

Supporting Information

1 Computational Methods

The frequency-dependent second-order nonlinear polarization P induced by an optical field can be expressed as:

$$P_a(2\omega) = \chi_{abc}^{(2)}(2\omega; \omega, \omega) E_b(\omega) E_c(\omega), \quad (1)$$

where $\chi_{abc}^{(2)}$ is the SHG susceptibility tensor (the superscript (2) denotes second-order nonlinearity), and a, b, c are Cartesian indices. $E_b(\omega)$ is the component of the incident electric field at frequency ω along direction b . We employ the *length-gauge* formalism for the calculation of SHG response, in which the total susceptibility $\chi_{abc}^{(2)}$ can be decomposed into three major contributions: the interband term $\chi_{abc}^{\text{inter}}$, the intraband term $\chi_{abc}^{\text{intra}}$, and the modulation term χ_{abc}^{mod} . Their relationship is given by

$$\chi_{abc}^{(2)}(2\omega; \omega, \omega) = \chi_{abc}^{\text{inter}}(2\omega; \omega, \omega) + \chi_{abc}^{\text{intra}}(2\omega; \omega, \omega) + \chi_{abc}^{\text{mod}}(2\omega; \omega, \omega). \quad (2)$$

Their explicit forms are as follows:

$$\chi_{abc}^{\text{inter}}(2\omega; \omega, \omega) = \frac{e^3}{\hbar^2} \sum_{nml} \int_{\text{BZ}} \frac{d\mathbf{k}}{4\pi^3} \frac{r_{nm}^a \{r_{ml}^b r_{ln}^c\}}{\omega_{ln} - \omega_{ml}} \left(\frac{2f_{nm}}{\omega_{mn} - 2\omega} + \frac{f_{ml}}{\omega_{ml} - \omega} + \frac{f_{ln}}{\omega_{ln} - \omega} \right), \quad (3)$$

$$\begin{aligned} \chi_{abc}^{\text{intra}}(2\omega; \omega, \omega) &= \frac{e^3}{\hbar^2} \int_{\text{BZ}} \frac{d\mathbf{k}}{4\pi^3} \left[\sum_{nml} \omega_{mn} r_{nm}^a \{r_{ml}^b r_{ln}^c\} \left(\frac{f_{nl}}{\omega_{ln}^2 (\omega_{ln} - \omega)} - \frac{f_{lm}}{\omega_{ml}^2 (\omega_{ml} - \omega)} \right) \right. \\ &\quad \left. - 8i \sum_{nm} \frac{f_{nm} r_{nm}^a \{\Delta_{mn}^b r_{mn}^c\}}{\omega_{mn}^2 (\omega_{mn} - 2\omega)} + 2 \sum_{nml} \frac{f_{nm} r_{nm}^a \{r_{ml}^b r_{ln}^c\} (\omega_{ml} - \omega_{ln})}{\omega_{mn}^2 (\omega_{mn} - 2\omega)} \right], \end{aligned} \quad (4)$$

$$\begin{aligned} \chi_{abc}^{\text{mod}}(2\omega; \omega, \omega) &= \frac{e^3}{2\hbar^2} \int_{\text{BZ}} \frac{d\mathbf{k}}{4\pi^3} \left[\sum_{nml} \frac{f_{nm}}{\omega_{mn}^2 (\omega_{mn} - \omega)} (\omega_{nl} r_{lm}^a \{r_{mn}^b r_{nl}^c\} - \omega_{lm} r_{nl}^a \{r_{ml}^b r_{mn}^c\}) \right. \\ &\quad \left. - i \sum_{nm} \frac{f_{nm} r_{nm}^a \{r_{mn}^b \Delta_{mn}^c\}}{\omega_{mn}^2 (\omega_{mn} - \omega)} \right]. \end{aligned} \quad (5)$$

Here, $\omega_{nm} = \omega_n - \omega_m$ denotes the energy difference between bands n and m , while $f_{nm} = f_n - f_m$ represents the difference of their Fermi–Dirac occupation factors. $r_{nm} = v_{nm}/(i\omega_{nm})$ is the position-operator matrix element, and $\{r_{ml}^b r_{ln}^c\} = \frac{1}{2}(r_{ml}^b r_{ln}^c + r_{ml}^c r_{ln}^b)$ is the symmetrized product. $\Delta_{mn}^b = v_{mm}^b - v_{nn}^b$ corresponds to the difference in the diagonal velocity matrix elements along the b direction. In the practical calculation, a small phenomenological broadening factor δ is introduced as $\omega \rightarrow \omega + i\delta$ with $\delta = 0.1$ eV, in order to avoid divergences when the photon energy coincides with band transition resonances.

2 Band Structure With Concentration of Defects

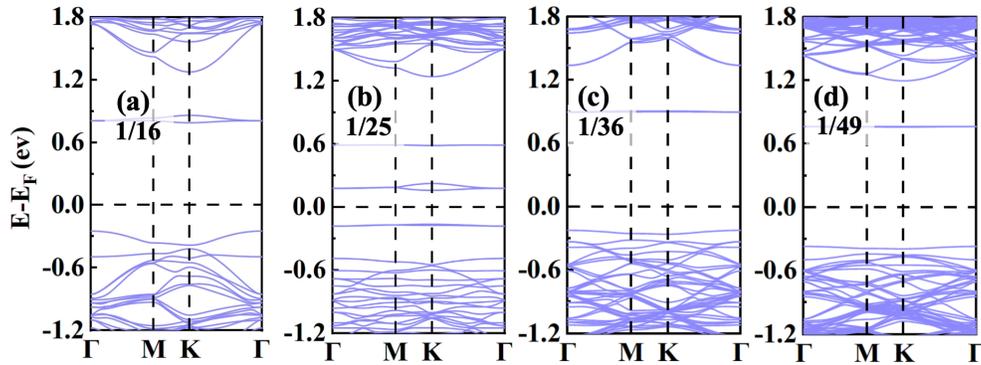


Figure 1: Band structures of the V_S defect type at different defect concentrations. (a) Band structure at 1/16 concentration. (b) Band structure at 1/25 concentration. (c) Band structure at 1/36 concentration. (d) Band structure at 1/49 concentration.

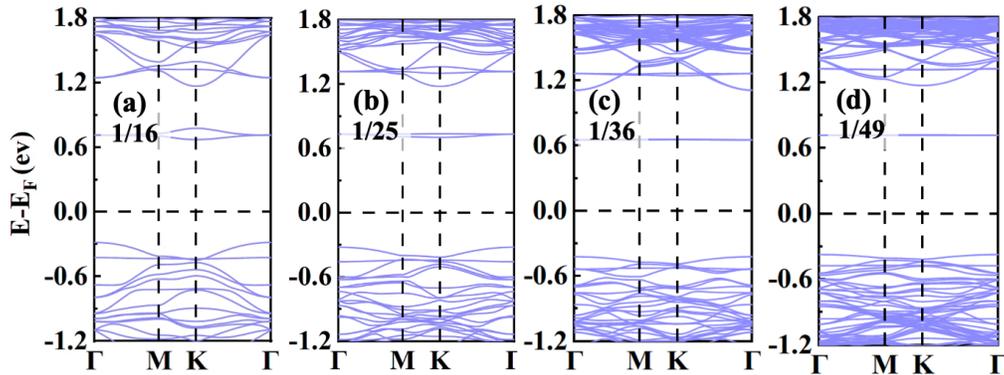


Figure 2: Band structures of the V_{S2} defect type at different defect concentrations. (a) Band structure at 1/16 concentration. (b) Band structure at 1/25 concentration. (c) Band structure at 1/36 concentration. (d) Band structure at 1/49 concentration.

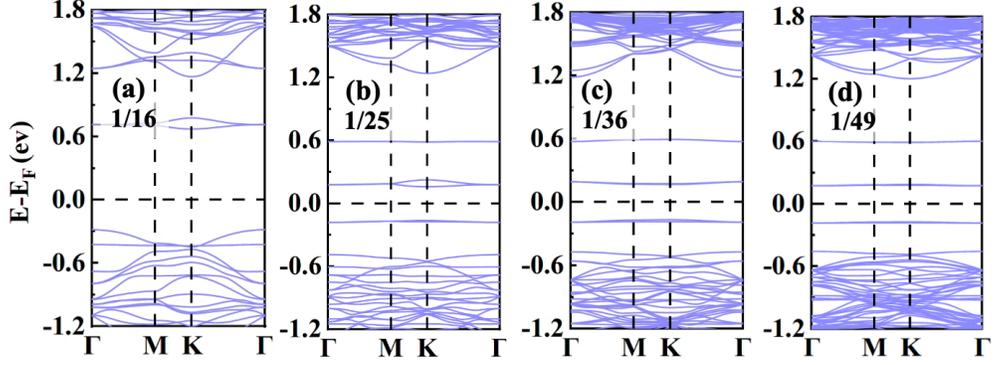


Figure 3: Band structures of the V_{Mo} defect type at different defect concentrations. (a) Band structure at 1/16 concentration. (b) Band structure at 1/25 concentration. (c) Band structure at 1/36 concentration. (d) Band structure at 1/49 concentration.

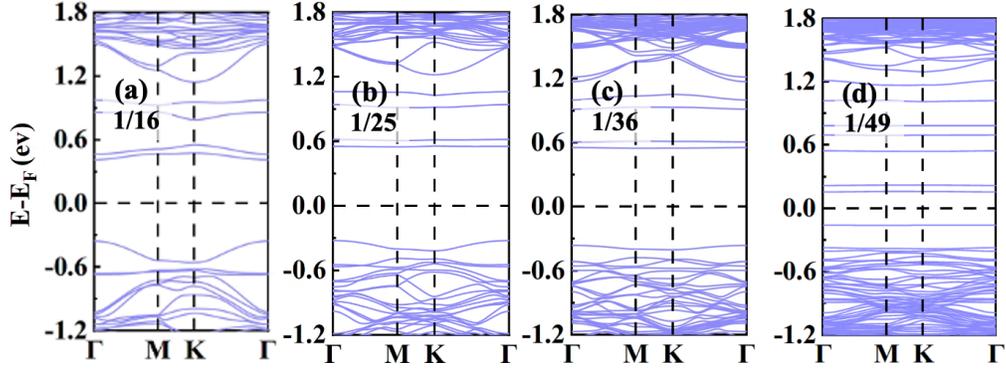


Figure 4: Band structures of the $V2_{S2}$ defect type at different defect concentrations. (a) Band structure at 1/16 concentration. (b) Band structure at 1/25 concentration. (c) Band structure at 1/36 concentration. (d) Band structure at 1/49 concentration.

3 SHG variation with defect concentration

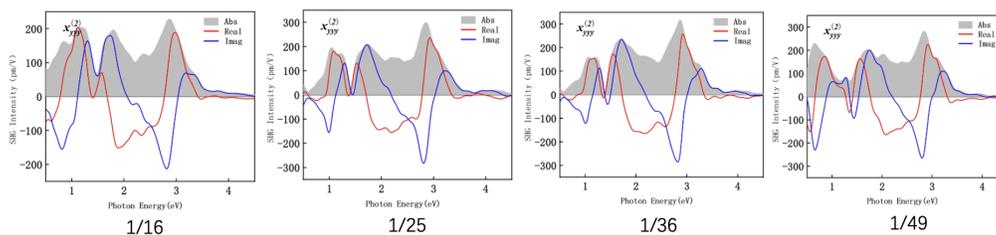


Figure 5: SHG response of the V_{2S_2} defect type at different defect concentrations.

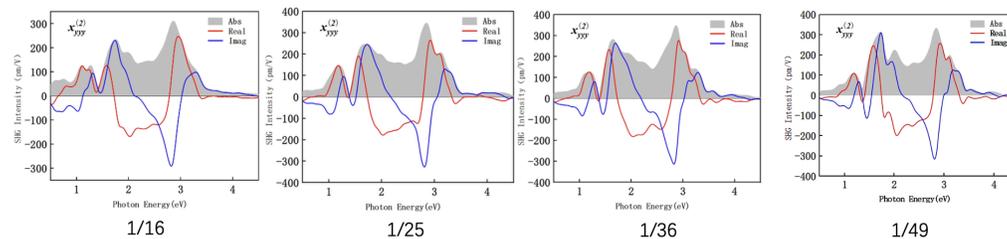


Figure 6: SHG response of the V_S defect type at different defect concentrations.

4 All SHG components with different defect types in 4X4 supercell

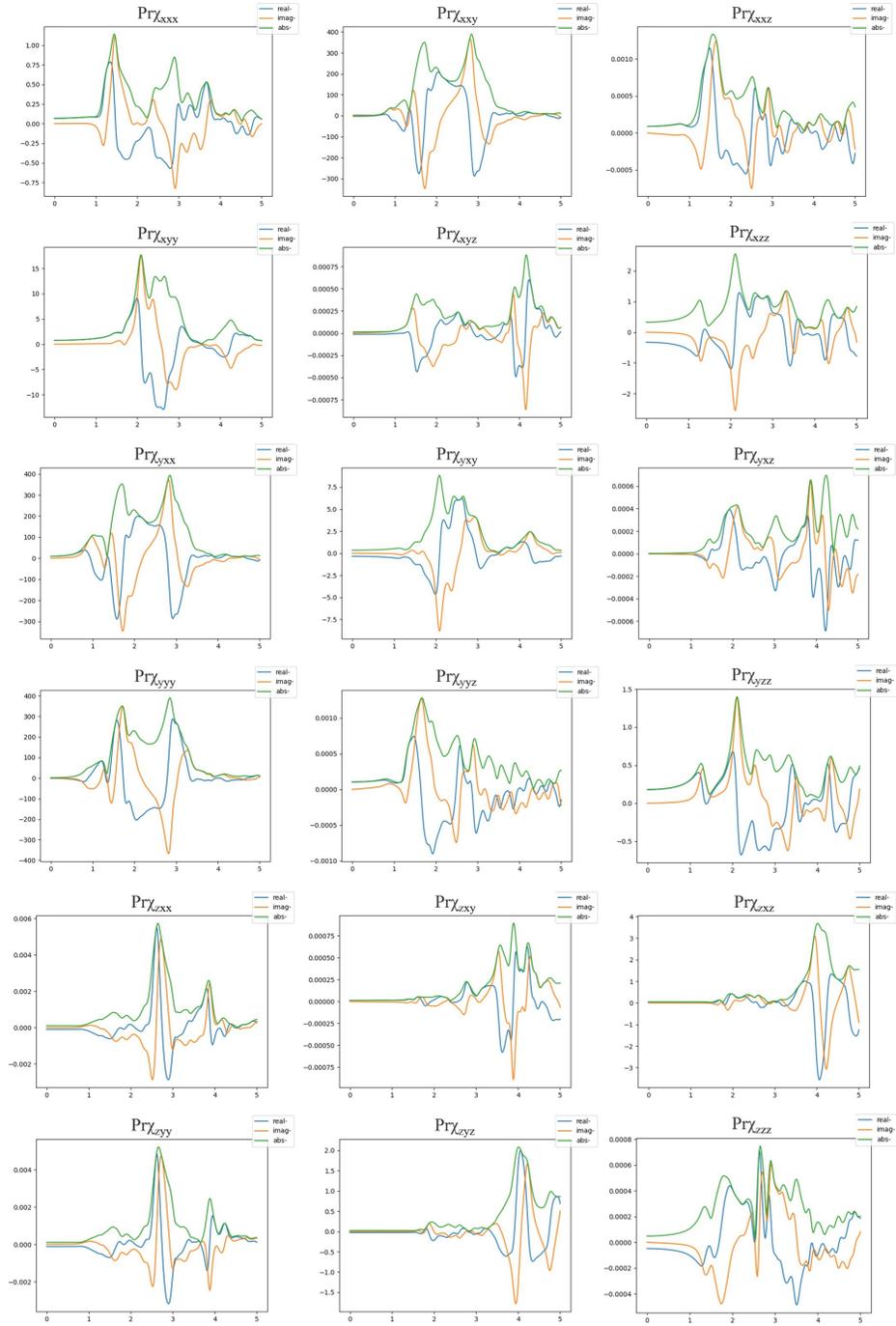


Figure 7: All SHG components with pristine MoS₂ in 4X4 supercell

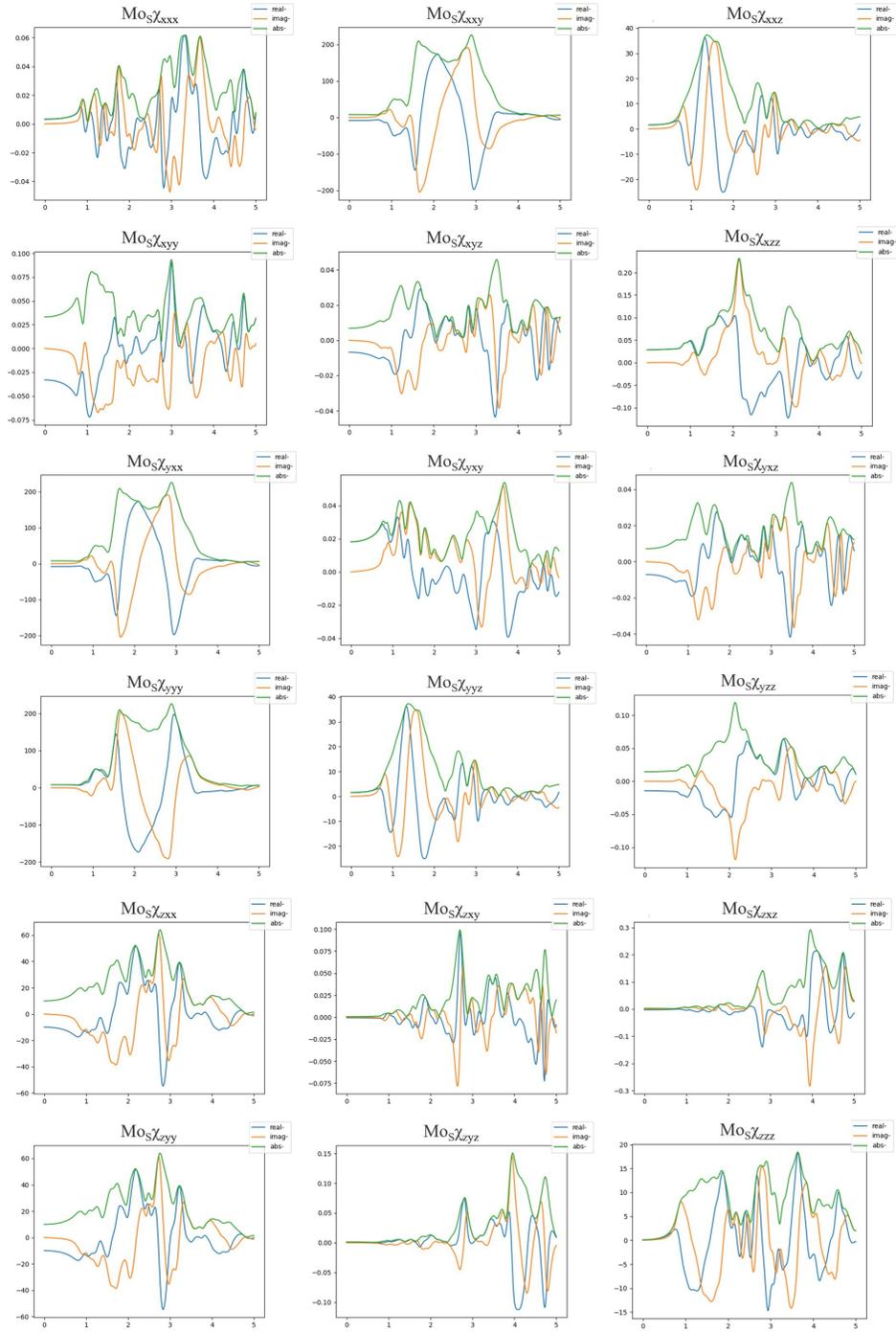


Figure 8: All SHG components with Mo_S with defects in 4X4 supercell

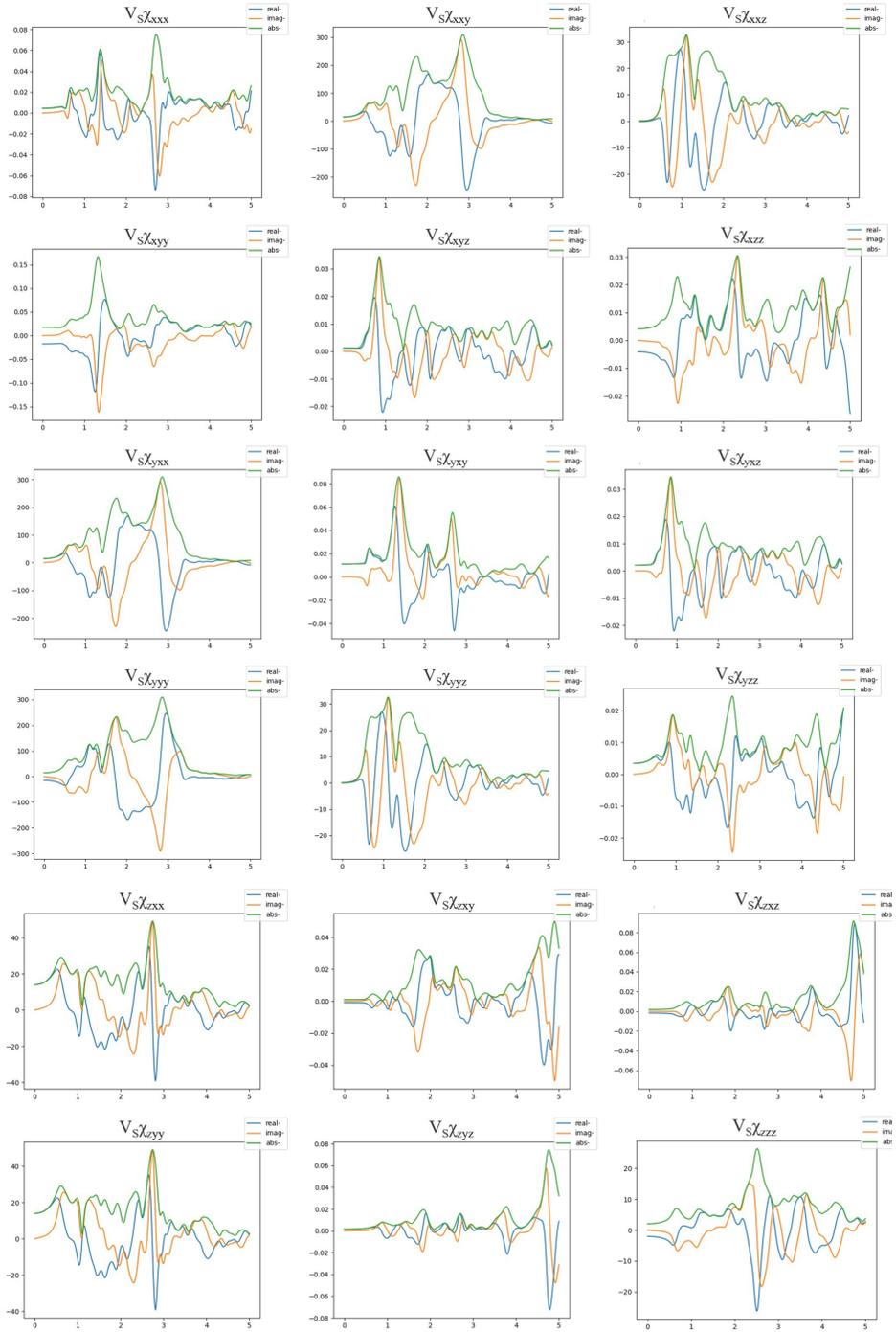


Figure 9: All SHG components with V_S with defects in 4X4 supercell

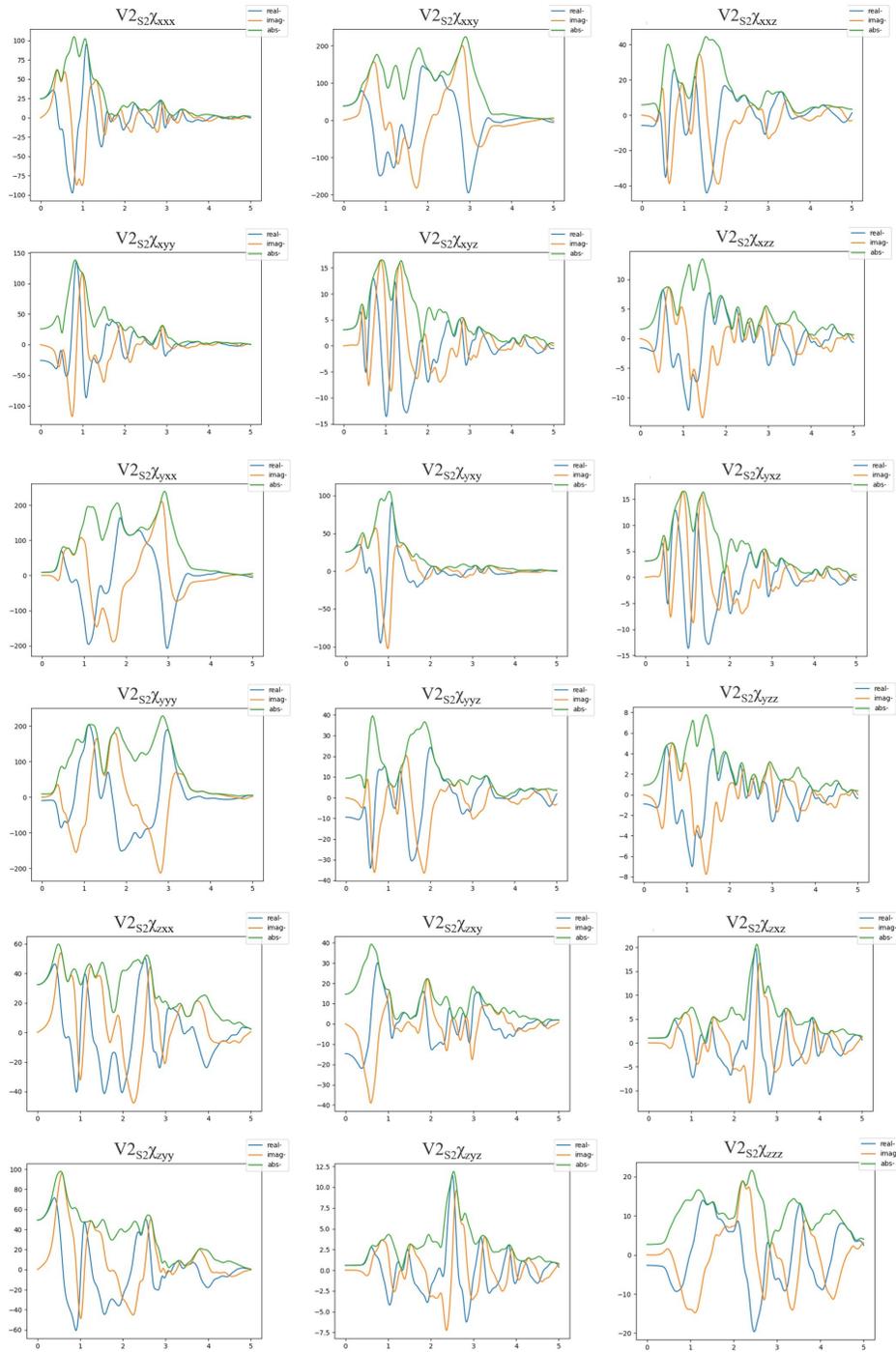


Figure 10: All SHG components with $V2_{S2}$ with defects in 4X4 supercell

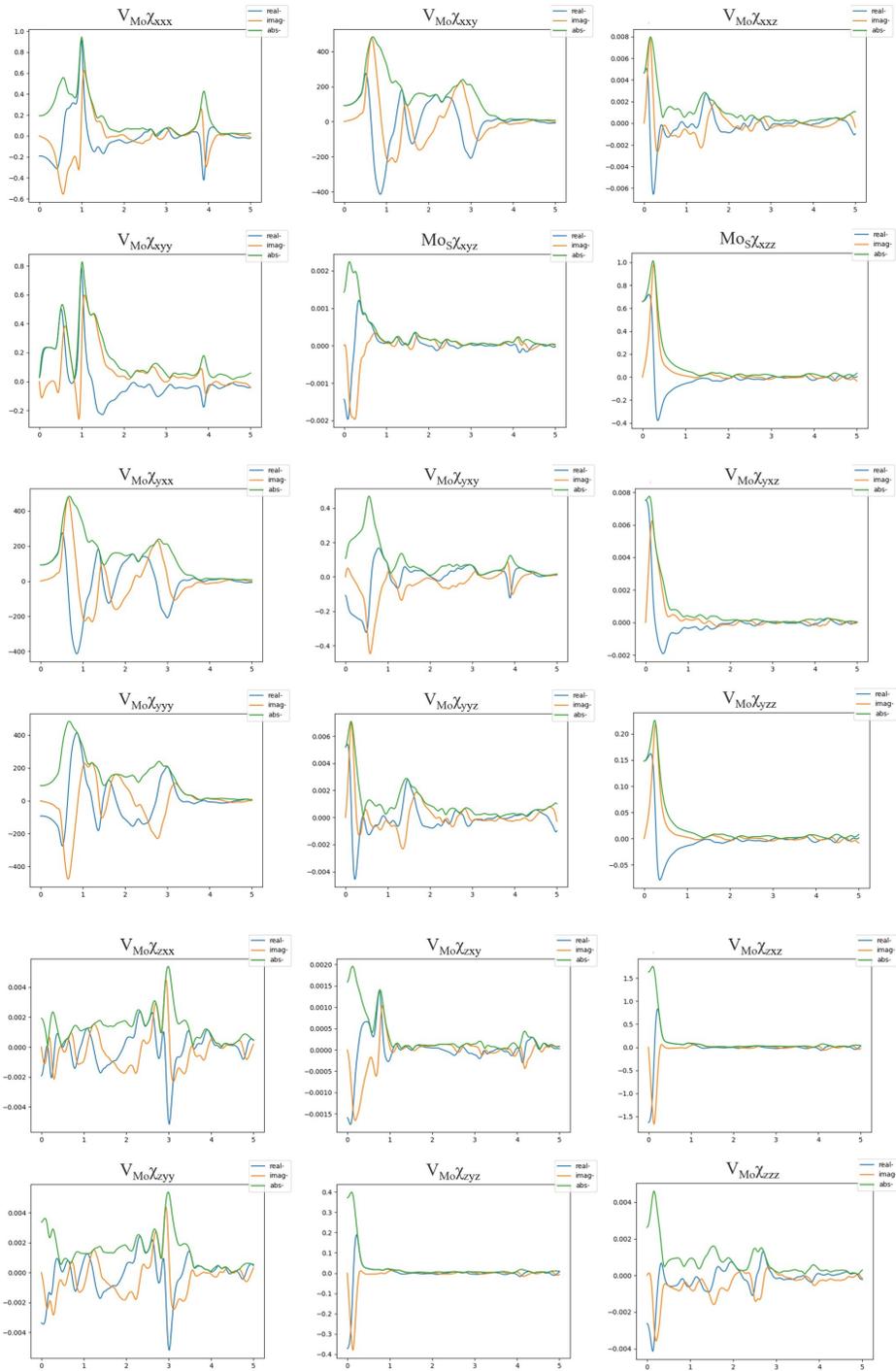


Figure 11: All SHG components with V_{Mo} with defects in 4X4 supercell