Supporting Information

Wrinkles in 2D TMD Heterostructures: Unlocking Enhanced Hydrogen Evolution Reaction Catalysis

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Figure S1. Atomic structure of the MoSSe/WSSe lateral heterostructure with (a) Armchair-1, (b) Armchair-2, (c) Zigzag-1 and (d) Zigzag-2 interface configuration, the gray ball at the interface indicates the possible adsorption sites for H. Red, blue, yellow and the cyan spheres represent the Mo, W, S and Se atoms, respectively.



Figure S2. Phonon dispersions of the (a) Armchair-1, (b) Armchair-2, (c) Zigzag-1 and (d) Zigzag-2 MoSSe/WSSe lateral heterostructures.



Figure S3. The energy variation of the (a) Armchair-1, (b) Armchair-2, (c) Zigzag-1 and (d) Zigzag-2 MoSSe/WSSe lateral heterostructures in AIMD calculations, the insets demonstrate the snapshots of the studied system.

The geometric expression of Janus heterostructure can be described by using Euler–Bournelli beam model:¹

$$\frac{d^2}{dx^2} \left(D \frac{d^2 w}{dx^2} \right) = 0, \qquad (S1)$$

where w represents the deflection, D is the effective bending stiffness. The boundary conditions are defined by the fractional coordinates of the atoms along the z-axis, with Δz representing the difference between the z-coordinates of all atoms and the initial atom. The resulting deflection can thus be described by Δz as:

$$w(x) = \frac{ML^2}{24D} \left[-\left(\frac{2x-L}{L}\right)^3 + 3\left(\frac{2x-L}{L}\right)^2 - 2\left(\frac{2x-L}{L}\right) \right],$$
 (S2)

where the L is the length of the heterostructure. One can see that the fitted function of the trajectory is consistent, suggesting the wrinkled MoSSe/WSSe interface can be tuned by Euler–Bournelli beam theory.



Figure S4. (a, b) The bond strain and the (c, d) the fitted Euler–Bournelli beam curve of the Armchair-2 MoSSe/WSSe lateral heterostructure by (a, c) 12 and (b, d) 16 unit-cells.

Table. S1 The binding energy (E_b , eV) of H on MoSSe/WSSe heterostructures at different adsorption sites as illustrated in Figure. S1. The Janus site denotes the opposite side of the adsorption position.

	Armchair-1	Armchair-2	Zigzag-1	Zigzag-2
P1	-0.35	0.19	2.15	2.71
P2	-0.30	0.03	2.23	2.78
P3	-0.37	0.02	3.12	3.68
P4	-0.28	0.06	2.41	2.96
P5	-0.47	0.25	2.12	1.73
P6	-0.30	-0.10	2.25	2.81
Janus P1	-0.22	-0.33	2.60	3.15
Janus P2	-0.12	0.03	2.78	3.33
Janus P3	-0.21	-0.14	3.01	3.56
Janus P4	-0.20	0.27	2.43	2.98
Janus P5	-0.22	-0.06	2.52	3.07
Janus P6	-0.20	0.24	2.64	3.19

The HER performance of these MoSSe/WSSe lateral heterostructures is examined by calculating the Gibbs free energy (ΔG_{H^*}) according to:²

$$\Delta G_{\mathrm{H}^*} = \Delta E + \Delta E_{\mathrm{zpe}} + T \Delta S, \qquad (S3)$$

where ΔE denotes the total energy of the H adsorbed MoSSe/WSSe heterostructure, ΔE_{zpe} represents the change in zero-point energy, *T* denotes the temperature and is treated as 298.15 K in this work. The difference in entropy after the adsorption of H is defined by ΔS . "*" is marked as the active site of the HER within the heterostructure. Furthermore, the HER process of the MoSSe/WSSe lateral heterostructure is studied by looking into the two elementary steps:³

$$^{*} + \mathrm{H}^{+} + e^{-} \rightarrow \mathrm{H}^{*}, \tag{S4}$$

$$\mathbf{H}^* + \mathbf{H}^+ + e^- \to \mathbf{H}_2 + ^*. \tag{S5}$$

The projection band center of the MoSSe/WSSe heterostructure is calculated by:

$$\varepsilon = \frac{\int_{-\infty}^{\infty} x\rho(x)dx}{\int_{-\infty}^{\infty} \rho(x)dx}$$
(S6)

The charge density difference $(\Delta \rho)$ of the H-absorbed MoSSe/WSSe lateral heterostructure is calculated by:

$$\Delta \rho = \rho_{\text{H-hetero}} - \rho_{\text{H}} - \rho_{\text{hetero}},\tag{S7}$$

where the $\rho_{\text{H-hetero}}$, ρ_{H} , and ρ_{hetero} represent the charge density of the H absorbed system, individual H, and the pure MoSSe/WSSe heterostructure, respectively.



Figure S5. The orbital contribution of pCOHP between intermediates (H*) on the Armchair-1 and Zigzag-2 MoSSe/WSSe heterostructures.

The atomic virial pressure of atoms in MoSSe/WSSe heterostructure is calculated as:⁴

$$\sigma_{i} = \frac{1}{\Omega_{i}} \left[m_{i} v_{i} \otimes v_{i} + \frac{1}{2} \sum_{j \neq i} F_{ij} \otimes r_{ij} \right]$$
(S8)

where the volume, mass, and velocity vectors for the atom *i* are respresented as Ω_i , m_i , and v_i , respectively. The F_{ij} represents the force of atom *i* applied on atom *j*. The r_{ij} shows the distance between atom *i* and *j*. Furthermore, the symmetric stress tensor (σ_i) can be is composed as σ_{xx} , σ_{yy} , σ_{zz} , σ_{xy} , and σ_{yz} , where σ_{xz} , σ_{yz} , and σ_{zz} can be ignored in the 2D MoSSe/WSSe heterostructure system.⁵



Figure S6. The calculated stress–strain curves of the MoSSe/WSSe lateral heterostructure with Armchair-2 (red line) and Zigzag-2 (navy line) interface.



Figure S7. The obtained H adsorbed Zigzag-2 MoSSe/WSSe heterostructure under the strain of 17%.

References

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