# **Supplementary Information**

## Engineering CoSe<sub>2</sub> phase transition in Br-induced confined space for

## high-performance electromagnetic wave absorption

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#### **Theoretical formula**

With the aid of the associated dielectric loss analyzed by Debye theory, the complex dielectric constant can be expressed as:<sup>1</sup>

$$\varepsilon' = \varepsilon_{\infty} + \frac{(\varepsilon_s - \varepsilon_{\infty})}{1 + (\omega \tau)^2}$$
 (Formula S1)

To clarify the source of the sample dielectric loss, the conduction loss and polarisation loss should be analysed in detail and can be measured using the following equation:<sup>2</sup>

$$\varepsilon^{''} = \varepsilon^{''}_{c} + \varepsilon^{''}_{p} = \frac{\sigma^{*}}{\omega\varepsilon_{0}} + \frac{(\varepsilon_{s} - \varepsilon_{\infty})\omega\tau}{1 + (\omega\tau)^{2}}$$
(Formula S2)

 $\varepsilon_c''$  and  $\varepsilon_p''$  represent the conduction loss and polarization loss, respectively.  $\sigma^*$ stands for electronic conductivity, while  $\varepsilon_0(\varepsilon_0=8.854\times10^{-12} \text{ F/m})$  represents the vacuum dielectric constant.  $\tau$  is the relaxation time, and  $\omega$  ( $\omega=2\pi f$ ) denotes angular frequency.  $\varepsilon_s$  and  $\varepsilon_{\infty}$  signify static permittivity and relative permittivity at infinite frequency, respectively.

The quarter-wavelength theory( $\lambda/4$ ) is used to explore the relationship between the matching thickness of the absorbing material ( $d_m$ ) and frequency( $f_m$ ), which obeys the following equation:<sup>3</sup>

$$d_m = \frac{n\lambda}{4} = \frac{nc}{4f_m\sqrt{|\varepsilon_r||\mu_r|}} (n = 1,3,5...)$$
(Formula S3)

Among them  $\lambda$  is the wavelength of microwave, *and c* is the speed of light in vacuum. Z represents the impedance matching characteristic, while the attenuation coefficient ( $\alpha$ ) can be used to evaluate the combined loss capability of the electromagnetic absorption materials. The equations for both these values are given

below:4

$$Z = \left| \frac{Z_{in}}{Z_0} \right| \qquad (Formula S4)$$
$$\alpha = \frac{\sqrt{2}\pi f}{c} \times \sqrt{\left(\mu^{''} \varepsilon^{''} - \mu^{'} \varepsilon^{'}\right) + \sqrt{\left(\mu^{''} \varepsilon^{''} - \mu^{'} \varepsilon^{'}\right)^2 + \left(\mu^{''} \varepsilon^{'} + \varepsilon^{''} \mu^{'}\right)^2}} \qquad (Formula S5)$$

### **Computational details**

First-principles calculations: We performed the first-principles calculations in the frame of density functional theory (DFT) with the Vienna ab initio simulation package (VASP), using the Perdew-Burke-Ernzerhof (PBE) form of generalized-gradient approximation (GGA) exchange-correlation energy functional. The structure optimizations of Br doping o-CoSe<sub>2</sub>/c-CoSe<sub>2</sub> and Br doping c-CoSe<sub>2</sub> have been carried out by allowing all atomic positions to vary and relaxing lattice parameters until the energy difference of successive atom configurations was less than 10-4eV. The force on each atom in the relaxed structures was less than 0.05 eV/Å. The cutoff energy for the plane-wave basis set was set to 450V. VASPKIT was used to produce K-points, and the slabs were optimized using a  $(1 \times 1 \times 2)$  Monkhorst-Pack grid.



Fig.S1. FESEM images of MF.



**Fig.S2**. (a) FT-IR spectra of ZnCo-ZIF, MF@ ZnCo-ZIF@PF and MF; (b) TG, DTA and DTG of MF@ ZnCo-ZIF@PF; (c) XRD patterns of Co@HC.



**Fig. S3**. FESEM images of (a) MF@ZnCo-ZIF@PF, (b) HC-Co, (c) o-CoSe<sub>2</sub>@HC, (d) c-CoSe<sub>2</sub>-Br<sub>2</sub>@HC, (e) ZnCo-ZIF and (f) C@C.



Fig.S4. Interplanar spacings of the inverse Fourier transform from o-CoSe<sub>2</sub> in o-CoSe<sub>2</sub>@HC.



Fig.S5. Interplanar spacings of the inverse Fourier transform from (a)  $o-CoSe_2$ , (b)  $c-CoSe_2$  in mix-CoSe<sub>2</sub>-Br<sub>1</sub>@HC.



Fig.S6. Interplanar spacings of the inverse Fourier transform from c-CoSe<sub>2</sub> in c-CoSe<sub>2</sub>-

Br<sub>2</sub>@HC.



**Fig.S7**. (a) PDOS of c-CoSe<sub>2</sub> and Br-doped c-CoSe<sub>2</sub>; (b) TDOS of c-CoSe<sub>2</sub> and Br-doped c-CoSe<sub>2</sub>; (c-d) work functions of Br-doped o-CoSe<sub>2</sub> and Br-doped c-CoSe<sub>2</sub>.



Fig.S8. Tanδ<sub>μ</sub> of o-CoSe<sub>2</sub>@HC, mix-CoSe<sub>2</sub>-Br<sub>1</sub>@HC, and c-CoSe<sub>2</sub>-Br<sub>2</sub>@ HC.



Fig.S9. (a-b) RL contour plots of o-CoSe<sub>2</sub>@HC and c-CoSe<sub>2</sub>-Br<sub>2</sub>@HC; (c-d) 3D RL plots of Co@C and Co@HC.



Fig.S10. 2D contour maps of  $|Z_{in}/Z_0|$  of (a) o-CoSe<sub>2</sub>@HC, (b) mix-CoSe<sub>2</sub>-Br<sub>1</sub>@HC and (c) c-CoSe<sub>2</sub>-Br<sub>2</sub>@HC.



**Fig.S11.** (a-b) Three-dimensional radar wave scattering signals of PEC, o-CoSe<sub>2</sub>@HC and c-CoSe<sub>2</sub>-Br<sub>2</sub>@ HC; (c) simulated RCS values of PEC, o-CoSe<sub>2</sub>@HC and c-CoSe<sub>2</sub>-Br<sub>2</sub>@ HC under certain detecting angles; (d) the RCS curves of mix-CoSe<sub>2</sub>-Br<sub>1</sub>@HC with 3.55 mm under typical frequency in the detection range of  $-90^{\circ} \le \theta \le 90^{\circ}$ .

### References

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