

Supplementary Information for

## 2D TeThCl with High Solar-to-Hydrogen Efficiency and Strain-Tunable Excitonic Properties

### Calculation methods

#### (1) Young's modulus and Poisson's ratio

For 2D material, its in-plane Young's modulus ( $Y$ ) and Poisson's ratio ( $\nu$ ) can be calculated using the formulas as follows:[1]

$$Y(\theta) = \frac{C_{11}C_{22} - C_{12}^2}{C_{11}\sin^4\theta + \left(\frac{C_{11}C_{22} - C_{12}^2}{C_{66}} - 2C_{12}\right)\sin^2\theta\cos^2\theta + C_{22}\cos^4\theta}$$

(S1)

$$\nu(\theta) = \frac{C_{12}\sin^4\theta - \left(C_{11} + C_{22} - \frac{C_{11}C_{22} - C_{12}^2}{C_{66}}\right)\sin^2\theta\cos^2\theta + C_{12}\cos^4\theta}{C_{11}\sin^4\theta + \left(\frac{C_{11}C_{22} - C_{12}^2}{C_{66}} - 2C_{12}\right)\sin^2\theta\cos^2\theta + C_{22}\cos^4\theta}$$

(S2)

Where  $C_{ij}$  is the elastic constant, and  $\theta$  is the angle with respect to  $a$ -axis.

#### (2) Carrier mobility

The carrier mobility ( $\mu$ ) can be calculated using the deformation potential theory (DPT) method, which can be described as:[2-4]

$$\mu = \frac{e\hbar^2 C^{2D}}{k_B T m^* m_d E_l^2} \quad (\text{S3})$$

where  $C^{2D}$  is in-plane stiffness,  $m^*$  is the carrier effective mass,  $m_d$  is the density-of-states effective mass defined as  $m_d = \sqrt{m_a^* m_b^*}$  ( $a$  and  $b$  represent the directions of the  $a$ -axis and  $b$ -axis), and  $E_l$  is the deformation potential constant defined by  $E_l = \partial E_{edge} / \partial E$ ,  $E_{edge}$  is the energy near CBM (for electrons) and VBM (for holes), and  $E$  is the uniaxial strain along  $a$ -axis and  $b$ -axis. Here we set the total uniaxial strain is  $-2\% \sim 2\%$  and with the step size of  $1\%$ .

According to Lang et al. 's research[4], the acoustic phonon-limited mobility (APM,  $\bar{\mu}$ ) of anisotropic 2D semiconductor can be corrected by its elastic constants ( $C_{ij}$ ), which can be expressed as below:

$$\bar{\mu} = \frac{e\hbar^3}{k_B T (m_x^*)^{\frac{3}{2}} (m_y^*)^{\frac{1}{2}}} \left( \frac{A + B - \sqrt{A^2 - B^2}}{B\sqrt{A^2 - B^2}} \right) \times \left( \frac{I + J - \sqrt{I^2 - J^2}}{J\sqrt{I^2 - J^2}} \right) \quad (\text{S4})$$

$$A = \bar{E}_l^2 + \frac{(\Delta E_l)^2}{2}, \quad B = \bar{E}_l \Delta E_l \quad (\text{S5})$$

$$\bar{E}_l = \frac{E_{lx} + E_{ly}}{2}, \quad \Delta E_l = \frac{E_{ly} - E_{lx}}{2} \quad (\text{S6})$$

$$\bar{C} = \frac{C_{11} + C_{22}}{2}, \quad \Delta C = \frac{C_{22} - C_{11}}{2} \quad (\text{S7})$$

$$I = \frac{1}{\sqrt{\bar{C}^2 - (\Delta C)^2}}, \quad J = \frac{\bar{C}}{\Delta C} \left( \frac{1}{\bar{C}} - \frac{1}{\sqrt{\bar{C}^2 - (\Delta C)^2}} \right) \quad (\text{S8})$$

### (3) Gibbs free energy

We calculated the Gibbs free energy ( $\Delta G$ ) for the adsorption of intermediates (\*H, \*OH, \*O, and \*OOH) on the surface of the monolayer

during the photocatalytic water splitting process, governed by the following relationships:[5-7]

$$\Delta G = \Delta E + \Delta E_{ZPE} - T\Delta S + \Delta G_{pH} - eU \quad (\text{S9})$$

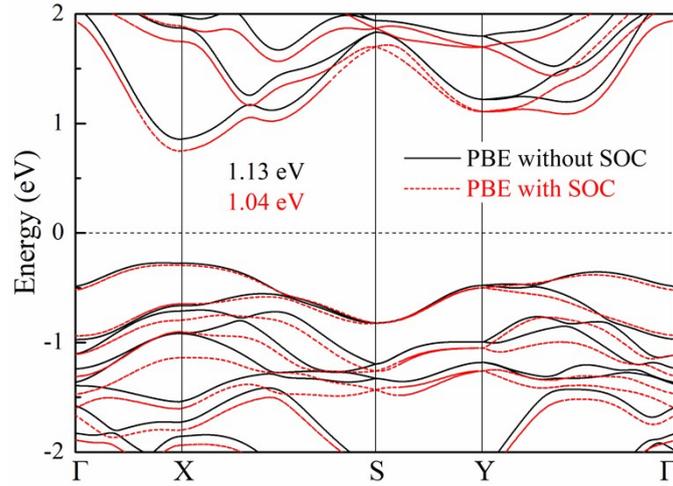
Where  $\Delta E$  represents the differential adsorption energy for each intermediate reaction, and  $\Delta E_{ZPE}$  is the difference in zero-point energy (ZPE).  $T$  denotes the temperature of 298.15 K, while  $\Delta S$  represents the entropy change obtained from vibrational frequency calculations.  $\Delta G_{pH}$  accounts for the change in free energy influenced by the pH value.  $U$  is the electrode potential relative to the standard water reduction potential.

For HER, the expression for  $T\Delta S$  is approximately  $-0.24$  eV. Therefore, the Gibbs free energy  $\Delta G_{*H}$  of  $*H$  can be simplified as follows:

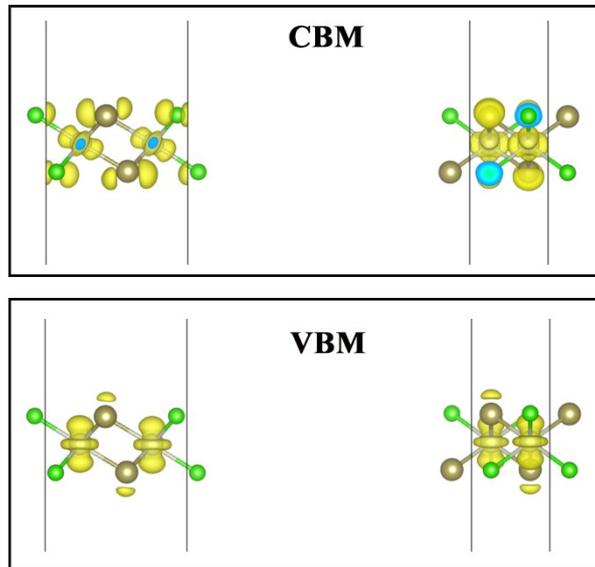
$$\Delta G_{*H} = \Delta E + 0.24 \quad (\text{S10})$$

**Fig. S1** Calculated cleavage energy ( $E_f$ , eV) of monolayer ThTeCl. The computation involves the total energies of the bulk unit cell ( $E_{bulk}$ , eV) and monolayer unit cell ( $E_{mono}$ , eV), and the in-plane area ( $S_{in-plane}$ ,  $\text{\AA}^2$ ). For comparison, the cleavage energies of monolayer graphene and MoS<sub>2</sub> are also provided.

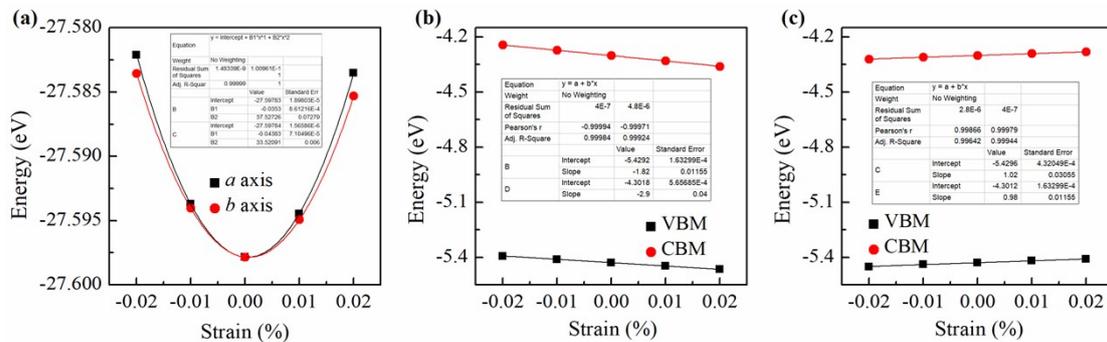
Materials	$E_{bulk}$	$E_{mono}$	$S_{in-plane}$	$E_f$
Graphene	-60.26	-19.98	5.26	0.33
MoS <sub>2</sub>	-113.22	-56.46	8.81	0.27
RhTeCl	-322.19	-160.68	21.60	0.31



**Fig. S1** Electronic band structures of the monolayer calculated using the PBE functional, with (red curve) and without (black curve) the inclusion of spin-orbit coupling (SOC).

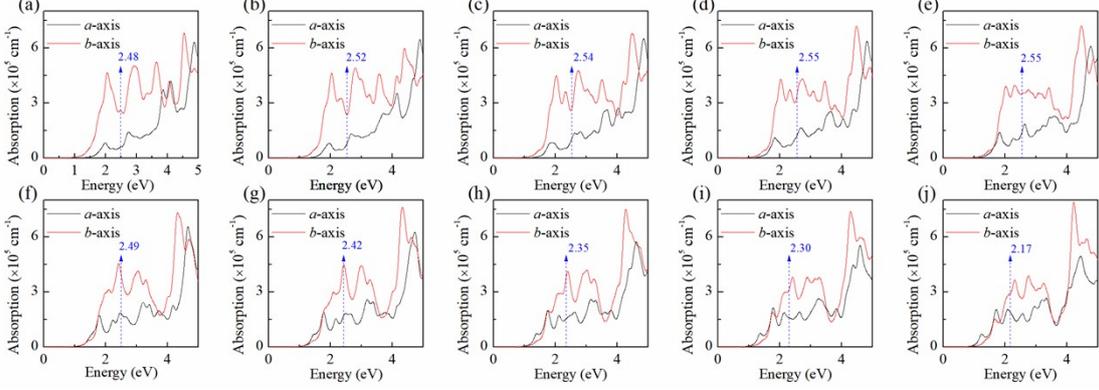


**Fig. S2** Real-space distribution of the electron wave functions at the valence band maximum (VBM) and conduction band minimum (CBM) of monolayer ThTeCl. Two different side-view perspectives are provided.

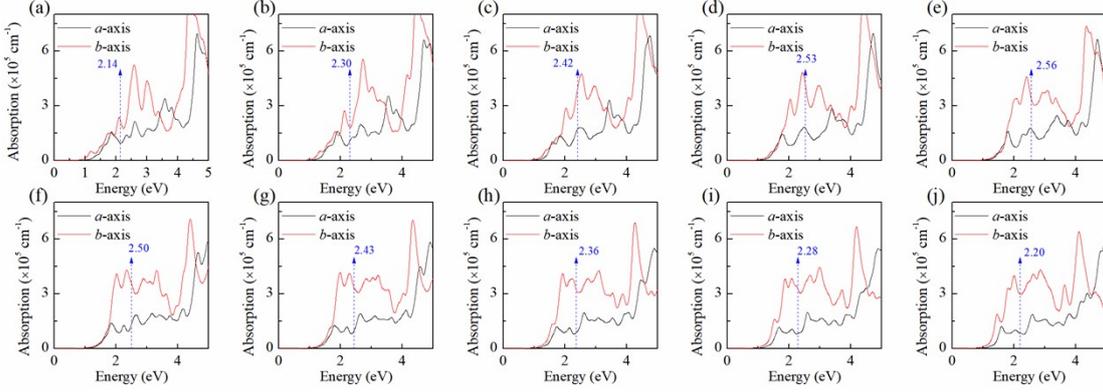


**Fig. S3** Supporting data for the carrier mobility calculation within the deformation potential theory

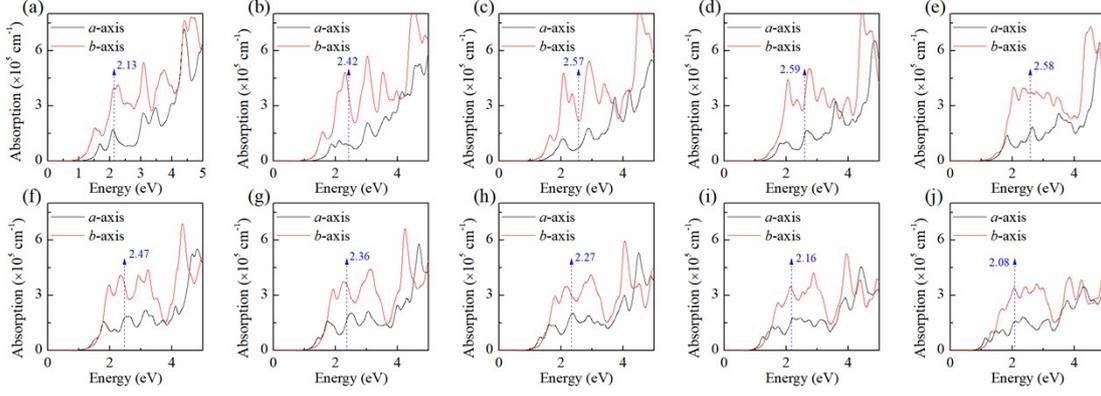
framework. (a) The quadratic fitting of the total energy as a function of uniaxial strain, used to derive the planar stiffness. (b, c) Linear fittings of the VBM and CBM energies versus strain along the  $a$ - and  $b$ -axes, used to obtain the deformation potential constants for holes and electrons, respectively.



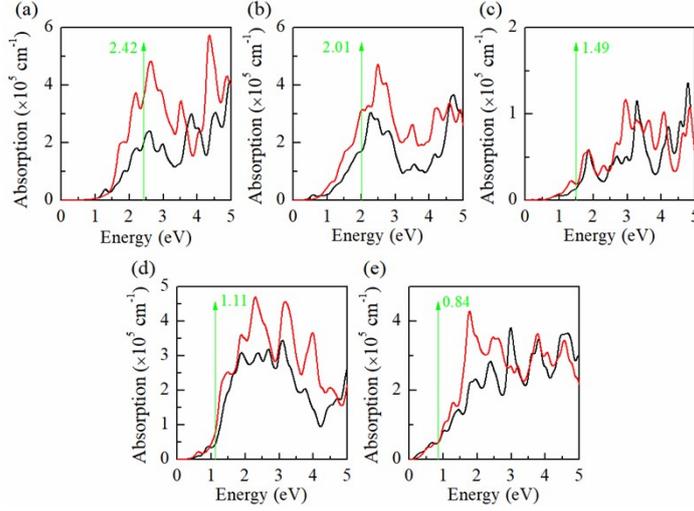
**Fig. S4** Strain-tunable anisotropic optical absorption in monolayer RhTeCl. The absorption coefficients along the  $a$ -axis (black curves) and  $b$ -axis (red curves) are presented. The sequence from panel (a) to (j) illustrates the spectral evolution under uniaxial strain applied along the  $a$ -axis, ranging from -10% to +10% with a 2% step.



**Fig. S5** Strain-tunable anisotropic optical absorption in monolayer RhTeCl. The absorption coefficients along the  $a$ -axis (black curves) and  $b$ -axis (red curves) are presented. The sequence from panel (a) to (j) illustrates the spectral evolution under uniaxial strain applied along the  $b$ -axis, ranging from -10% to +10% with a 2% step.



**Fig. S6** Strain-tunable anisotropic optical absorption in monolayer RhTeCl. The absorption coefficients along the  $a$ -axis (black curves) and  $b$ -axis (red curves) are presented. The spectra from (a) to (j) correspond to the results under biaxial strain ranging from -10% to +10% with an increment of 2%.



**Fig. S7** Modulation of the anisotropic optical absorption in monolayer RhTeCl by an external electric field. The absorption coefficients along the  $a$ -axis (black curves) and  $b$ -axis (red curves) are presented. The sequence from panel (a) to (e) corresponds to the calculated spectra under electric fields of 0.1, 0.2, 0.3, 0.4, and 0.5 V/Å, respectively.

**Table S2** The total energy  $E_t$ , the position of VBM ( $E_V$ ) and CBM ( $E_C$ ) with respect to vacuum level at uniaxial strains of -4% ~ 4%, where the step size is 1%.

Direction	Strain	$E_t$ (eV)	$E_V$ (eV)	$E_C$ (eV)
$a$ -axis	-4%	-27.582	-5.393	-4.245
	-2%	-27.594	-5.411	-4.272
	0%	-27.598	-5.429	-4.301
	2%	-27.594	-5.447	-4.330
	4%	-27.584	-5.466	-4.361
$b$ -axis	-4%	-27.584	-5.451	-4.321
	-2%	-27.594	-5.439	-4.311
	0%	-27.598	-5.429	-4.301
	2%	-27.595	-5.419	-4.291

**Table S3** The calculation details of the Gibbs free energy change of the intermediate product in monolayer during the OER. Here,  $E_{DFT}$  represents the total energy of the intermediate product adsorbed at the most probable position on the superlattice surface,  $G_{tot}$  represents the total Gibbs free energy, and the  $T$  is set to 298.15 K.

Materials	$E_{DFT}$	$\Delta E_{zpe}$	$T\Delta S$	Molecular	$G_{tot}$	$\Delta G (U = 0.0 V)$	$\Delta G (U = 1.23 V)$
G*	-165.657	0.000	0.000	H <sub>2</sub> O	-194.125	0.000	0.000
G*OH	-174.829	0.318	0.163	H <sub>2</sub> O + (H <sup>+</sup> + e <sup>-</sup> )	-192.311	1.815	0.585
G*O	-170.711	0.062	0.080	H <sub>2</sub> O + 2(H <sup>+</sup> + e <sup>-</sup> )	-191.770	2.355	-0.105
G*OOH	-179.008	0.406	0.193	3(H <sup>+</sup> + e <sup>-</sup> )	-189.005	5.120	1.430
G*	-165.657	0.000	0.000	O <sub>2</sub> + 4(H <sup>+</sup> + e <sup>-</sup> )	-189.205	4.920	0.000

## References

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