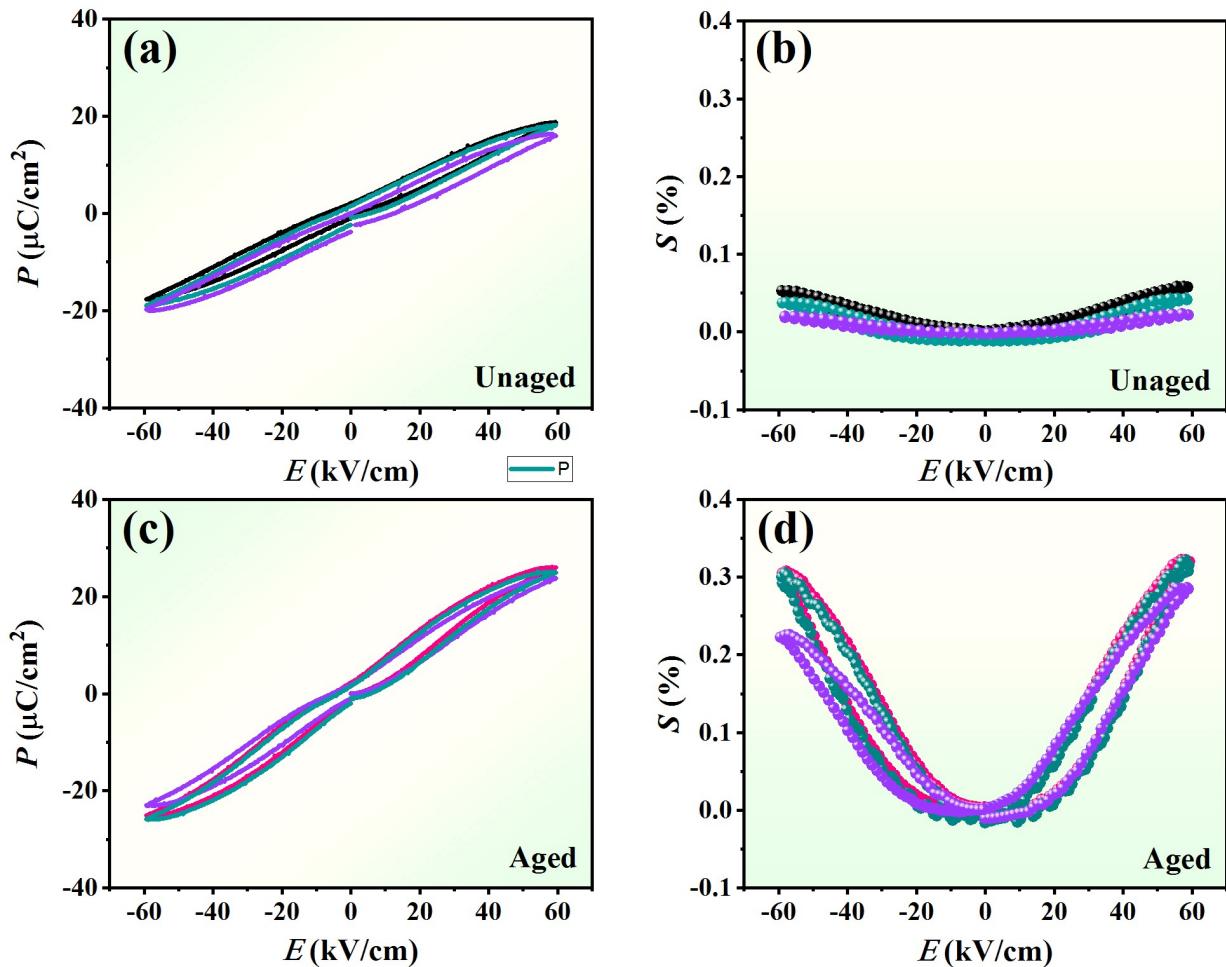


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139 **Fig. S9.** Bipolar strain-electric field (*S-E*) loops of  $(1-x)[\text{BNT-SBT}]-x\text{STZ}$  ceramics, measured at  
140 RT under an electric field of 50 kV/cm and a frequency of 1 Hz.

141 Fig. S9 presents the *S-E* loops of the  $(1-x)[\text{BNT-SBT}]-x\text{STZ}$  ceramics measured under an  
142 electric field of 50 kV/cm at 1 Hz for compositions  $x = 0.00, 0.05, 0.10, 0.15$ , and  $0.20$ . As

143  $x$  increases, the maximum strain decreases, and the loops become progressively slimmer,  
 144 indicating reduced ferroelectric ordering and enhanced relaxor behavior. For  $x \geq 0.15$ , the loops  
 145 approach a nearly linear shape, consistent with a dominant ergodic relaxor phase and minimal  
 146 remanent strain.



147

148 **Fig. S10.** (a-b)  $P$ - $E$  hysteresis loops and (c-d)  $S$ - $E$  loops for aged samples ( $x = 0.15$ ),  
 149 measured under an electric field of 60 kV/cm at 1 Hz from three repeated specimens ( $n = 3$ ).

150 To ensure statistically reliable strain measurements,  $P$ - $E$  and  $S$ - $E$  loops were recorded  
 151 three times. Fig. S10(a-b) shows the  $P$ - $E$  hysteresis loops and bipolar  $S$ - $E$  loops for unaged  
 152 samples ( $x = 0.15$ ), measured at RT under an electric field of 60 kV/cm at 1 Hz from three

153 repeated specimens (n = 3). Before annealing, the strain values were 0.03%, 0.05%, and 0.06%,  
154 with an average of 0.05% (Fig. S10(b)). After annealing, the strain values increased to 0.29%,  
155 0.32%, and 0.33%, resulting in an average of 0.31% (Fig. S10(d)).

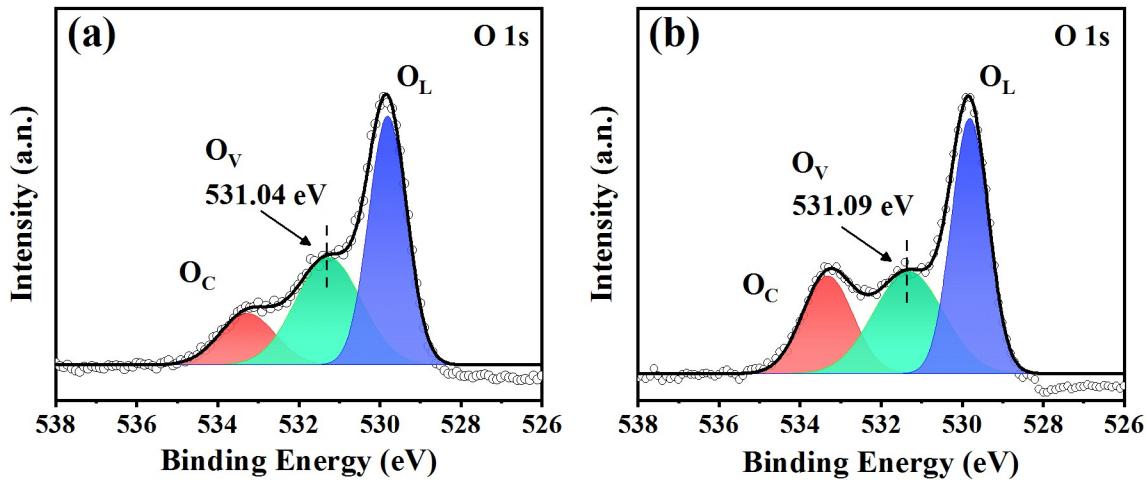
156 **Table S3.** Comparison of cell viability assay at 24 h.

Sample	Cell viability (%)	Concentration ( $\mu\text{g}/\text{ml}$ )	Results	Ref.
BNT-SBT-0.15STZ	>95	0.1 – 1000	Non-toxic	<i>This work</i>
10Co/HA/T coating	67	-	Non-toxic	84
HAp/AlN nanocomposites	>85	1000	Non-toxic	85
BG 1%FeS <sub>2</sub> and 3%FeS <sub>2</sub>	>90	1000	Non-toxic	86
Glass-ceramics (50SiO <sub>2</sub> -(45-X)CaO-(XTiO <sub>2</sub> )-5P <sub>2</sub> O <sub>5</sub> , X = 5, 7.5 and 10 wt %)	>90	12 – 50	Non-toxic	87
ZnO and DEP + ZnO NPs	-	15 – 25	Cytotoxicity	88

157

158 Table S3 provides a comparative summary of reported 24-h cell-viability results for  
159 different materials. The table lists the sample type, measured viability percentage, testing  
160 concentration range, and overall cytotoxicity outcome. The BNT-SBT-0.15STZ composition  
161 developed in this work shows >95% cell viability across a wide concentration range (0.1 – 1000  
162  $\mu\text{g}/\text{mL}$ ), indicating a non-toxic response comparable to or better than previously reported materials  
163 such as Co/HA/Ti coatings, HAp/AlN nanocomposites, FeS<sub>2</sub>-modified bioglasses, and SiO<sub>3</sub>CaO-  
164 TiO<sub>3</sub> glass-ceramics. In contrast, ZnO and DEP+/ZnO nanoparticles exhibit cytotoxic behavior

165 even at relatively low concentrations. This comparison highlights the favorable biological  
166 performance and safety of the tested ceramics.



167

168 **Fig. S11.** High resolution XPS spectra of O 1s core levels of (a) unaged and (b) aged of BNT-  
169 SBT-0.15STZ ceramics.

170 Figure S11 shows the high-resolution XPS O 1s spectra of the studied ceramics before  
171 and after thermal aging at 120 °C. The O 1s peak can be separated into three components: lattice  
172 oxygen ( $O_L$ ), adsorption oxygen ( $O_C$ ), and the oxygen vacancy component ( $O_V$ ) located near 531  
173 eV. In the unaged sample, the presence of the  $O_V$  peak at 531.04 eV indicates that oxygen  
174 vacancies were already introduced during sintering. After thermal aging, the  $O_V$  peak position  
175 (531.09 eV) and its relative intensity remain almost unchanged. This indicates that aging at 120  
176 °C, performed without any electrical treatment, is not sufficient to create new oxygen vacancies.  
177 Instead, the results suggest that the existing vacancies undergo only slight redistribution or  
178 stabilization during low-temperature aging. These XPS findings confirm that the changes in  
179 electrical behavior after aging originate from vacancy reorganization rather than an increase in  
180 vacancy concentration.

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