

## Supplementary Materials

### **Size-dependent electrobending in piezoceramics mediated by gradient defect dipoles**

Haoyu Gu<sup>1</sup>, Zehua Deng<sup>2</sup>, Siqing He<sup>1</sup>, Hu Zhao<sup>1</sup>, Shaoqi Huang<sup>1</sup>, Shuwen Zhang<sup>1</sup>,  
Linlan Li<sup>1</sup>, Shubao Shao<sup>1,\*</sup> and Minglong Xu<sup>1,\*\*</sup>

*1. State Key Laboratory for Strength and Vibration of Mechanical Structure, School  
of Aerospace Engineering, Xi'an Jiaotong University, Xi'an, 710049, China*

*2. Shanghai Institute of Satellite Engineering, Shanghai 200240, China*

\* Corresponding author. Email: [ssb\\_xjtu@xjtu.edu.cn](mailto:ssb_xjtu@xjtu.edu.cn)

\*\* Corresponding author. Email: [mlxu@xjtu.edu.cn](mailto:mlxu@xjtu.edu.cn)

## Formula derivation section

### Determination of parameters in general solution

Based on Kirchhoff's thin plate theory, the governing equation for the circular thin plate is given by

$$D\nabla^4 w(r) = q(r), \#(S1)$$

where  $D = \frac{Et^3}{12(1-\nu^2)}$  represents the bending stiffness,  $\nu$  is Poisson's ratio,  $w(r)$  denotes the deflection, and  $q(r)$  is the applied force. Since the circular thin plate is solely subjected to a bending moment, i.e.  $q(r) = 0$ , Eq. S1 can be expressed as

$$\frac{d}{dr}\left(r\frac{d}{dr}\left(\frac{1}{r}\frac{d}{dr}\left(r\frac{dw}{dr}\right)\right)\right) = 0. \#(S2)$$

The general solution to Eq. S2 is given by

$$w(r) = C_1 r^2 \ln r + C_2 r^2 + C_3 \ln r + C_4, \#(S3)$$

here  $C_1$ ,  $C_2$ ,  $C_3$ , and  $C_4$  are pending parameters that can be determined by boundary conditions.

Since the deflection is finite at the center of the disk,  $C_3 = 0$ , otherwise the function  $w(r)$  diverges at  $r = 0$ . The bending moment  $M(r)$  is also finite at  $r = 0$ , and

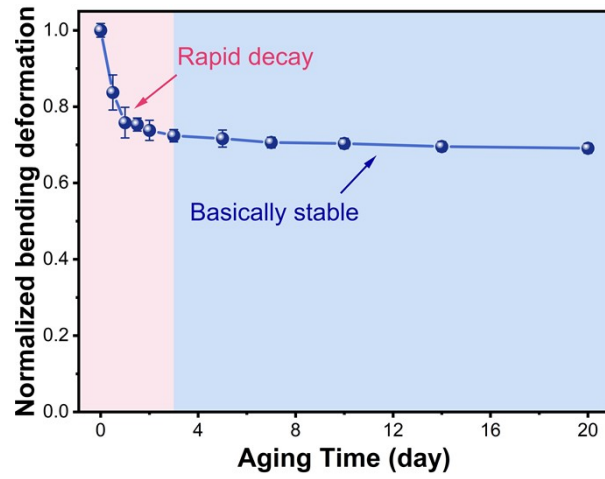
$$M(r) = -D\left(\frac{d^2 w}{dr^2} + \frac{\nu dw}{r dr}\right) = -D[2C_1(1+\nu)\ln r + (3+\nu)C_1 + 2C_2(1+\nu)]. \#(S4)$$

Therefore,  $C_1 = 0$ , otherwise  $M(r)$  also diverges at  $r = 0$ . Furthermore, at  $r = R$ ,  $w(R) = 0$  and  $M(R) = M$ , i.e.

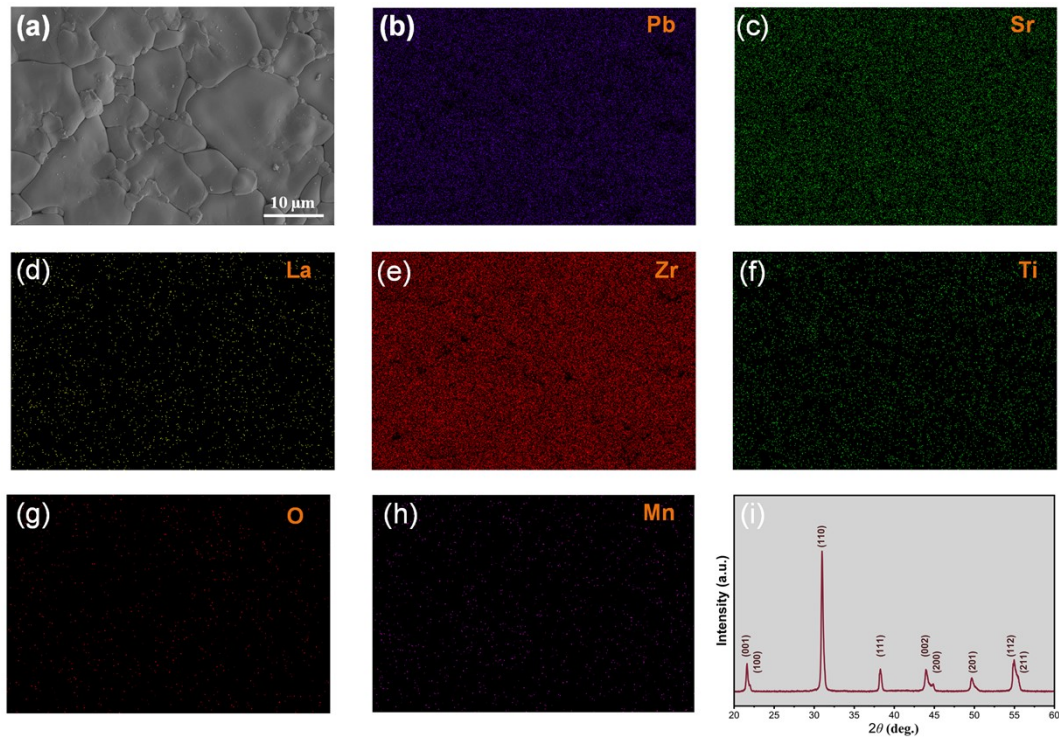
$$\begin{cases} C_2 R^2 + C_4 = 0 \\ -2D(1+\nu)C_2 = M \end{cases} \#(S5)$$

Based on the above boundary conditions, the pending parameters can be determined as follows:

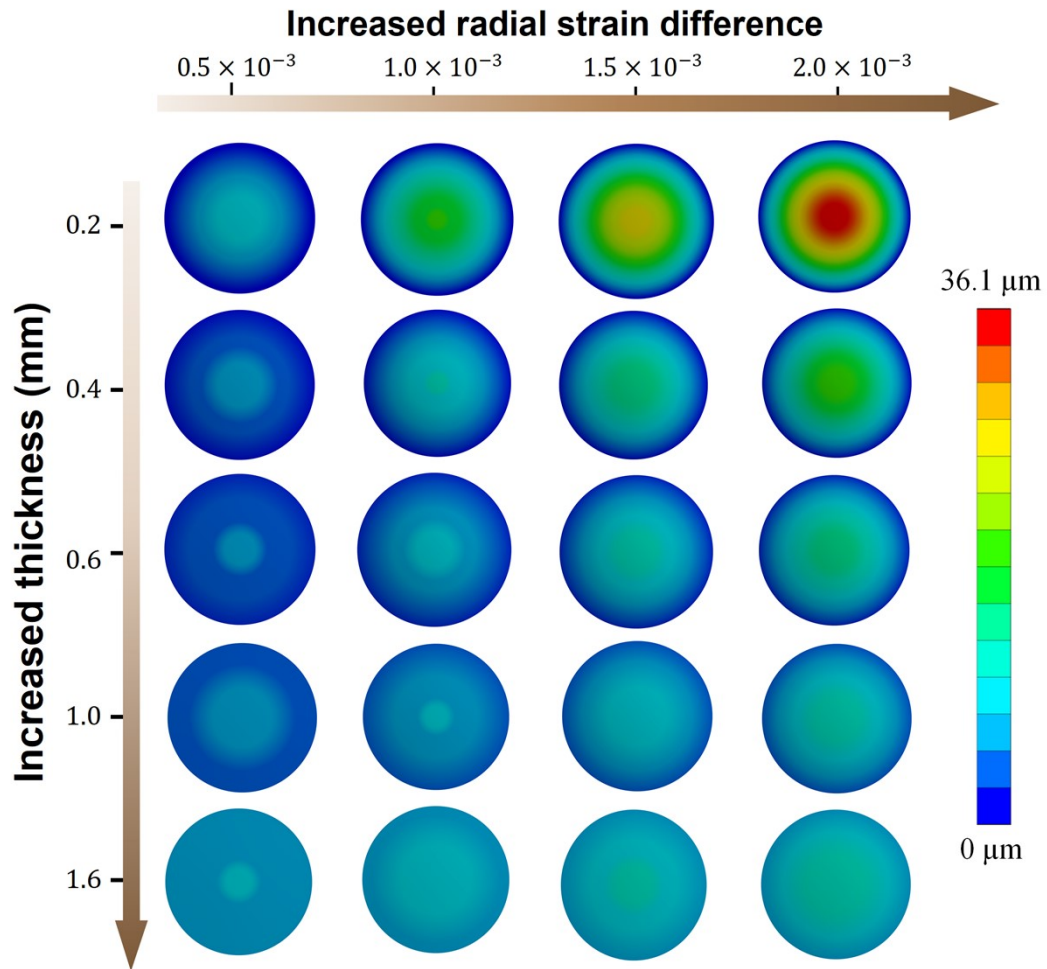
$$\begin{cases} C_1 = C_3 = 0 \\ C_2 = -\frac{M}{2D(1+\nu)} \#(S6) \\ C_4 = \frac{MR^2}{2D(1+\nu)} \end{cases}$$



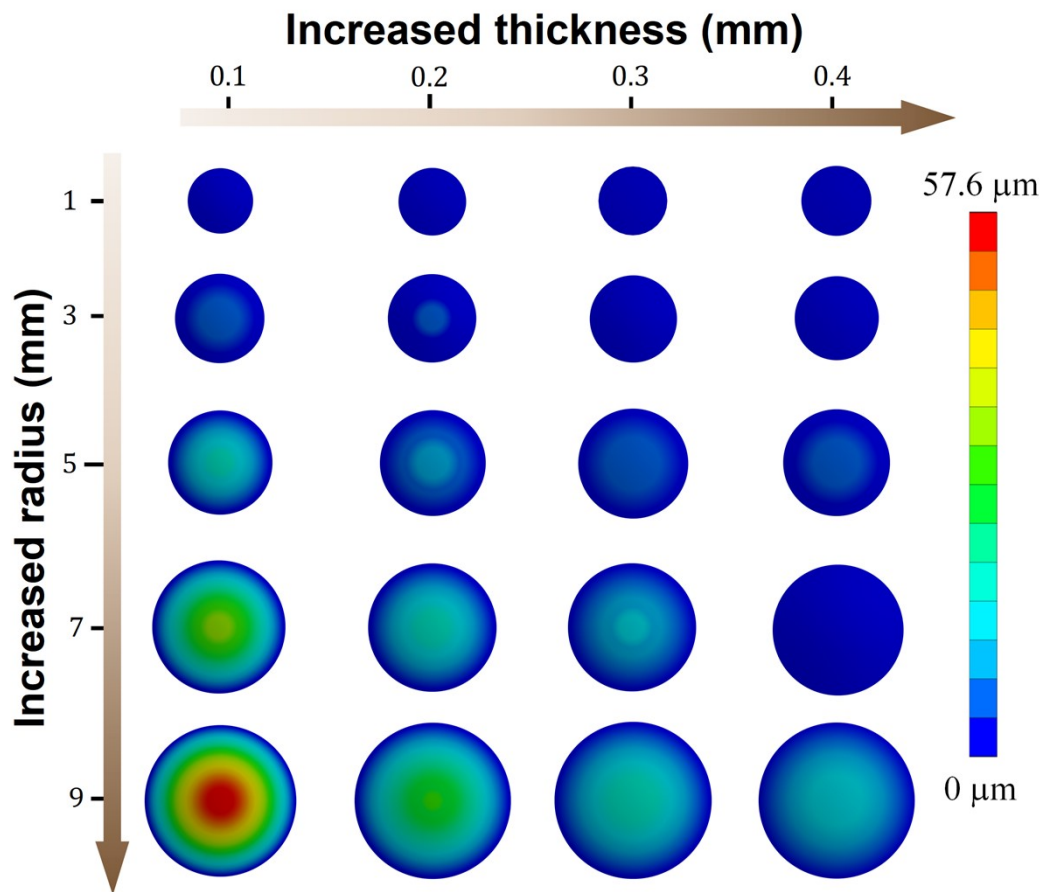
**Fig. S1.** The relationship between normalized bending deformation and aging time.



**Fig. S2.** Structure characterization of PSLZT-0.10Mn ceramics. (a) SEM images of the ceramic surface. (b-h) The in-situ energy-dispersive spectrum (EDS) mapping of Pb, Sr, La, Zr, Ti, O and Mn, respectively. (i) XRD patterns of PSLZT-0.10Mn ceramic.



**Fig. S3.** FEM results of samples with various radial strain difference and thickness, with a radius of 10 mm.



**Fig. S4.** FEM results of samples with various thickness and radius, with a radial strain difference of  $2.0 \times 10^{-3}$ .

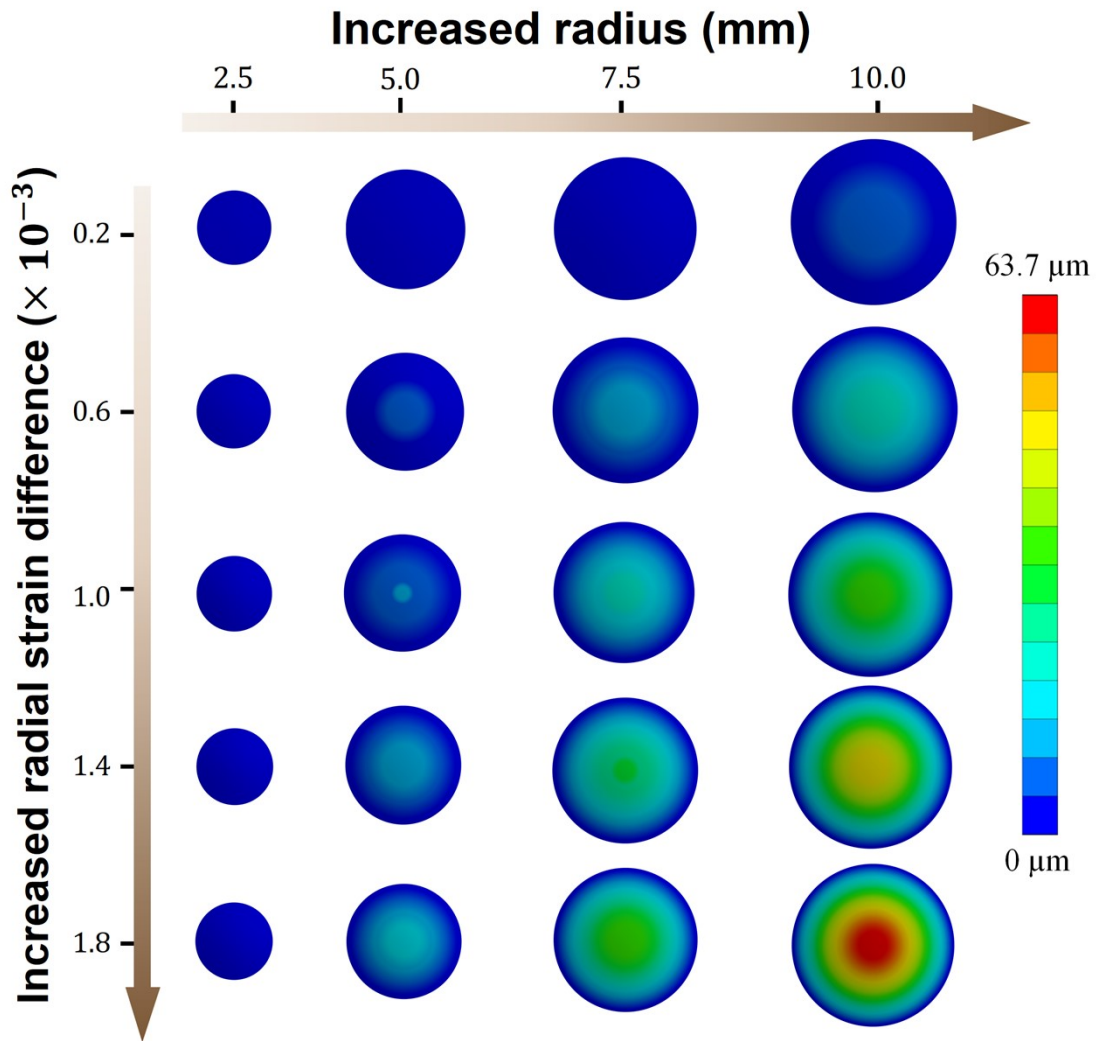
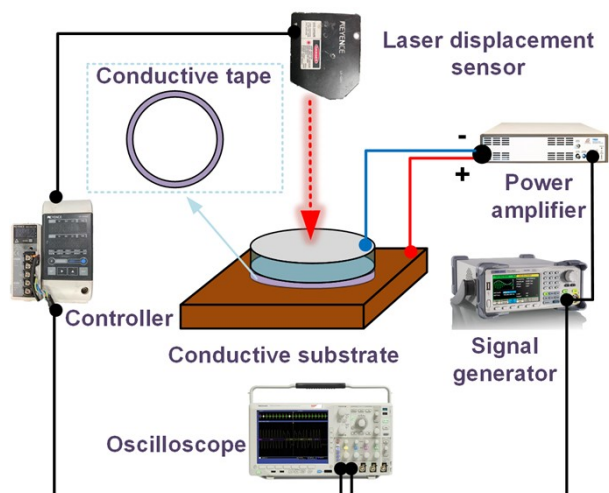
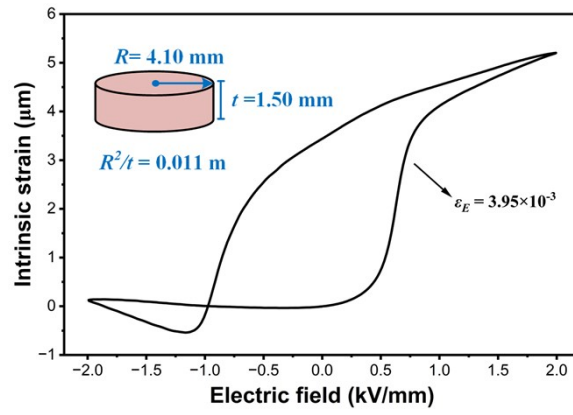


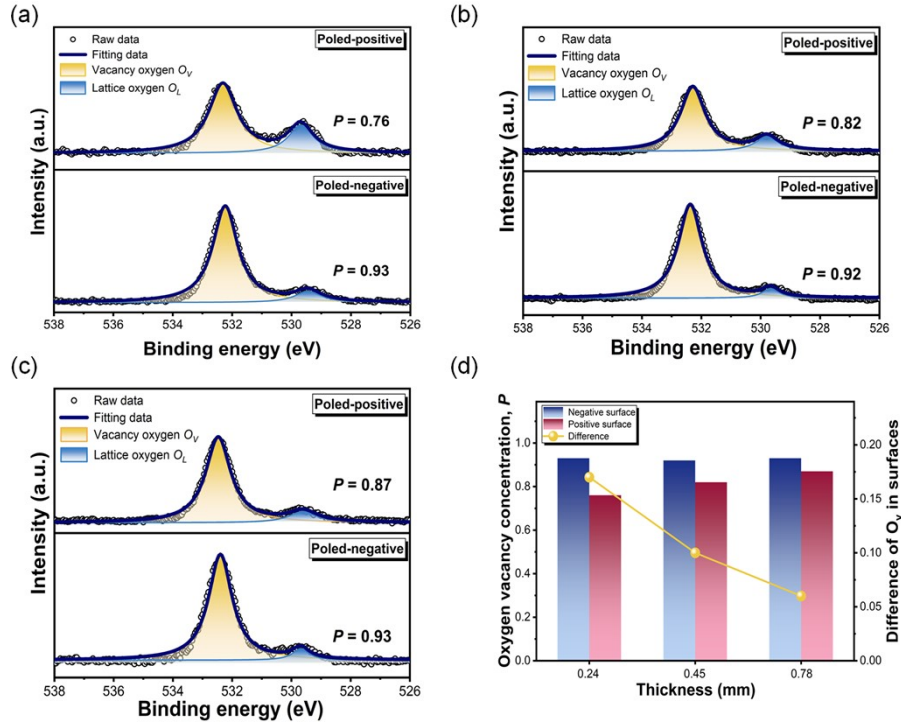
Fig. S5. FEM results of samples with various radial strain difference and radius, with a thickness of 0.1 mm.



**Fig. S6.** Schematic diagram of the custom-built displacement measurement system.



**Fig. S7.** Intrinsic  $S$ - $E$  curve of PSLZT-0.10Mn piezoceramic under an electric field of  $\pm 2$  kV/mm.



**Fig. S8.** The thickness dependence of oxygen vacancy concentration distribution. XPS of the O1 s orbitals for positive and negative surface of samples with (a)  $t = 0.24$  mm, (b)  $t = 0.45$  mm, and (c)  $t = 0.78$  mm. (d) The concentration and difference of oxygen vacancies on the positive and negative surfaces vary with the thickness.