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## Supporting information

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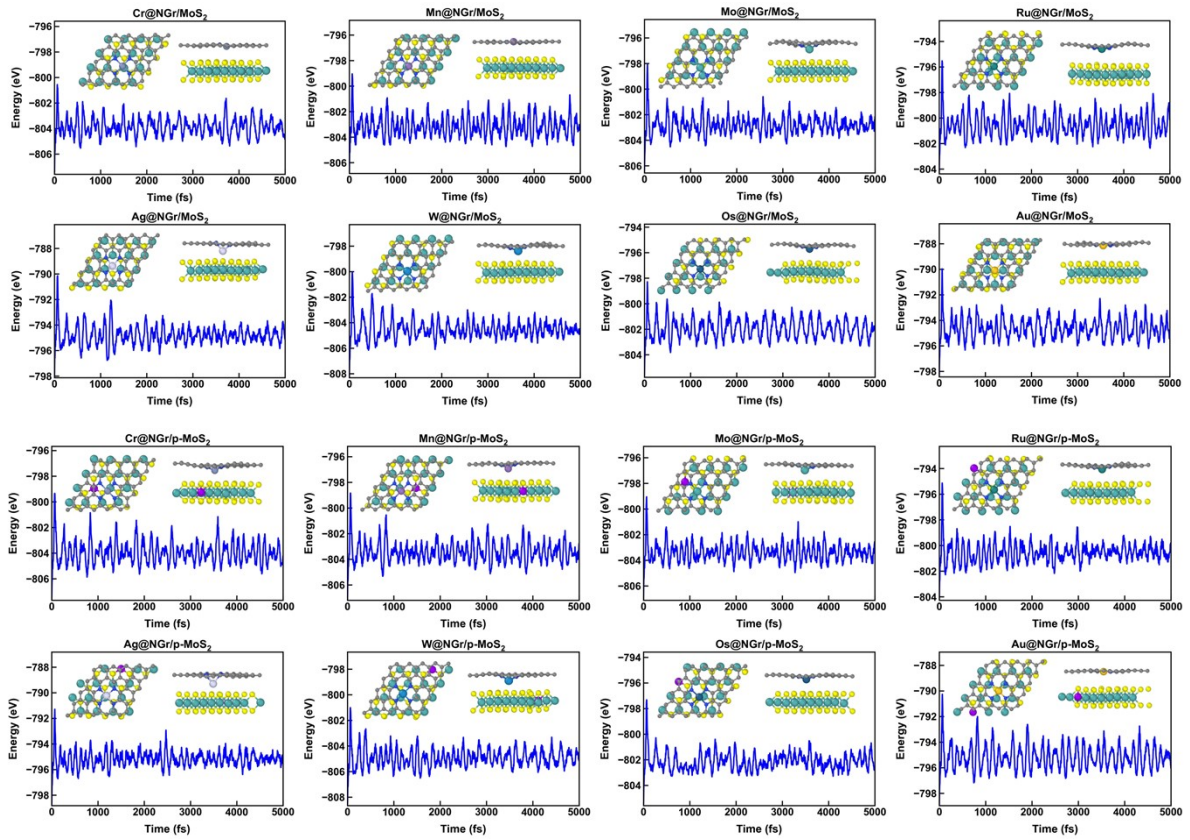
### 3 **Breaking the stability-activity trade-off in single-atom catalysis via interlayer electric** 4 **fields in van der Waals heterostructure**

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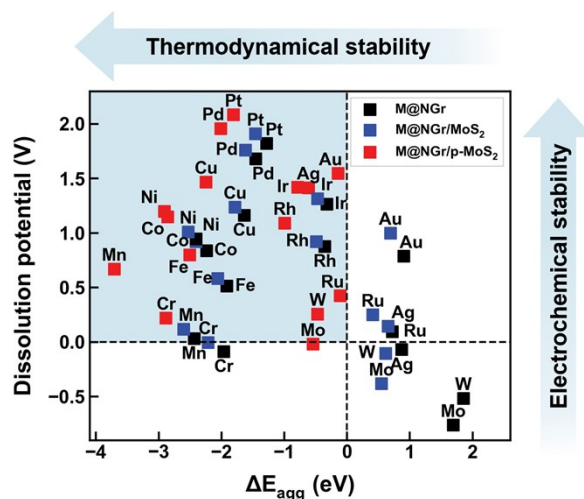
9 \*Corresponding author: [bchan@yonsei.ac.kr](mailto:bchan@yonsei.ac.kr)



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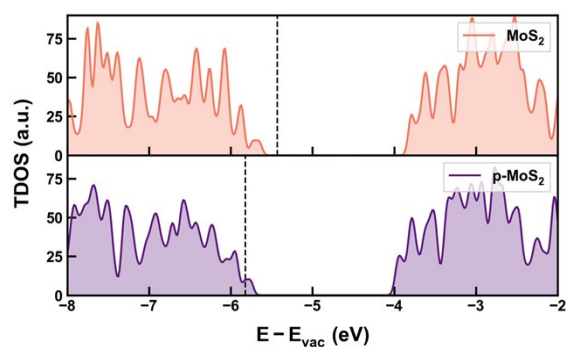
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3 **Figure S1.** The AIMD simulation results for  $M@NGr/MoS_2$  and  $M@NGr/p-MoS_2$  structures.  
 4 The inset displays the final structure after 5 ps. During the simulation,  $M@NGr$  and TMD  
 5 layers moved in  $x$  and  $y$  directions, which is a normal phenomenon considering that the two  
 6 layers interact with weak dispersion forces.



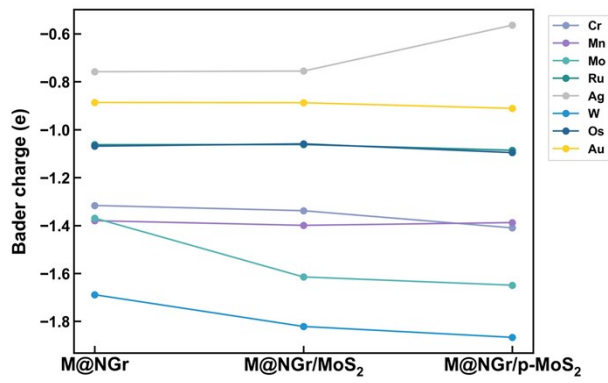
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2 **Figure S2.** Stability map of all transition metal single atom in M@NGr, M@NGr/MoS<sub>2</sub>, and  
 3 M@NGr/p-MoS<sub>2</sub> structures illustrating resistance against aggregation and electrochemical  
 4 dissolution. The blue quadrant denotes the optimal region where the single atom is both  
 5 stable against aggregation and dissolution during HER in acidic conditions.



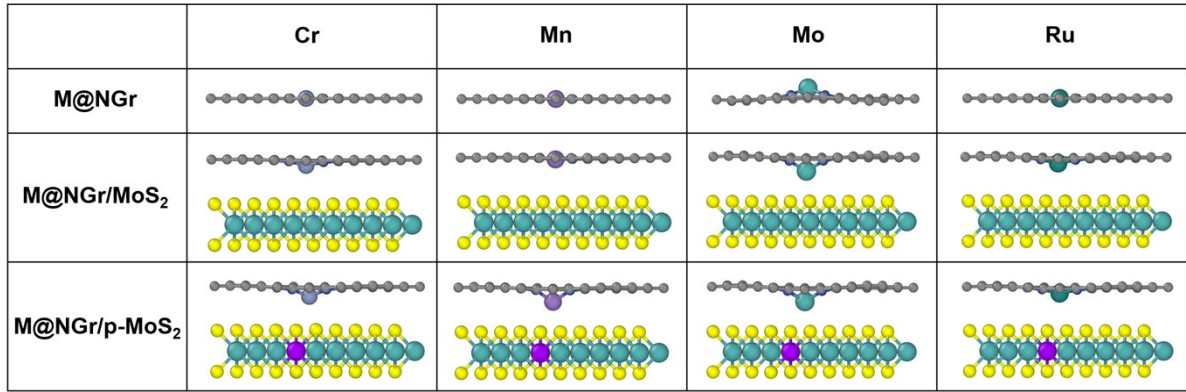
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2 **Figure S3.** Total density of states (TDOS) of isolated monolayers of MoS<sub>2</sub> and Nb-doped p-  
 3 type MoS<sub>2</sub> (p-MoS<sub>2</sub>). The energies are referenced to the vacuum level ( $E_{\text{vac}}$ ), and the dashed  
 4 lines indicate the respective Fermi level ( $E_{\text{F}}$ ) for each system. Nb doping shifts  $E_{\text{F}}$  to lower  
 5 energy, increasing the work function of p-MoS<sub>2</sub> relative to undoped MoS<sub>2</sub> (5.82 eV vs. 5.43  
 6 eV; see Table S5) and thereby amplifying the work function mismatch with M@NGr upon  
 7 heterostructure formation.

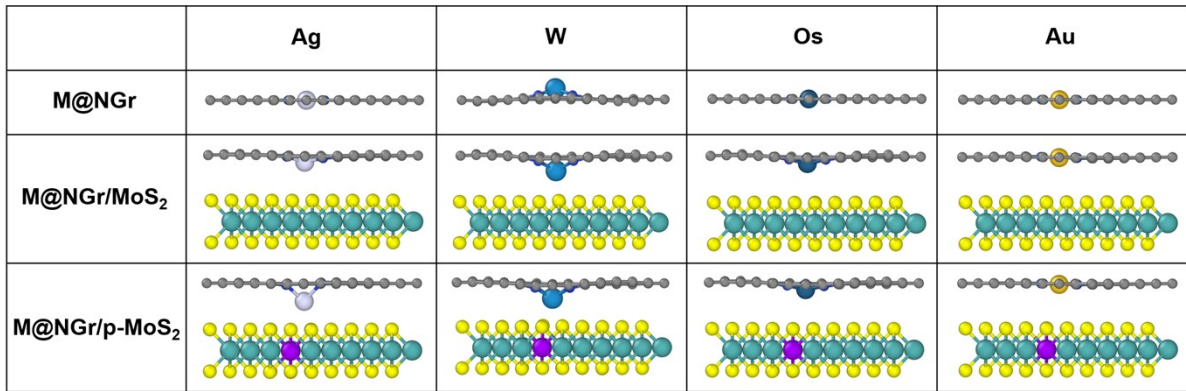


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2 **Figure S4.** Bader charge of metal single atom in M@NGr, M@NGr/MoS<sub>2</sub>, and M@NGr/p-  
 3 MoS<sub>2</sub>. The interfacial interactions do not significantly alter the charge state of metal atoms,  
 4 suggesting that the interlayer charge transfer mainly occurs from C/N atoms in M@NGr to  
 5 TMD.



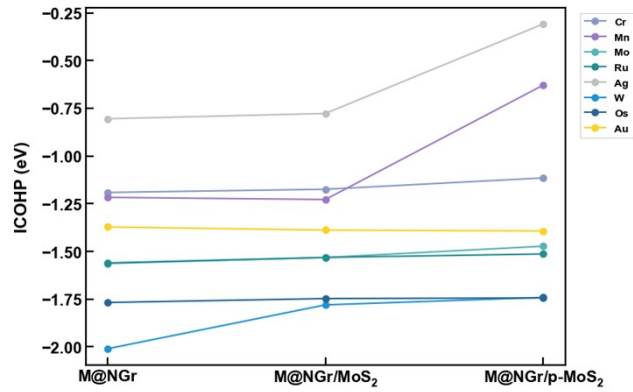
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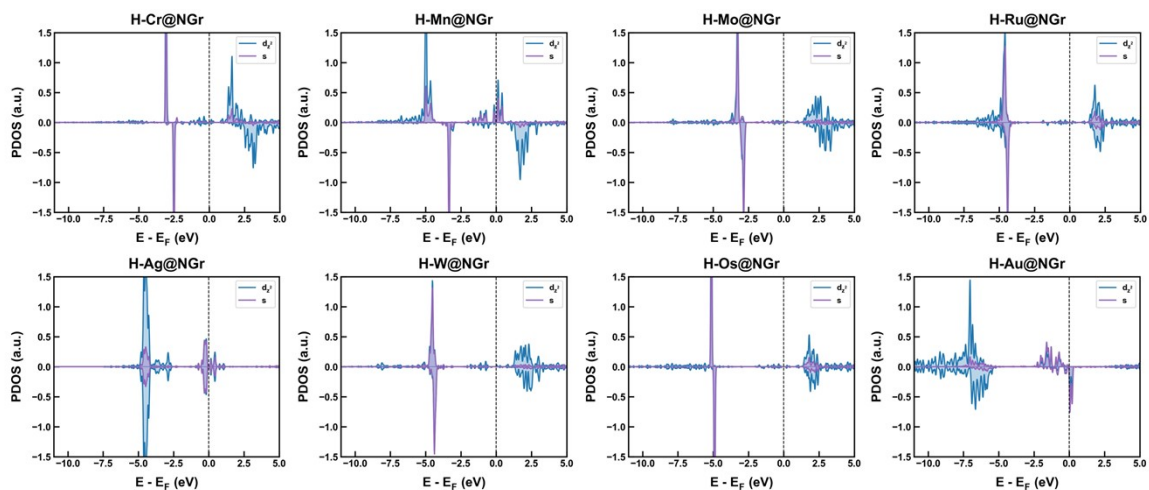
3 **Figure S5.** Side views of the DFT-optimized geometric structures of M@NGr, M@NGr/MoS<sub>2</sub>,  
 4 and M@NGr/p-MoS<sub>2</sub> for selected transition metals, showing the downward displacement ( $\Delta z$ )  
 5 of the single metal atom along out-of-plane direction upon heterostructure formation. The  
 6 progressive increase in  $\Delta z$  from M@NGr/MoS<sub>2</sub> to M@NGr/p-MoS<sub>2</sub> is consistent with the  
 7 stronger built-in electric field and greater electrostatic attraction toward the TMD layer in the  
 8 p-doped system. Color code: gray (C), blue (N), yellow (S), green (Mo), purple (Nb), and  
 9 element-specific color for the single metal atom.

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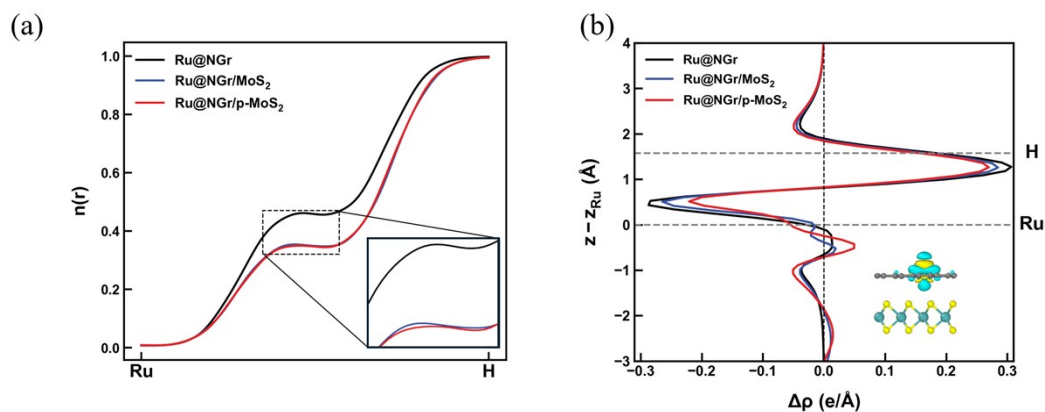
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2 **Figure S6.** Integrated Crystal Orbital Hamilton Population (ICOHP) values for metal-N bond  
 3 in M@NGr, M@NGr/MoS<sub>2</sub>, and M@NGr/p-MoS<sub>2</sub>, quantifying metal-N bond strength. The  
 4 weakened ICOHP values are consistent with increased M-N bond lengths. For most metals, the  
 5 M-N bond strengths remains nearly unchanged, showing that the enhanced single atom stability  
 6 is not due to stronger M-N interactions but because of electrostatic attraction from interlayer  
 7 built-in electric field.



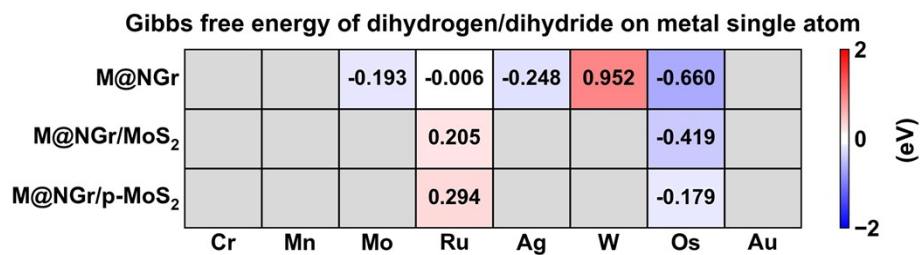
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2 **Figure S7.** Projected density of states (PDOS) of the metal  $d_{22}$  orbital and the adsorbed H 1s  
 3 orbital in H-adsorbed M@NGr structures for all metals investigated. The energy levels of  
 4 bonding and antibonding orbitals are consistent with the projected crystal orbital Hamiltonian  
 5 populations (pCOHP) in Figure 6c.



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2 **Figure S8.** Charge redistribution and bonding character analysis of the Ru-H interaction in H-  
 3 adsorbed Ru@NGr, Ru@NGr/MoS<sub>2</sub>, and Ru@NGr/p-MoS<sub>2</sub>. (a) Electron localization function  
 4 (ELF) profiles computed along the line passing through the Ru-H bond axis, demonstrating the  
 5 progressive reduction in covalency of Ru-H bonding upon heterostructure formation. (b) Plane-  
 6 averaged and three-dimensional charge density difference ( $\Delta\rho = \rho_{\text{total}} - \rho_{\text{slab}} - \rho_{\text{H}^*}$ ) for each  
 7 system (yellow: charge accumulation; cyan: charge depletion). The progressive attenuation of  
 8 charge accumulation along the Ru-H axis from Ru@NGr to Ru@NGr/p-MoS<sub>2</sub> is consistent  
 9 with the increasing Ru  $d_{22}$  orbital delocalization induced by interlayer  $d_{22}$ - $p_z$  hybridization and  
 10 weakening of the Ru-H covalent bond.



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3 **Figure S9.** Free energy of dihydrogen/dihydride on metal single atoms in M@NGr,  
4 M@NGr/MoS<sub>2</sub>, and M@NGr/p-MoS<sub>2</sub> systems. Shaded regions correspond to cases where H<sub>2</sub>  
5 was only weakly physisorbed with large metal-H distance.

1 **Table S1.** Optimized in-plane lattice parameters of graphene (Gr), and MoS<sub>2</sub> monolayers

System	Lattice constant (Å)	Note
Gr	2.467	GGA-PBE optimized
MoS <sub>2</sub> (unstrained)	3.162	GGA-PBE optimized
MoS <sub>2</sub> (strained -2.47%)	3.084	Applied to match Gr supercell

2 (unstrained and strained).

1 **Table S2.** Aggregation energy ( $\Delta E_{agg}$ ) and single atom dissolution potential ( $U_{diss,SAC}^{n+}$ ) values  
 2 for all transition metal SACs in the M@NGr reference system used for Figure 2a, Figure 2b

<b>M</b>	$E_{coh}$	$E_{emb}$	$\Delta E_{agg}$	$U_{diss,bulk}^{n+}$			$U_{diss,SAC}^{n+}$		
				<b>1+</b>	<b>2+</b>	<b>3+</b>	<b>1+</b>	<b>2+</b>	<b>3+</b>
<b>Cr</b>	-4.150	-6.113	-1.963		-0.913	-0.744		0.068	-0.090
<b>Mn</b>	-2.920	-5.352	-2.432		-1.185	-0.276		0.031	0.535
<b>Fe</b>	4.320	-6.235	-1.915		-0.447	-0.037		0.511	0.601
<b>Co</b>	4.470	-6.703	-2.233		-0.280	0.453		0.836	1.198
<b>Ni</b>	4.480	-6.885	-2.405		-0.257			0.945	
<b>Cu</b>	3.510	-5.145	-1.635	0.521	0.342	1.028	2.156	1.159	1.573
<b>Mo</b>	-6.840	-5.445	1.395			-0.200			-0.665
<b>Ru</b>	-6.800	-6.074	0.726		0.455	0.386		0.092	0.144
<b>Rh</b>	5.760	-6.110	-0.350	0.600		0.758	0.950		0.875
<b>Pd</b>	3.930	-5.383	-1.453		0.951			1.678	
<b>Ag</b>	-2.960	-2.089	0.871	0.800	1.390	1.527	-0.071	0.955	1.236
<b>W</b>	-8.830	-6.975	1.855			0.100			-0.518
<b>Os</b>	-8.220	-7.271	0.949						
<b>Ir</b>	6.960	-7.277	-0.317			1.156			1.262
<b>Pt</b>	5.870	-7.148	-1.278		1.180			1.819	
<b>Au</b>	-3.830	-2.924	0.906	1.692	1.746	1.498	0.786	1.293	1.196

3 and Figure S2.

1 **Table S3.** Aggregation energy ( $\Delta E_{agg}$ ) and single atom dissolution potential ( $U_{diss,SAC}^{n+}$ ) values

<b>M</b>	$E_{coh}$	$E_{emb}$	$\Delta E_{agg}$	$U_{diss,bulk}^{n+}$			$U_{diss,SAC}^{n+}$		
				<b>1+</b>	<b>2+</b>	<b>3+</b>	<b>1+</b>	<b>2+</b>	<b>3+</b>
<b>Cr</b>	-4.150	-6.360	-2.210		-0.913	-0.744		0.192	-0.007
<b>Mn</b>	-2.920	-5.521	-2.601		-1.185	-0.276		0.115	0.591
<b>Fe</b>	4.320	-6.376	-2.056		-0.447	-0.037		0.581	0.648
<b>Co</b>	4.470	-6.871	-2.401		-0.280	0.453		0.920	1.254
<b>Ni</b>	4.480	-7.008	-2.528		-0.257			1.007	
<b>Cu</b>	3.510	-5.294	-1.784	0.521	0.342	1.028	2.305	1.234	1.623
<b>Mo</b>	-6.840	-6.290	0.550			-0.200			-0.383
<b>Ru</b>	-6.800	-6.386	0.414		0.455	0.386		0.248	0.248
<b>Rh</b>	5.760	-6.247	-0.487	0.600		0.758	1.087		0.920
<b>Pd</b>	3.930	-5.543	-1.613		0.951			1.758	
<b>Ag</b>	-2.960	-2.303	0.657	0.800	1.390	1.527	0.142	1.061	1.307
<b>W</b>	-8.830	-8.214	0.616			0.100			-0.105
<b>Os</b>	-8.220	-7.581	0.639						
<b>Ir</b>	6.960	-7.427	-0.467			1.156			1.312
<b>Pt</b>	5.870	-7.325	-1.455		1.180			1.908	
<b>Au</b>	-3.830	-3.135	0.695	1.692	1.746	1.498	0.997	1.399	1.266

2 for all transition metal SACs in M@NGr/MoS<sub>2</sub> structure used for Figure 2b and Figure S2.

1 **Table S4.** Aggregation energy ( $\Delta E_{agg}$ ) and single atom dissolution potential ( $U_{diss,SAC}^{n+}$ ) values

<b>M</b>	$E_{coh}$	$E_{emb}$	$\Delta E_{agg}$	$U_{diss,bulk}^{n+}$			$U_{diss,SAC}^{n+}$		
				<b>1+</b>	<b>2+</b>	<b>3+</b>	<b>1+</b>	<b>2+</b>	<b>3+</b>
<b>Cr</b>	-4.150	-7.032	-2.882		-0.913	-0.744		0.528	0.217
<b>Mn</b>	-2.920	-6.623	-3.703		-1.185	-0.276		0.666	0.958
<b>Fe</b>	4.320	-6.822	-2.502		-0.447	-0.037		0.804	0.797
<b>Co</b>	4.470	-7.325	-2.855		-0.280	0.453		1.147	1.405
<b>Ni</b>	4.480	-7.388	-2.908		-0.257			1.197	
<b>Cu</b>	3.510	-5.756	-2.246	0.521	0.342	1.028	2.767	1.465	1.777
<b>Mo</b>	-6.840	-7.379	-0.539			-0.200			-0.020
<b>Ru</b>	-6.800	-6.908	-0.108		0.455	0.386		0.509	0.422
<b>Rh</b>	5.760	-6.749	-0.989	0.600		0.758	1.589		1.088
<b>Pd</b>	3.930	-5.940	-2.010		0.951			1.956	
<b>Ag</b>	-2.960	-3.573	-0.613	0.800	1.390	1.527	1.413	1.696	1.731
<b>W</b>	-8.830	-9.296	-0.466			0.100			0.255
<b>Os</b>	-8.220	-8.045	0.175						
<b>Ir</b>	6.960	-7.743	-0.783			1.156			1.417
<b>Pt</b>	5.870	-7.676	-1.806		1.180			2.083	
<b>Au</b>	-3.830	-3.969	-0.139	1.692	1.746	1.498	1.831	1.816	1.544

2 for all transition metal SACs in M@NGr/p-MoS<sub>2</sub> structure used for Figure 2b and Figure S2.

1 **Table S5.** Calculated work functions of isolated MoS<sub>2</sub>, p-MoS<sub>2</sub>, and M@NGr monolayers for  
2 selected transition metals. Work functions were computed from the plane-averaged  
3 electrostatic potential using DFT-D3-corrected GGA-PBE+U functional. The consistently  
4 lower work functions of all M@NGr systems relative to MoS<sub>2</sub> and p-MoS<sub>2</sub> confirm that  
5 electrons flow from M@NGr to the TMD layer upon heterostructure contact for all metals  
6 investigated.

<b>System</b>	<b>Work function (eV)</b>
MoS <sub>2</sub>	5.43
p-MoS <sub>2</sub>	5.82
Cr@NGr	3.80
Mn@NGr	3.83
Mo@NGr	3.85
Ru@NGr	3.81
Ag@NGr	3.78
W@NGr	3.76
Os@NGr	4.03
Au@NGr	3.58

1 **Table S6.** Calculated magnetic moments ( $\mu_B$ ) of the single-atom metal center in M@NGr,  
2 M@NGr/MoS<sub>2</sub>, and M@NGr/p-MoS<sub>2</sub>. For Mn, heterostructure formation induces a low-spin  
3 to high-spin transition in M@NGr/p-MoS<sub>2</sub>, driven by reduced crystal field splitting upon out-  
4 of-plane metal displacement. For Ag, the magnetic moment vanishes in M@NGr/p-MoS<sub>2</sub>,  
5 indicating complete stabilization of the d<sup>10</sup> closed-shell configuration upon Ag-N bond  
6 elongation. These element-specific spin-state responses provide supplementary stabilization

<b>M</b>	<b>M@NGr</b> <b>(<math>\mu_B</math>)</b>	<b>M@NGr/MoS<sub>2</sub></b> <b>(<math>\mu_B</math>)</b>	<b>M@NGr/p-MoS<sub>2</sub></b> <b>(<math>\mu_B</math>)</b>	<b>Spin state transition</b>
<b>Cr</b>	3.575	3.545	3.453	—
<b>Mn</b>	3.262	3.293	4.441	Low spin→High-spin
<b>Mo</b>	2.925	2.166	2.295	—
<b>Ru</b>	1.645	1.597	1.623	—
<b>Ag</b>	0.441	0.413	0.000	High spin→Low-spin
<b>W</b>	1.816	1.478	1.598	—
<b>Os</b>	1.928	1.771	1.814	—
<b>Au</b>	0.000	0.000	0.000	—

7 contributions beyond the general electrostatic mechanism.

1 **Table S7.** Average interlayer separation between the M@NGr layer and the top S plane of the  
2 TMD support, and the shortest M-S distance after structural relaxation. The interlayer  
3 separation is defined as the average height difference between the atoms in the M@NGr layer  
4 and the upper S atoms of the MoS<sub>2</sub> or p-MoS<sub>2</sub> layer. Underlined values denote M-S distances  
5 shorter than the sum of the corresponding covalent radii, indicating possible interfacial  
6 chemical bonding rather than purely vdW interaction.

<b>M</b>	<b>Interlayer distance (Å)</b>		<b>M-S distance (Å)</b>	
	<b>M@NGr/MoS<sub>2</sub></b>	<b>M@NGr/p-MoS<sub>2</sub></b>	<b>M@NGr/MoS<sub>2</sub></b>	<b>M@NGr/p-MoS<sub>2</sub></b>
<b>Cr</b>	3.38	3.35	3.12	2.71
<b>Mn</b>	3.40	3.38	3.42	2.52
<b>Fe</b>	3.40	3.41	3.45	3.36
<b>Co</b>	3.40	3.41	3.47	3.42
<b>Ni</b>	3.41	3.42	3.51	3.50
<b>Cu</b>	3.41	3.41	3.47	3.41
<b>Mo</b>	3.49	3.49	<u>2.41</u>	<u>2.46</u>
<b>Ru</b>	3.42	3.42	3.14	3.12
<b>Rh</b>	3.41	3.39	3.48	3.25
<b>Pd</b>	3.41	3.43	3.51	3.52
<b>Ag</b>	3.40	3.49	3.08	<u>2.38</u>
<b>W</b>	3.45	3.29	<u>2.35</u>	<u>2.55</u>
<b>Os</b>	3.48	3.45	3.16	3.16
<b>Ir</b>	3.41	3.43	3.52	3.55
<b>Pt</b>	3.41	3.43	3.52	3.55
<b>Au</b>	3.41	3.41	3.52	3.48