

Supplementary Information

Ion dynamics in *closo*-hydroborates $\text{Ag}_2[\text{B}_{10}\text{H}_{10}]$ and $\text{Tl}_2[\text{B}_{10}\text{H}_{10}]$:
 ^1H , ^{11}B and ^{205}Tl NMR studies

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Table S1. Results of the elemental analysis for $\text{Ag}_2[\text{B}_{10}\text{H}_{10}]$

Elements	H	B	Ag
Calc. (wt. %)	3.02	32.38	64.61
Exp. (wt. %)	2.97	32.35	64.59

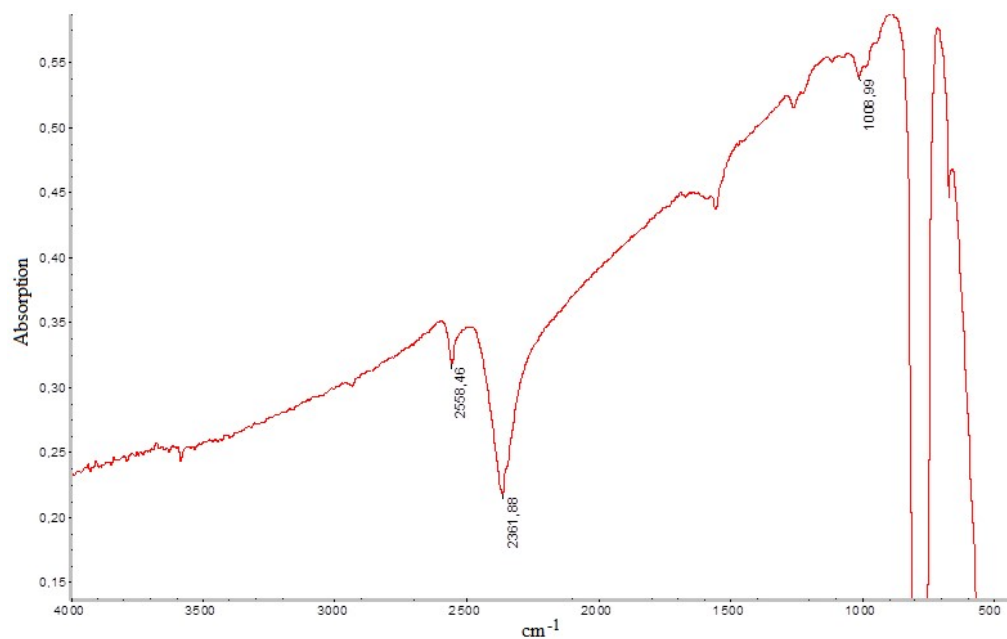


Fig. S1. IR spectrum of $\text{Ag}_2[\text{B}_{10}\text{H}_{10}]$: $\nu(\text{BH})_{\text{MHB}} = 2558 \text{ cm}^{-1}$, $\nu(\text{BH}) = 2362 \text{ cm}^{-1}$, $\delta(\text{BBH}) = 1009 \text{ cm}^{-1}$.

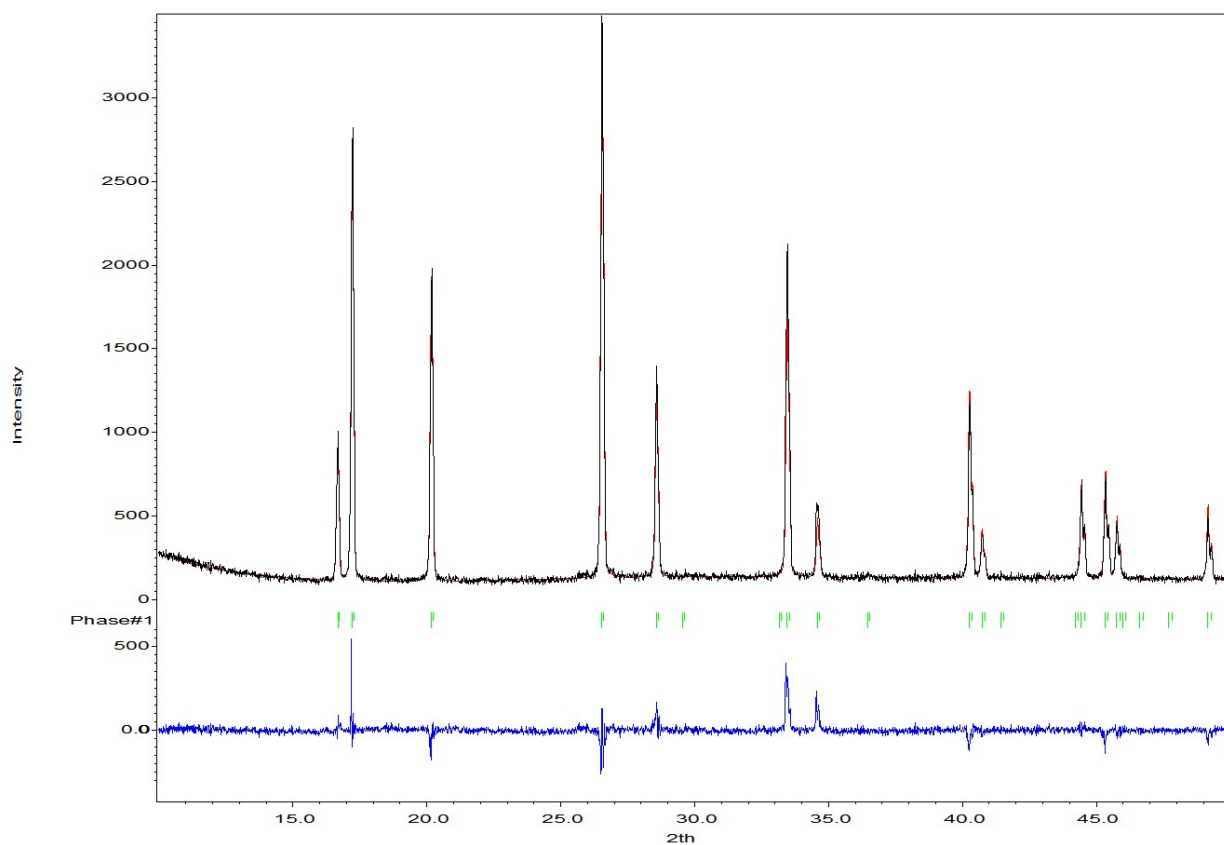


Fig. S2. Room-temperature X-ray powder diffraction pattern for $\text{Ag}_2[\text{B}_{10}\text{H}_{10}]$. Black: experimental; red: calculated; blue: the difference. Vertical bars show the calculated positions of Bragg peaks.

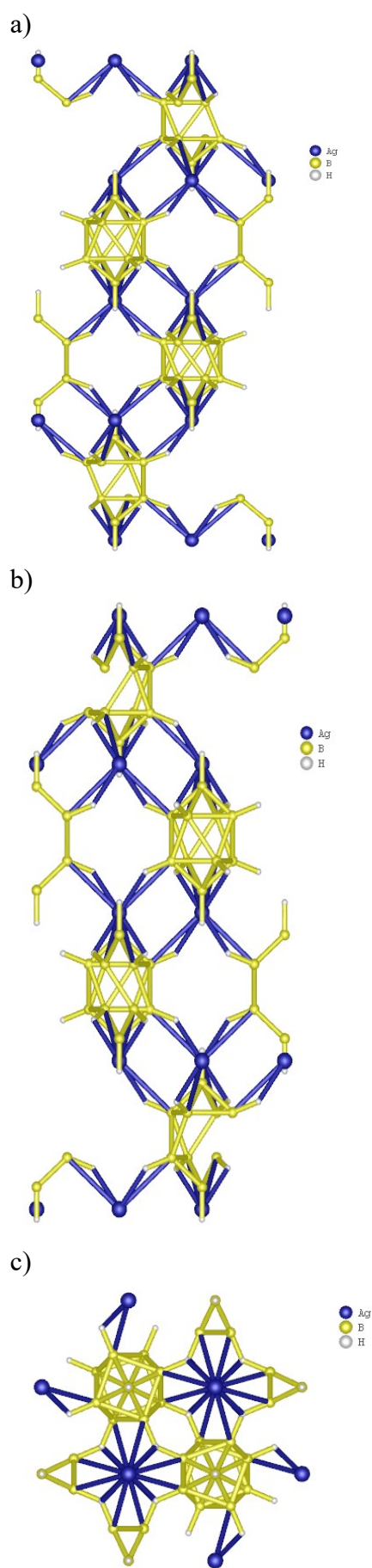


Fig. S3. Projections of the crystal structure of $\text{Ag}_2[\text{B}_{10}\text{H}_{10}]$ along 100 (a), 010 (b), and 001 (c). Data are from cif file 1790981.^{S1}

Expressions used for analysis of the proton spin-lattice relaxation rates

The standard model based on Lorentzian spectral densities^{S2,S3} for fluctuating dipole-dipole interactions relates the proton spin-lattice relaxation rate R_1^H with the parameters of these fluctuations. According to our estimates for the studied compounds, the “rigid-lattice” second moment of the ^1H NMR line is dominated by the $^1\text{H} - ^1\text{H}$ and $^1\text{H} - ^{11}\text{B}$ dipole-dipole interactions, while contributions due to interactions of ^1H with other nuclear spins can be neglected. In this case, the proton spin-lattice relaxation rate can be written as the sum of the $^1\text{H} - ^1\text{H}$ and $^1\text{H} - ^{11}\text{B}$ contributions:

$$R_1^H = \frac{2\Delta M_{HH}\tau_c}{3} \left[\frac{1}{1 + \omega_H^2\tau_c^2} + \frac{4}{1 + 4\omega_H^2\tau_c^2} \right] + \frac{\Delta M_{HB}\tau_c}{2} \left[\frac{1}{1 + (\omega_H - \omega_B)^2\tau_c^2} + \frac{3}{1 + \omega_H^2\tau_c^2} + \frac{6}{1 + (\omega_H + \omega_B)^2\tau_c^2} \right] \quad (\text{S1})$$

where ω_H and ω_B are the resonance frequencies of ^1H and ^{11}B , respectively, and ΔM_{HH} and ΔM_{HB} are the parts of the dipolar second moments due to $^1\text{H} - ^1\text{H}$ and $^1\text{H} - ^{11}\text{B}$ interactions that are caused to fluctuate by the reorientational process. The dipolar correlation times τ_c are simply related to the corresponding mean residence times τ between successive H jumps: $\tau_c = \tau$ for H – H and H – B interactions within the same reorienting group, and $\tau_c = \tau/2$ for the interactions between different reorienting groups; the temperature dependence of τ_i is governed by the Arrhenius relation:

$$\tau = \tau_0 \exp(E_a / k_B T) \quad . \quad (\text{S2})$$

Since each of the H – H and H – B terms in Equation (S1) exhibits nearly the same temperature and frequency dependence, it is practically impossible to determine the amplitude factors ΔM_{HH} and ΔM_{HB} *independently* from the fits to the experimental spin-lattice relaxation rate data. For the studied compounds, the estimates of the “rigid-lattice” second moments show that both the H – H and H – B contributions are of the same order of magnitude. Therefore, for the fits we have assumed that $\Delta M_{HH} = \Delta M_{HB} = \Delta M/2$.

For $\text{Ag}_2[\text{B}_{10}\text{H}_{10}]$, the experimental proton spin-relaxation data suggest the presence of a jump rate distribution, and we have used the model based on a Gaussian distribution of the activation energies.^{S4} For this model, the relaxation rate is expressed as

$$R_1^H = \int R_1^H(E_a)G(E_a, \bar{E}_a, \Delta E_a)dE_a, \quad (\text{S4})$$

where $G(E_a, \bar{E}_a, \Delta E_a)$ is the Gaussian distribution centered on \bar{E}_a with the dispersion ΔE_a , and $R_1^H(E_a)$ is given by equations (S1) and (S2).

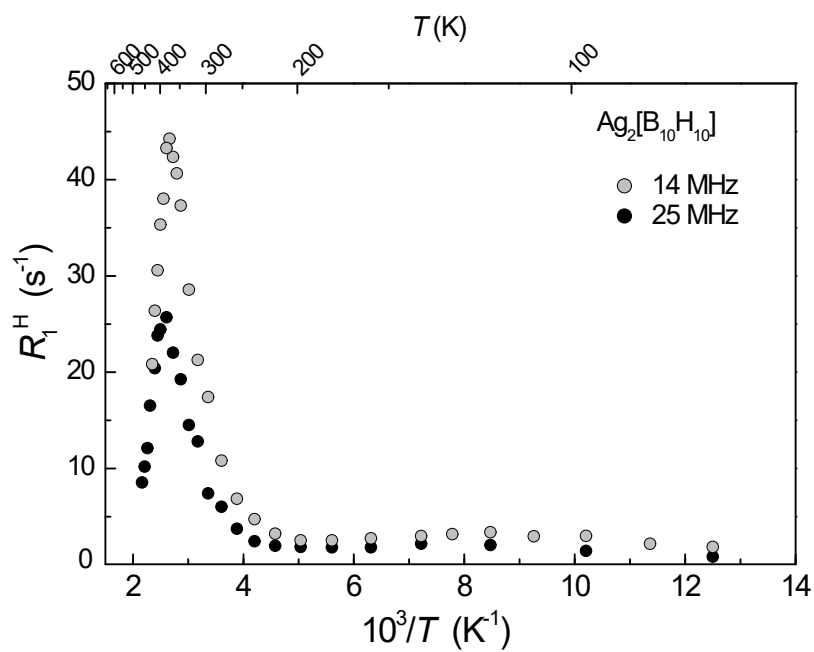


Fig. S4. Proton spin-lattice relaxation rates measured at 14 and 25 MHz for $Ag_2[B_{10}H_{10}]$ at the linear relaxation rate scale.

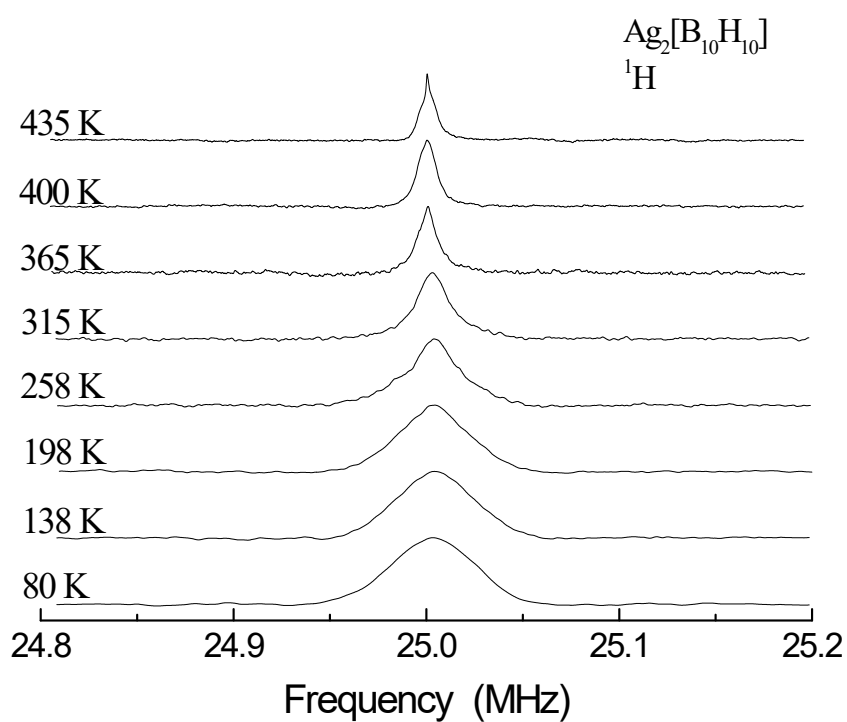


Fig. S5. Evolution of the 1H NMR spectra with temperature.

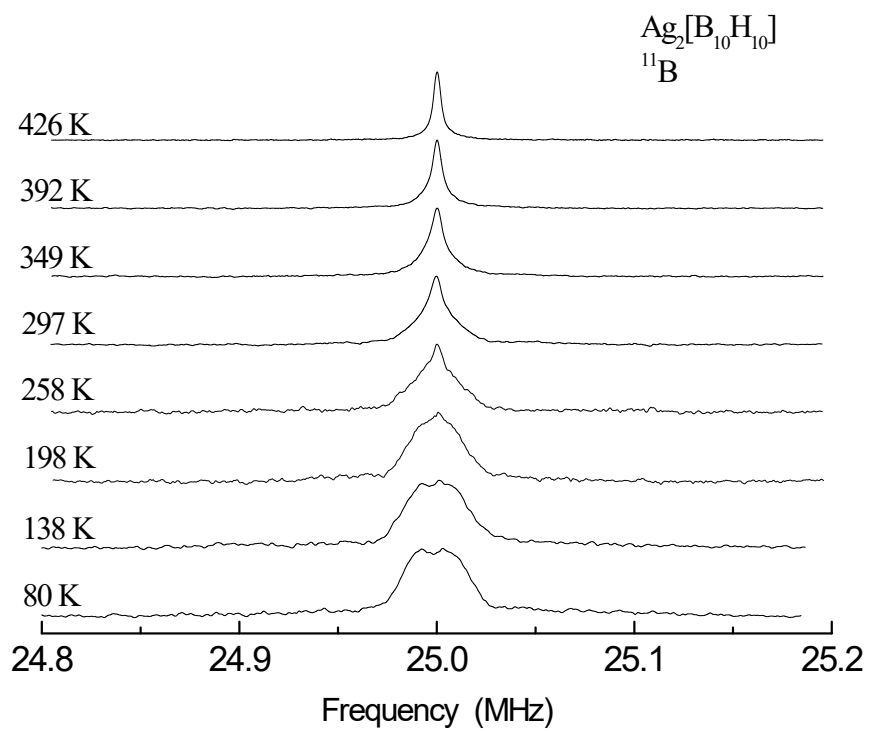


Fig. S6. Evolution of the ^{11}B NMR spectra with temperature.

Table S2. Results of the elemental analysis for $\text{Tl}_2[\text{B}_{10}\text{H}_{10}]$

Elements	H	B	Tl
Calc. (wt. %)	1.91	20.52	77.57
Exp. (wt. %)	1.89	20.48	77.51

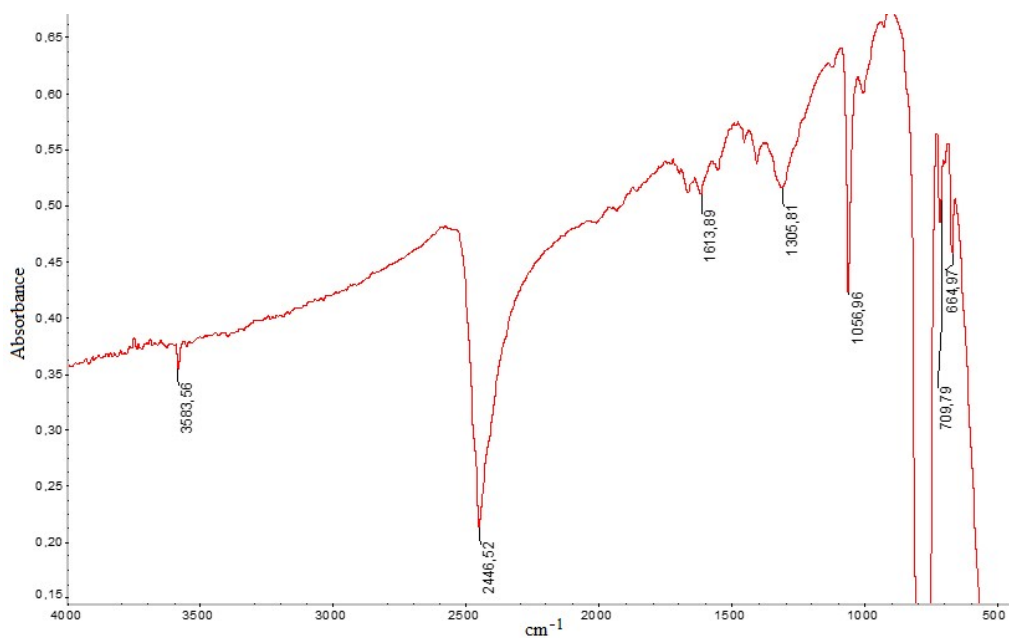


Fig. S7. IR spectrum of $\text{Tl}_2[\text{B}_{10}\text{H}_{10}]$: $\nu(\text{BH}) = 2446 \text{ cm}^{-1}$, $\delta(\text{BBH}) = 1057 \text{ cm}^{-1}$.

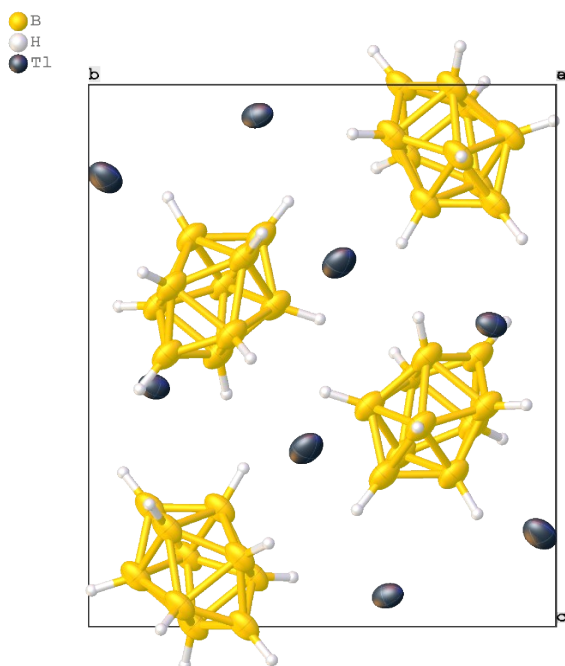


Fig. S8. Schematic view of the unit cell of $\text{Tl}_2[\text{B}_{10}\text{H}_{10}]$ along a axis.^{S5}

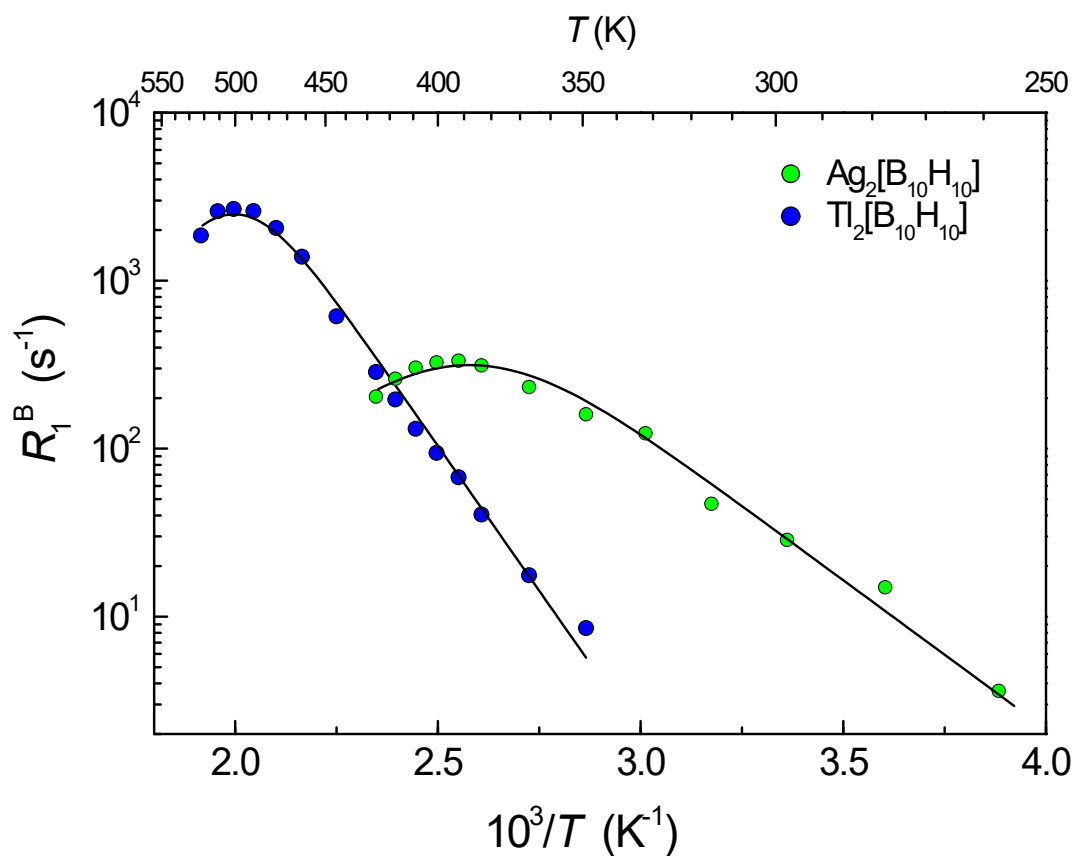


Fig. S9. Combined presentation of the ^{11}B spin-lattice relaxation rates measured at 25 MHz for both $\text{Ag}_2[\text{B}_{10}\text{H}_{10}]$ and $\text{Tl}_2[\text{B}_{10}\text{H}_{10}]$ as functions of the inverse temperature. Solid curves show the fits of the standard model to the data.

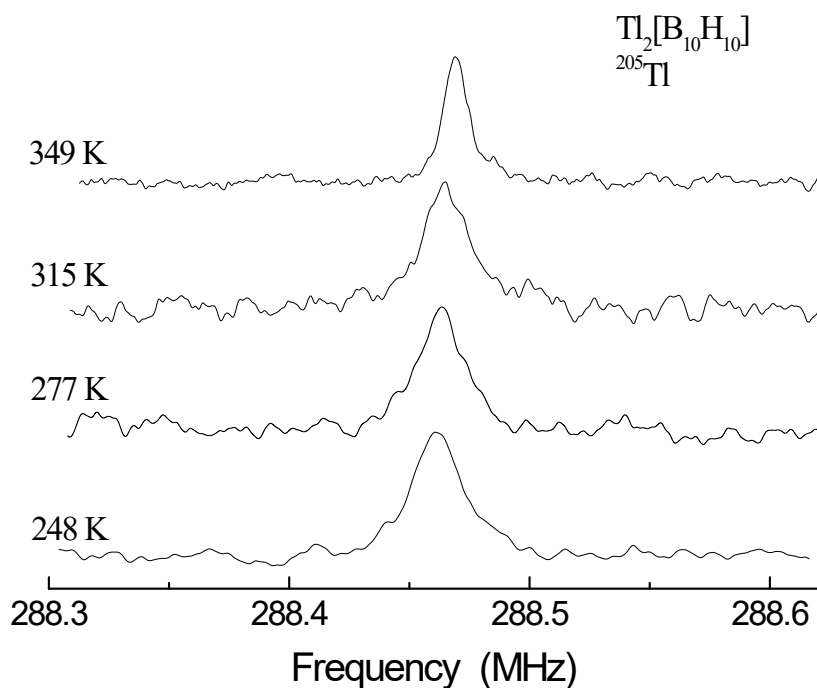


Fig. S10. Evolution of ^{205}Tl NMR spectrum for $\text{Tl}_2[\text{B}_{10}\text{H}_{10}]$ with temperature.

References

- S1 M. Paskevicius, B.R.S. Hansen, M. Jørgensen, B. Richter and T.R. Jensen, *Nat. Commun.*, 2017, **8**, 15136.
- S2 N. Bloembergen, E.M. Purcell, R.M. Pound, *Phys. Rev.*, 1948, **73**, 679-712.
- S3 A. Abragam, *The Principles of Nuclear Magnetism*, Clarendon Press, Oxford, 1961.
- S4 J.T. Markert, E.J. Cotts, R.M. Cotts, *Phys. Rev. B*, 1988, **37**, 6446-6452.
- S5 V.V. Avdeeva, A.S. Kubasov, S.E. Korolenko, L.V. Goeva, E.A. Malinina and N.T. Kuznetsov, *Russ. J. Inorg. Chem.*, 2022, **67**, 582-590.