

Supplementary Information to: Intrinsic Point Defects and Frenkel Pair Formation in Photovoltaic Absorber Zn_3P_2 : Regulating p -type Conductivity

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1 Pristine lattice constants

The structural properties of the pristine Zn_3P_2 unit cell were established by relaxing the lattice and the ionic positions using different exchange-correlation functionals and fitting the resulting energies to the Birch–Murnaghan equation of state [1, 2], see Fig. S11. The resulting ground state properties, highlighting the sensitivity to the choice of functional and motivating the selection of PBEsol for subsequent calculations, are summarized in Table S11. We chose to proceed with the PBEsol-optimized model. This choice stems from the well-documented tendency of LDA to overbind, leading to an underestimation of lattice parameters, and a similar tendency of PBE to overestimate cell dimensions.

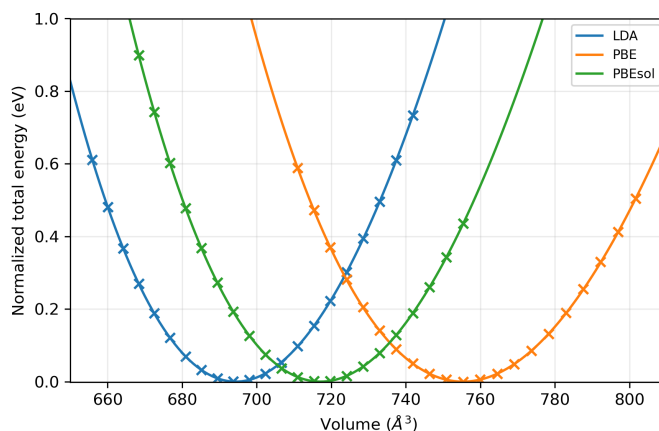


Figure S11 Total energy as a function of volume computed using LDA, PBE, and PBEsol. Markers denote calculated data points, while solid lines show fits to the Birch-Murnaghan equation of state. Energies are normalized such that the minimum of each fit is set to zero.

Table S11 Structural properties and Equation of State (EOS) parameters of the 40-atom Zn_3P_2 unit cell calculated using various exchange-correlation functionals. Lattice constants (a , c) and the c/a ratio are compared with available literature values. Reference data are followed by the respective functional used in the cited study in brackets. The cohesive energy E_c is referenced here to spin-polarised free atoms (including atomic spin-polarisation corrections).

| Method | a (Å) | c (Å) | c/a | B_0 (GPa) | B'_0 | E_c/atom (eV) |
|---------------|---------|---------|-----------------|-------------|--------|------------------------|
| LDA | 7.898 | 11.130 | 1.409 | 82.2 | 4.80 | 3.096 |
| PBEsol | 7.982 | 11.255 | 1.410 | 76.2 | 5.21 | 2.671 |
| PBE | 8.119 | 11.457 | 1.411 | 64.2 | 5.06 | 2.255 |
| Lit. (Theory) | 8.029 | 11.336 | 1.412 (PBE) [3] | | | |
| | 8.108 | 11.416 | 1.408 (PBE) [4] | | | |
| | 8.076 | 11.388 | 1.410 (HSE) [5] | | | |
| | 8.080 | 11.380 | 1.408 (HSE) [6] | | | |
| | 8.160 | 11.342 | 1.390 (HSE) [4] | | | |
| Lit. (Expt.) | 8.097 | 11.450 | 1.414 [7] | | | |
| | 8.0785 | 11.3966 | 1.4107 [8] | | | |
| | 8.02 | 11.67 | 1.46 [9] | | | |

2 Regarding the errors reported for the effective masses

In principle, effective masses can be obtained by fitting a parabolic function to the electronic band dispersion in the vicinity of a band extremum at a high-symmetry point, along a given crystallographic direction. The effective mass m^* is then extracted from the curvature of the band according to

$$\frac{1}{m^*} = \frac{1}{\hbar^2} \frac{\partial^2 E(k)}{\partial k^2}, \tag{1}$$

where $E(k)$ is the band energy and \hbar is the reduced Planck's constant.

In practice, however, the numerical determination of effective masses involves several methodological choices that can significantly influence the resulting values, particularly for bands with small curvature (heavy electrons or holes). One key parameter is the size of the k -space window around the band extremum, which must be chosen small enough to avoid non-parabolic features of the dispersion, yet large enough to allow for a numerically stable fit. Additional ambiguity arises from the choice of fitting function. While a second-order polynomial is sufficient in theory, higher-order fits are sometimes employed to improve robustness with respect to finite k -point sampling. Furthermore, different fitting strategies may either allow all polynomial coefficients to vary freely or constrain lower-order terms, such as fixing the band-edge energy right at the high-symmetry point to its computed value.

To account for these sources of uncertainty, and given that we find no single fitting protocol can be regarded as uniquely optimal, we evaluated a set of eight different fitting schemes. Specifically, we performed fits using either the ten k -points closest to the band extremum (see Fig. 3 of the main work) or only the five closest points. For each of these two data sets, four fitting functions were considered: (i) a second-order polynomial with all coefficients treated as free parameters, (ii) a third-order polynomial with all coefficients free, (iii) a second-order polynomial in which only the quadratic coefficient was fitted while the two lower-order terms were fixed, and (iv) a third-order polynomial in which only the quadratic and cubic coefficients were fitted.

The effective masses reported in this work correspond to the average over all eight fitting schemes. The associated uncertainties are defined by the maximum deviation of the individual results from this average, i.e., the difference between the average value and the smallest or largest effective mass obtained across all fitting approaches. This procedure provides a conservative estimate of the uncertainty arising from the numerical extraction of effective masses.

3 Influence of exchange–correlation functional and spin–orbit coupling on effective masses

To assess the robustness of the effective masses reported in the main text, we evaluated the influence of the explicit inclusion of spin-orbit coupling (SOC), and the choice of exchange-correlation functional. For that purpose, we performed a complementary calculation in which SOC was neglected, as well as a calculation using the PBE functional.

The resulting band structures close to the band gap at Γ are shown in Fig. S12 and corresponding effective masses are listed in Table S12 for direct comparison. We observe that spin-orbit coupling has a significant impact on the electronic structure near the band extrema, leading to band splitting and modified curvatures that directly influence the effective

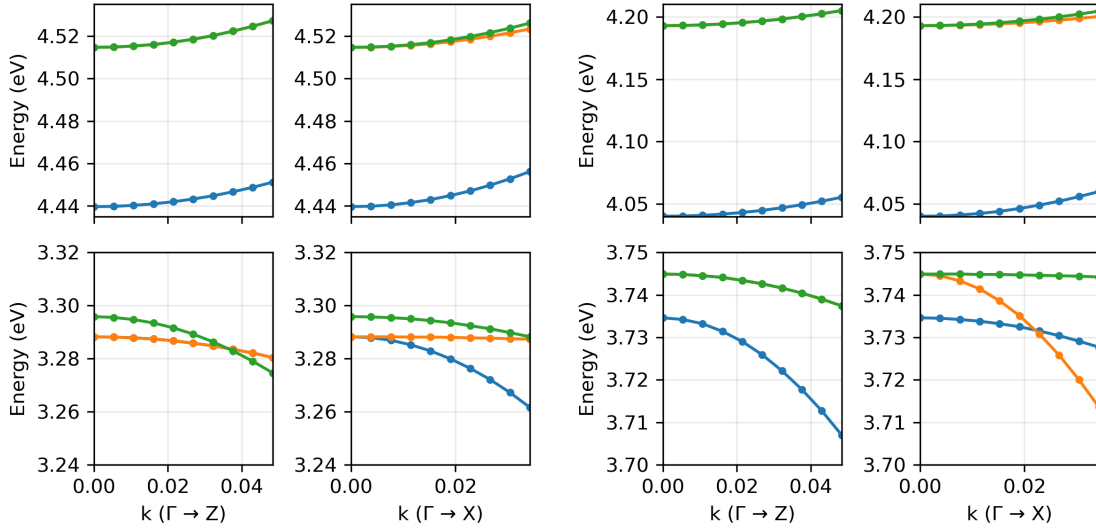


Figure S12 Electronic band structures near the valence- and conduction-band edges around the Γ point. Left: HSE06 calculations without spin-orbit coupling. Right: PBE calculations without spin-orbit coupling. Effective masses reported in Table S12 were extracted by fitting the band curvature in the immediate vicinity of Γ .

masses. Neglecting SOC therefore introduces additional quantitative errors. Furthermore, we find that the use of the PBE functional leads to noticeable deviations in the band curvatures compared to HSE06. These differences are accompanied by the well-known underestimation of the band gap within PBE.

Based on these observations, we conclude that the effective masses reported in the main text, obtained using the HSE06 functional with explicit inclusion of spin-orbit coupling, provide the most reliable description within the present computational framework. The additional results presented here serve to illustrate the sensitivity of effective masses to methodological choices and to justify the approach adopted in the main work.

Table S12 Effective carrier masses at the Γ point computed using different exchange-correlation functionals (HSE06 and PBE) and with or without spin-orbit coupling (wSOC, woSOC). For reference, the values reported in the main text correspond to the HSE06 calculations including spin-orbit coupling (HSE06, wSOC).

| path | band | m^* (HSE06, wSOC) | m^* (HSE06, woSOC) | m^* (PBE, woSOC) |
|------------------------|------------------------|---------------------|----------------------|--------------------|
| $\Gamma \rightarrow X$ | VB3 | -0.26 ± 0.09 | -0.10 ± 0.01 | -0.44 ± 0.10 |
| | VB2 | -0.24 ± 0.10 | < -2.4 | -0.09 ± 0.01 |
| | VB1 | -0.50 ± 0.20 | -0.38 ± 0.03 | < -3.8 |
| | CB1 | 0.17 ± 0.01 | 0.17 ± 0.01 | 0.14 ± 0.01 |
| | CB2 | 0.28 ± 0.02 | 0.33 ± 0.06 | 0.46 ± 0.10 |
| | CB3 | 0.27 ± 0.02 | 0.24 ± 0.04 | 0.23 ± 0.01 |
| | $\Gamma \rightarrow Z$ | VB3 | -0.31 ± 0.09 | -0.47 ± 0.13 |
| VB2 | | -0.35 ± 0.01 | -0.47 ± 0.13 | -0.39 ± 0.02 |
| VB1 | | -0.16 ± 0.02 | -0.14 ± 0.01 | -0.39 ± 0.02 |
| CB1 | | 0.31 ± 0.09 | 0.23 ± 0.02 | 0.18 ± 0.01 |
| CB2 | | 0.27 ± 0.05 | 0.23 ± 0.02 | 0.22 ± 0.02 |
| CB3 | | 0.27 ± 0.05 | 0.23 ± 0.02 | 0.22 ± 0.02 |

4 Formation energies of various defects computed in this work

This section provides the comprehensive numerical data for the intrinsic defect calculations presented in the main text. Tabulated values include raw energy differences (ΔE), finite-size supercell corrections (E_{corr}), and final formation energies (E^{form}) for all investigated charge states. To facilitate the exact reconstruction of the transition energy diagrams, these data should be used in combination with the chemical potentials defined in the manuscript and the computed valence band maximum (VBM) of 3.289 eV.

As shown, certain defect-charge state combinations exhibit an exceedingly high Kumagai-Oba correction uncertainty (ΔE_{corr}) beyond 50 meV. These elevated uncertainties are primarily associated with unstable or shallow defect states, which are intrinsically difficult to model within the periodic supercell approach. In this study, such instances correspond to highly charged species (e.g., P_i in charge state ± 5) that possess significantly higher formation energies than the dominant defect species. Given their negligible thermodynamic concentrations and the convergence limitations of the supercell model for these states, these entries were excluded from further refinement and the resulting transition level analysis.

Table S13 Detailed breakdown of the energy components used to calculate the formation energies (E^{form}) of intrinsic defects in Zn_3P_2 . All values are reported in eV. Following the doped convention, inequivalent defect sites are uniquely identified by their point group symmetry and nearest-neighbor distance. The here reported values for E^{form} are calculated at the P-rich stability limit with the Fermi level set to 0.12 eV above the valence band maximum (VBM). ΔE is the raw DFT energy difference between the defect and host supercells. E_{corr} is the finite-size supercell charge correction (Kumagai-Oba). ΔE_{corr} represents the estimated error in the charge correction based on the potential sampling variance.

| Defect | Charge State | ΔE | E_{corr} | ΔE_{corr} | E^{form} |
|--------------------------|--------------|------------|-------------------|--------------------------|-------------------|
| v_Zn_Cs_P2.32 | +1 | 0.275 | -0.135 | 0.011 | 2.098 |
| v_Zn_Cs_P2.32 | ± 0 | 3.229 | 0.000 | 0.000 | 1.779 |
| v_Zn_Cs_P2.32 | -1 | 6.308 | 0.156 | 0.005 | 1.605 |
| v_Zn_Cs_P2.32 | -2 | 9.393 | 0.327 | 0.012 | 1.452 |
| v_Zn_Cs_P2.35P2.38P2.70a | +1 | 0.264 | -0.133 | 0.010 | 2.089 |
| v_Zn_Cs_P2.35P2.38P2.70a | ± 0 | 3.162 | 0.000 | 0.000 | 1.711 |
| v_Zn_Cs_P2.35P2.38P2.70a | -1 | 6.201 | 0.129 | 0.007 | 1.471 |
| v_Zn_Cs_P2.35P2.38P2.70a | -2 | 9.275 | 0.252 | 0.017 | 1.259 |
| v_Zn_Cs_P2.35P2.38P2.70b | +1 | 0.564 | -0.153 | 0.010 | 2.369 |
| v_Zn_Cs_P2.35P2.38P2.70b | ± 0 | 3.470 | 0.000 | 0.000 | 2.019 |
| v_Zn_Cs_P2.35P2.38P2.70b | -1 | 6.504 | 0.157 | 0.007 | 1.802 |
| v_Zn_Cs_P2.35P2.38P2.70b | -2 | 9.554 | 0.314 | 0.015 | 1.599 |
| v_P_C2_Zn2.38 | +3 | -0.823 | -0.325 | 0.032 | 2.237 |
| v_P_C2_Zn2.38 | +2 | 2.074 | -0.228 | 0.015 | 1.823 |
| v_P_C2_Zn2.38 | +1 | 4.966 | -0.117 | 0.005 | 1.416 |

| | | | | | |
|------------------------------|----|---------|--------|-------|-------|
| v_P_C2_Zn2.38 | ±0 | 8.951 | 0.000 | 0.000 | 2.110 |
| v_P_C2_Zn2.38 | -1 | 12.654 | 0.153 | 0.013 | 2.557 |
| v_P_C2v_Zn2.32Zn2.35Zn2.70a | +3 | -1.017 | -0.355 | 0.032 | 2.013 |
| v_P_C2v_Zn2.32Zn2.35Zn2.70a | +2 | 1.853 | -0.229 | 0.018 | 1.601 |
| v_P_C2v_Zn2.32Zn2.35Zn2.70a | +1 | 4.750 | -0.113 | 0.008 | 1.204 |
| v_P_C2v_Zn2.32Zn2.35Zn2.70a | ±0 | 8.653 | 0.000 | 0.000 | 1.812 |
| v_P_C2v_Zn2.32Zn2.35Zn2.70a | -1 | 12.151 | 0.156 | 0.013 | 2.057 |
| v_P_C2v_Zn2.32Zn2.35Zn2.70b | +3 | -0.928 | -0.328 | 0.032 | 2.129 |
| v_P_C2v_Zn2.32Zn2.35Zn2.70b | +2 | 1.933 | -0.215 | 0.017 | 1.695 |
| v_P_C2v_Zn2.32Zn2.35Zn2.70b | +1 | 4.832 | -0.111 | 0.007 | 1.288 |
| v_P_C2v_Zn2.32Zn2.35Zn2.70b | ±0 | 8.656 | 0.000 | 0.000 | 1.815 |
| v_P_C2v_Zn2.32Zn2.35Zn2.70b | -1 | 12.064 | 0.145 | 0.013 | 1.959 |
| Zn_P_C2v_Zn2.32Zn2.35Zn2.70a | +5 | -8.599 | 0.026 | 0.072 | 3.080 |
| Zn_P_C2v_Zn2.32Zn2.35Zn2.70a | +4 | -5.687 | 0.040 | 0.046 | 2.598 |
| Zn_P_C2v_Zn2.32Zn2.35Zn2.70a | +3 | -2.723 | 0.031 | 0.028 | 2.144 |
| Zn_P_C2v_Zn2.32Zn2.35Zn2.70a | +2 | 0.882 | -0.022 | 0.016 | 2.287 |
| Zn_P_C2v_Zn2.32Zn2.35Zn2.70a | +1 | 4.313 | -0.032 | 0.007 | 2.299 |
| Zn_P_C2v_Zn2.32Zn2.35Zn2.70a | ±0 | 8.112 | 0.000 | 0.000 | 2.721 |
| Zn_P_C2v_Zn2.32Zn2.35Zn2.70b | +5 | -8.440 | 0.015 | 0.065 | 3.229 |
| Zn_P_C2v_Zn2.32Zn2.35Zn2.70b | +4 | -5.524 | 0.021 | 0.041 | 2.741 |
| Zn_P_C2v_Zn2.32Zn2.35Zn2.70b | +3 | -2.562 | -0.005 | 0.024 | 2.268 |
| Zn_P_C2v_Zn2.32Zn2.35Zn2.70b | +2 | 1.008 | -0.051 | 0.015 | 2.384 |
| Zn_P_C2v_Zn2.32Zn2.35Zn2.70b | +1 | 4.438 | -0.049 | 0.008 | 2.406 |
| Zn_P_C2v_Zn2.32Zn2.35Zn2.70b | ±0 | 8.331 | 0.000 | 0.000 | 2.940 |
| Zn_P_C2_Zn2.38 | +5 | -8.789 | 0.094 | 0.072 | 2.959 |
| Zn_P_C2_Zn2.38 | +4 | -5.880 | 0.098 | 0.044 | 2.463 |
| Zn_P_C2_Zn2.38 | +3 | -2.937 | 0.087 | 0.025 | 1.986 |
| Zn_P_C2_Zn2.38 | +2 | 0.623 | -0.013 | 0.013 | 2.037 |
| Zn_P_C2_Zn2.38 | +1 | 4.216 | -0.045 | 0.007 | 2.189 |
| Zn_P_C2_Zn2.38 | ±0 | 8.071 | 0.000 | 0.000 | 2.680 |
| P_Zn_Cs_P2.32 | +3 | -14.493 | 0.535 | 0.043 | 1.660 |
| P_Zn_Cs_P2.32 | +2 | -10.463 | 0.237 | 0.030 | 1.983 |
| P_Zn_Cs_P2.32 | +1 | -7.348 | 0.121 | 0.014 | 1.573 |
| P_Zn_Cs_P2.32 | ±0 | -2.503 | 0.000 | 0.000 | 2.888 |
| P_Zn_Cs_P2.32 | -1 | 0.971 | -0.020 | 0.011 | 2.933 |
| P_Zn_Cs_P2.32 | -2 | 5.539 | 0.035 | 0.021 | 4.147 |
| P_Zn_Cs_P2.32 | -3 | 9.777 | 0.206 | 0.044 | 5.147 |
| P_Zn_Cs_P2.32 | -4 | 14.413 | 0.302 | 0.070 | 6.470 |
| P_Zn_Cs_P2.32 | -5 | 19.145 | 0.320 | 0.101 | 7.811 |
| P_Zn_Cs_P2.35P2.38P2.70a | +3 | -14.593 | 0.450 | 0.031 | 1.474 |
| P_Zn_Cs_P2.35P2.38P2.70a | +2 | -10.547 | 0.245 | 0.019 | 1.907 |
| P_Zn_Cs_P2.35P2.38P2.70a | +1 | -7.134 | 0.079 | 0.012 | 1.744 |
| P_Zn_Cs_P2.35P2.38P2.70a | ±0 | -2.910 | 0.000 | 0.000 | 2.481 |
| P_Zn_Cs_P2.35P2.38P2.70a | -1 | 0.688 | -0.034 | 0.013 | 2.636 |
| P_Zn_Cs_P2.35P2.38P2.70a | -2 | 5.365 | -0.003 | 0.022 | 3.935 |
| P_Zn_Cs_P2.35P2.38P2.70a | -3 | 9.245 | 0.135 | 0.052 | 4.543 |
| P_Zn_Cs_P2.35P2.38P2.70a | -4 | 13.894 | 0.357 | 0.071 | 6.006 |
| P_Zn_Cs_P2.35P2.38P2.70a | -5 | 18.563 | 0.631 | 0.087 | 7.540 |
| P_Zn_Cs_P2.35P2.38P2.70b | +3 | -14.707 | 0.527 | 0.040 | 1.437 |
| P_Zn_Cs_P2.35P2.38P2.70b | +2 | -10.245 | 0.225 | 0.017 | 2.189 |
| P_Zn_Cs_P2.35P2.38P2.70b | +1 | -7.077 | 0.102 | 0.007 | 1.824 |
| P_Zn_Cs_P2.35P2.38P2.70b | ±0 | -2.456 | 0.000 | 0.000 | 2.935 |
| P_Zn_Cs_P2.35P2.38P2.70b | -1 | 1.069 | -0.013 | 0.014 | 3.038 |
| P_Zn_Cs_P2.35P2.38P2.70b | -2 | 5.544 | 0.031 | 0.023 | 4.148 |
| P_Zn_Cs_P2.35P2.38P2.70b | -3 | 9.241 | 0.217 | 0.035 | 4.623 |
| P_Zn_Cs_P2.35P2.38P2.70b | -4 | 14.128 | 0.295 | 0.098 | 6.179 |
| P_Zn_Cs_P2.35P2.38P2.70b | -5 | 17.984 | 0.489 | 0.143 | 6.820 |

| | | | | | |
|-----------------|----|---------|--------|-------|-------|
| Zn_i_C1 | +2 | -8.107 | 0.283 | 0.013 | 0.444 |
| Zn_i_C1 | +1 | -3.393 | 0.120 | 0.005 | 1.587 |
| Zn_i_C1 | ±0 | 1.327 | 0.000 | 0.000 | 2.777 |
| Zn_i_C2v | +2 | -8.097 | 0.282 | 0.013 | 0.454 |
| Zn_i_C2v | +1 | -3.401 | 0.121 | 0.005 | 1.580 |
| Zn_i_C2v | ±0 | 1.889 | 0.000 | 0.000 | 3.339 |
| Zn_i_Ci | +2 | -7.377 | 0.288 | 0.015 | 1.179 |
| Zn_i_Ci | +1 | -2.812 | 0.104 | 0.005 | 2.151 |
| Zn_i_Ci | ±0 | 1.821 | 0.000 | 0.000 | 3.272 |
| Zn_i_Cs_P2.45 | +2 | -8.098 | 0.283 | 0.013 | 0.453 |
| Zn_i_Cs_P2.45 | +1 | -3.400 | 0.121 | 0.005 | 1.580 |
| Zn_i_Cs_P2.45 | ±0 | 1.327 | 0.000 | 0.000 | 2.778 |
| Zn_i_Cs_Zn2.30 | +2 | -8.098 | 0.282 | 0.013 | 0.453 |
| Zn_i_Cs_Zn2.30 | +1 | -3.403 | 0.121 | 0.005 | 1.577 |
| Zn_i_Cs_Zn2.30 | ±0 | 1.325 | 0.000 | 0.000 | 2.776 |
| Zn_i_D2d_Zn2.16 | +2 | -8.097 | 0.282 | 0.013 | 0.454 |
| Zn_i_D2d_Zn2.16 | +1 | -3.397 | 0.120 | 0.005 | 1.583 |
| Zn_i_D2d_Zn2.16 | ±0 | 1.328 | 0.000 | 0.000 | 2.779 |
| Zn_i_D2d_Zn2.56 | +2 | -8.123 | 0.285 | 0.013 | 0.431 |
| Zn_i_D2d_Zn2.56 | +1 | -2.473 | 0.118 | 0.005 | 2.505 |
| Zn_i_D2d_Zn2.56 | ±0 | 1.831 | 0.000 | 0.000 | 3.281 |
| P_i_C1 | +5 | -21.096 | 1.098 | 0.085 | 3.888 |
| P_i_C1 | +4 | -17.293 | 0.874 | 0.062 | 4.058 |
| P_i_C1 | +3 | -14.056 | 0.673 | 0.045 | 3.684 |
| P_i_C1 | +2 | -10.524 | 0.302 | 0.015 | 3.437 |
| P_i_C1 | +1 | -7.300 | 0.155 | 0.007 | 3.105 |
| P_i_C1 | ±0 | -3.551 | 0.000 | 0.000 | 3.291 |
| P_i_C1 | -1 | 0.165 | -0.075 | 0.010 | 3.523 |
| P_i_C1 | -2 | 5.081 | -0.121 | 0.017 | 4.983 |
| P_i_C1 | -3 | 9.872 | -0.132 | 0.036 | 6.354 |
| P_i_C2v | +5 | -21.199 | 1.074 | 0.079 | 3.761 |
| P_i_C2v | +4 | -16.779 | 0.702 | 0.045 | 4.399 |
| P_i_C2v | +3 | -13.490 | 0.464 | 0.028 | 4.042 |
| P_i_C2v | +2 | -10.529 | 0.304 | 0.016 | 3.433 |
| P_i_C2v | +1 | -7.311 | 0.157 | 0.007 | 3.097 |
| P_i_C2v | ±0 | -3.656 | 0.000 | 0.000 | 3.185 |
| P_i_C2v | -1 | 0.076 | -0.078 | 0.013 | 3.431 |
| P_i_C2v | -2 | 4.973 | -0.126 | 0.025 | 4.871 |
| P_i_C2v | -3 | 9.869 | -0.136 | 0.037 | 6.348 |
| P_i_Ci | +5 | -21.096 | 1.097 | 0.085 | 3.887 |
| P_i_Ci | +4 | -17.064 | 0.756 | 0.065 | 4.170 |
| P_i_Ci | +3 | -13.916 | 0.520 | 0.046 | 3.672 |
| P_i_Ci | +2 | -10.737 | 0.355 | 0.028 | 3.277 |
| P_i_Ci | +1 | -7.135 | 0.156 | 0.006 | 3.271 |
| P_i_Ci | ±0 | -3.514 | 0.000 | 0.000 | 3.327 |
| P_i_Ci | -1 | 0.222 | -0.059 | 0.010 | 3.596 |
| P_i_Ci | -2 | 5.181 | -0.085 | 0.021 | 5.119 |
| P_i_Ci | -3 | 10.147 | -0.076 | 0.037 | 6.686 |
| P_i_Cs_P2.45 | +5 | -20.018 | 0.879 | 0.097 | 4.747 |
| P_i_Cs_P2.45 | +4 | -18.165 | 0.872 | 0.085 | 3.183 |
| P_i_Cs_P2.45 | +3 | -13.727 | 0.422 | 0.047 | 3.763 |
| P_i_Cs_P2.45 | +2 | -10.549 | 0.276 | 0.030 | 3.386 |
| P_i_Cs_P2.45 | +1 | -7.308 | 0.158 | 0.007 | 3.100 |
| P_i_Cs_P2.45 | ±0 | -3.656 | 0.000 | 0.000 | 3.186 |
| P_i_Cs_P2.45 | -1 | 0.075 | -0.077 | 0.012 | 3.431 |
| P_i_Cs_P2.45 | -2 | 4.972 | -0.127 | 0.025 | 4.868 |
| P_i_Cs_P2.45 | -3 | 9.870 | -0.136 | 0.037 | 6.348 |

| | | | | | |
|----------------|---------|---------|--------|-------|-------|
| P_i_Cs_Zn2.30 | +5 | -21.286 | 1.085 | 0.118 | 3.684 |
| P_i_Cs_Zn2.30 | +4 | -18.166 | 0.887 | 0.085 | 3.199 |
| P_i_Cs_Zn2.30 | +3 | -13.723 | 0.441 | 0.030 | 3.786 |
| P_i_Cs_Zn2.30 | +2 | -10.527 | 0.303 | 0.015 | 3.435 |
| P_i_Cs_Zn2.30 | +1 | -7.299 | 0.156 | 0.007 | 3.108 |
| P_i_Cs_Zn2.30 | ± 0 | -3.661 | 0.000 | 0.000 | 3.180 |
| P_i_Cs_Zn2.30 | -1 | 0.166 | -0.075 | 0.010 | 3.523 |
| P_i_Cs_Zn2.30 | -2 | 4.968 | -0.122 | 0.024 | 4.870 |
| P_i_Cs_Zn2.30 | -3 | 9.873 | -0.132 | 0.036 | 6.355 |
| P_i_D2d_Zn2.16 | +5 | -21.283 | 1.086 | 0.118 | 3.689 |
| P_i_D2d_Zn2.16 | +4 | -18.163 | 0.887 | 0.085 | 3.202 |
| P_i_D2d_Zn2.16 | +3 | -15.018 | 0.681 | 0.060 | 2.730 |
| P_i_D2d_Zn2.16 | +2 | -10.920 | 0.391 | 0.027 | 3.130 |
| P_i_D2d_Zn2.16 | +1 | -7.302 | 0.156 | 0.007 | 3.105 |
| P_i_D2d_Zn2.16 | ± 0 | -3.652 | 0.000 | 0.000 | 3.190 |
| P_i_D2d_Zn2.16 | -1 | 0.077 | -0.077 | 0.013 | 3.432 |
| P_i_D2d_Zn2.16 | -2 | 4.972 | -0.126 | 0.025 | 4.869 |
| P_i_D2d_Zn2.16 | -3 | 9.869 | -0.136 | 0.036 | 6.347 |
| P_i_D2d_Zn2.56 | +5 | -21.288 | 1.083 | 0.118 | 3.681 |
| P_i_D2d_Zn2.56 | +4 | -18.165 | 0.887 | 0.085 | 3.199 |
| P_i_D2d_Zn2.56 | +3 | -15.017 | 0.671 | 0.059 | 2.722 |
| P_i_D2d_Zn2.56 | +2 | -10.528 | 0.304 | 0.016 | 3.435 |
| P_i_D2d_Zn2.56 | +1 | -7.305 | 0.157 | 0.007 | 3.103 |
| P_i_D2d_Zn2.56 | ± 0 | -3.655 | 0.000 | 0.000 | 3.187 |
| P_i_D2d_Zn2.56 | -1 | 0.075 | -0.077 | 0.012 | 3.431 |
| P_i_D2d_Zn2.56 | -2 | 4.970 | -0.126 | 0.024 | 4.867 |
| P_i_D2d_Zn2.56 | -3 | 9.870 | -0.135 | 0.037 | 6.349 |

5 Defect Transition Levels and Thermodynamic Concentrations

This section provides a detailed tabulation of the thermodynamic transition levels and the corresponding defect concentrations for the intrinsic species in Zn_3P_2 discussed in the main text. The transition levels, denoted as $\varepsilon(q/q')$, represent the Fermi level positions where the formation energies of two charge states q and q' are equal. The concentrations reported in Table SI4 were calculated using the determined formation energies under P-rich growth conditions, assuming an annealing temperature of $T = 650\text{K}$ and a quenching/operating temperature of $T = 300\text{K}$. Under these specific conditions, the self-consistent Fermi level was determined to be 0.12eV above the valence band maximum (VBM).

Table SI4 Calculated thermodynamic transition levels $\varepsilon(q/q')$ and associated total defect concentrations for intrinsic defects in Zn_3P_2 . Transition energies are reported in eV relative to the valence band maximum. Concentration were computed for P-rich growth conditions and a Fermi-level of 0.12eV above the valence band maximum.

| Defect | Charge Transition | Transition Energy (eV) | Total Concentration (cm^{-3}) |
|--------------------------|--------------------------|------------------------|--|
| v_Zn_Cs_P2.35P2.38P2.70b | $\varepsilon(+1, \pm 0)$ | -0.23 | $3.29\text{e}+13$ |
| v_Zn_Cs_P2.35P2.38P2.70b | $\varepsilon(\pm 0, -1)$ | -0.10 | $3.29\text{e}+13$ |
| v_Zn_Cs_P2.35P2.38P2.70b | $\varepsilon(-1, -2)$ | -0.08 | $3.29\text{e}+13$ |
| v_Zn_Cs_P2.32 | $\varepsilon(+1, \pm 0)$ | -0.20 | $9.10\text{e}+14$ |
| v_Zn_Cs_P2.32 | $\varepsilon(\pm 0, -1)$ | -0.05 | $9.10\text{e}+14$ |
| v_Zn_Cs_P2.32 | $\varepsilon(-1, -2)$ | -0.03 | $9.10\text{e}+14$ |
| v_Zn_Cs_P2.35P2.38P2.70a | $\varepsilon(+1, \pm 0)$ | -0.26 | $2.89\text{e}+16$ |
| v_Zn_Cs_P2.35P2.38P2.70a | $\varepsilon(\pm 0, -1)$ | -0.12 | $2.89\text{e}+16$ |
| v_Zn_Cs_P2.35P2.38P2.70a | $\varepsilon(-1, -2)$ | -0.09 | $2.89\text{e}+16$ |
| Zn_i_D2d_Zn2.56 | $\varepsilon(+1, \pm 0)$ | 0.90 | $1.96\text{e}+01$ |
| Zn_i_Ci | $\varepsilon(+2, +1)$ | 1.09 | $2.16\text{e}+09$ |
| Zn_i_Ci | $\varepsilon(+1, \pm 0)$ | 1.24 | $2.16\text{e}+09$ |
| Zn_i_C1 | $\varepsilon(+2, +1)$ | 1.27 | $6.14\text{e}+15$ |
| Zn_i_C1 | $\varepsilon(+1, \pm 0)$ | 1.32 | $6.14\text{e}+15$ |

| | | | |
|------------------------------|-----------------------|-------|----------|
| v_P_C2_Zn2.38 | $\epsilon(+3, +2)$ | -0.29 | 2.70e+09 |
| v_P_C2_Zn2.38 | $\epsilon(+2, +1)$ | -0.29 | 2.70e+09 |
| v_P_C2_Zn2.38 | $\epsilon(+1, -1)$ | 0.69 | 2.70e+09 |
| v_P_C2v_Zn2.32Zn2.35Zn2.70b | $\epsilon(+3, +2)$ | -0.31 | 1.41e+10 |
| v_P_C2v_Zn2.32Zn2.35Zn2.70b | $\epsilon(+2, +1)$ | -0.29 | 1.41e+10 |
| v_P_C2v_Zn2.32Zn2.35Zn2.70b | $\epsilon(+1, -1)$ | 0.46 | 1.41e+10 |
| v_P_C2v_Zn2.32Zn2.35Zn2.70a | $\epsilon(+3, +2)$ | -0.29 | 6.02e+10 |
| v_P_C2v_Zn2.32Zn2.35Zn2.70a | $\epsilon(+2, +1)$ | -0.28 | 6.02e+10 |
| v_P_C2v_Zn2.32Zn2.35Zn2.70a | $\epsilon(+1, -1)$ | 0.55 | 6.02e+10 |
| P_Zn_Cs_P2.35P2.38P2.70b | $\epsilon(+3, +1)$ | 0.31 | 2.10e+06 |
| P_Zn_Cs_P2.35P2.38P2.70b | $\epsilon(+1, -1)$ | 0.73 | 2.10e+06 |
| P_Zn_Cs_P2.35P2.38P2.70b | $\epsilon(-1, -3)$ | 0.87 | 2.10e+06 |
| P_Zn_Cs_P2.35P2.38P2.70b | $\epsilon(-3, -5)$ | 1.26 | 2.10e+06 |
| P_Zn_Cs_P2.35P2.38P2.70a | $\epsilon(+3, +1)$ | 0.26 | 7.88e+06 |
| P_Zn_Cs_P2.35P2.38P2.70a | $\epsilon(+1, -1)$ | 0.57 | 7.88e+06 |
| P_Zn_Cs_P2.35P2.38P2.70a | $\epsilon(-1, -5)$ | 1.41 | 7.88e+06 |
| P_Zn_Cs_P2.32 | $\epsilon(+3, +1)$ | 0.08 | 1.64e+08 |
| P_Zn_Cs_P2.32 | $\epsilon(+1, -1)$ | 0.80 | 1.64e+08 |
| P_Zn_Cs_P2.32 | $\epsilon(-1, -3)$ | 1.23 | 1.64e+08 |
| P_Zn_Cs_P2.32 | $\epsilon(-3, -4)$ | 1.44 | 1.64e+08 |
| Zn_P_C2v_Zn2.32Zn2.35Zn2.70b | $\epsilon(+5, +4)$ | -0.37 | 3.01e+01 |
| Zn_P_C2v_Zn2.32Zn2.35Zn2.70b | $\epsilon(+4, +3)$ | -0.35 | 3.01e+01 |
| Zn_P_C2v_Zn2.32Zn2.35Zn2.70b | $\epsilon(+3, +1)$ | 0.19 | 3.01e+01 |
| Zn_P_C2v_Zn2.32Zn2.35Zn2.70b | $\epsilon(+1, \pm 0)$ | 0.65 | 3.01e+01 |
| Zn_P_C2v_Zn2.32Zn2.35Zn2.70a | $\epsilon(+5, +4)$ | -0.36 | 2.18e+02 |
| Zn_P_C2v_Zn2.32Zn2.35Zn2.70a | $\epsilon(+4, +3)$ | -0.33 | 2.18e+02 |
| Zn_P_C2v_Zn2.32Zn2.35Zn2.70a | $\epsilon(+3, +1)$ | 0.20 | 2.18e+02 |
| Zn_P_C2v_Zn2.32Zn2.35Zn2.70a | $\epsilon(+1, \pm 0)$ | 0.54 | 2.18e+02 |
| Zn_P_C2_Zn2.38 | $\epsilon(+5, +4)$ | -0.38 | 3.30e+03 |
| Zn_P_C2_Zn2.38 | $\epsilon(+4, +3)$ | -0.36 | 3.30e+03 |
| Zn_P_C2_Zn2.38 | $\epsilon(+3, +2)$ | 0.17 | 3.30e+03 |
| Zn_P_C2_Zn2.38 | $\epsilon(+2, +1)$ | 0.27 | 3.30e+03 |
| Zn_P_C2_Zn2.38 | $\epsilon(+1, \pm 0)$ | 0.61 | 3.30e+03 |
| P_i_D2d_Zn2.16 | $\epsilon(+5, +4)$ | -0.38 | 6.87e-02 |
| P_i_D2d_Zn2.16 | $\epsilon(+4, +3)$ | -0.34 | 6.87e-02 |
| P_i_D2d_Zn2.16 | $\epsilon(+3, \pm 0)$ | 0.27 | 6.87e-02 |
| P_i_D2d_Zn2.16 | $\epsilon(+0, -1)$ | 0.37 | 6.87e-02 |
| P_i_D2d_Zn2.16 | $\epsilon(-1, -2)$ | 1.56 | 6.87e-02 |
| P_i_D2d_Zn2.16 | $\epsilon(-2, -3)$ | 1.60 | 6.87e-02 |

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