

Titanocene as a precursor for a cyclopentadienyl-free titanium(III)-manganese(II) clusters: a new approach for nano-size materials

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1.1 Preparation of Compound

General: All reactions were conducted under dinitrogen using standard Schlenk techniques. Manganese (powder, 99.99%) and 2-methoxyethanol (anhydrous liquid, 99.8%) were obtained from Aldrich and used without further purification unless stated otherwise. Bis(cyclopentadienyl)titanium dichloride (Cp_2TiCl_2 , powder, 97+%) was obtained from ABCR GmbH & Co. and used without further purification. GC/MS analyses were recorded on a gas chromatograph with a mass detector HP 5971A and an infrared detector HP5965B (Hewlett Packard). Microanalyses were conducted with an ARL Model 3410 + ICP spectrometer (Fisons Instruments) and a VarioEL III CHNS (in-house).

[Mn₄Ti₄(μ₂-Cl)₂Cl₆(μ₃,η²-OC₂H₅OMe)₂(μ₂,η²-OC₂H₅OMe)₁₀]·2C₆H₅CH₃ (1·2C₆H₅CH₃). A Schlenk flask was charged with Cp_2TiCl_2 (1.29 g; 5.18 mmol), metallic Mn (1.32 g; 24.0 mmol), 40 mL of MeOC₂H₅OH (38.6 g; 507 mmol) and toluene C₆H₅CH₃ (30 mL). The reaction mixture was stirred for 24 h at 60 °C. The resulting blue solution was filtered to remove Mn, and the clear solution was reduced in volume to approximately 15 mL. After 72 h blue block crystals were formed. They were filtered off, washed with hexanes (3 × 10 mL) and dried under vacuum (0.8 g, 0.46mmol, 36%). Calcd for C₅₀H₁₀₀Cl₈Mn₄O₂₄Ti₄ (MW, 1780.26): H 5.66 C 33.73 Cl 15.93 Ti 10.76 Mn 12.34. Found: H 5.89 C 34.43 Cl 15.21 Ti 10.46 Mn 12.01. GC/MS: CpH (MW, 66), CpH dimer (traces), 1-methylcyclohexa-1,4-diene (traces), cyclopentene (traces).

1.2 X-Ray Structure Determinations.

Crystal data and refinement details for the title compound are given in Table S1. The crystal was mounted on glass fibers and then flash-frozen to 100 K or 160 K (Oxford Cryosystem-Cryostream Cooler). Preliminary examination and intensities data collections were carried out on a Kuma KM4CCD κ-axis diffractometer with graphite-monochromated MoK_α radiation. Data were corrected for Lorentz, polarization and absorption effects. Data reduction and analysis were carried out with the Kuma Diffraction programs¹. The structure was solved by direct methods and refined by the full-matrix least-squares method on all F² data using the SHELXTL software². Carbon bonded hydrogen atoms were included in calculated positions and refined in the riding mode using SHELXTL default parameters. A correction for diffuse effects due to the inclusion of disordered solvent molecules in the crystal structure was made using the SQUEEZE option of PLATON.³ The total potential solvent volume per unit cell was calculated to be 256.00 Å³ (12.4 % of the cell volume).

Table S1. Crystal data and structure refinement for **1**.

| | |
|--|--|
| No CCDC | 671160 |
| Chemical formula | C ₃₆ H ₈₄ Cl ₈ Mn ₄ O ₂₄ Ti ₄ ·2(C ₇ H ₈) |
| <i>M</i> _r | 1780.26 |
| Cell setting, space group | Triclinic, P $\overline{1}$ |
| Temperature (K) | 100 (2) |
| <i>a</i> , <i>b</i> , <i>c</i> (Å) | 10.700(5), 10.787(6), 18.960(6) |
| α , β , γ (°) | 85.12(3), 74.68(3), 77.63(3) |
| <i>V</i> (Å ³) | 2060.7(16) |
| <i>Z</i> | 1 |
| <i>D</i> _x (Mg m ⁻³) | 1.435 |
| Radiation type | Mo $K\alpha$ radiation |
| μ (mm ⁻¹) | 1.28 |
| Crystal form, colour | Block, colorless |
| Crystal size (mm) | 0.21 × 0.10 × 0.09 |
| Data collection | |
| Diffractometer | Kuma KM-4 CCD κ-axis diffractometer |
| Data collection method | ω scans |
| Absorption correction | Analytical |
| <i>T</i> _{min} | 0.713 |
| <i>T</i> _{max} | 0.831 |
| No. of measured, independent and observed reflections | 28065, 9425, 5195 |
| Criterion for observed reflections | <i>I</i> > 2σ(<i>I</i>) |
| <i>R</i> _{int} | 0.069 |
| θ _{max} (°) | 29.1 |
| Refinement | |
| Refinement on | <i>F</i> ² |
| <i>R</i> [<i>F</i> ² > 2σ(<i>F</i> ²)], <i>wR</i> (<i>F</i> ²), <i>S</i> | 0.039, 0.075, 0.86 |
| No. of reflections | 9425 |
| No. of parameters | 413 |
| H-atom treatment | Mixture of independent and constrained refinement |
| Weighting scheme | <i>w</i> = 1/[σ ² (<i>F</i> _o ²) + (0.0398 <i>P</i>) ²] |
| (Δ/σ) _{max} | 0.002 |
| Δρ _{max} , Δρ _{min} (e Å ⁻³) | 0.42, -0.34 |

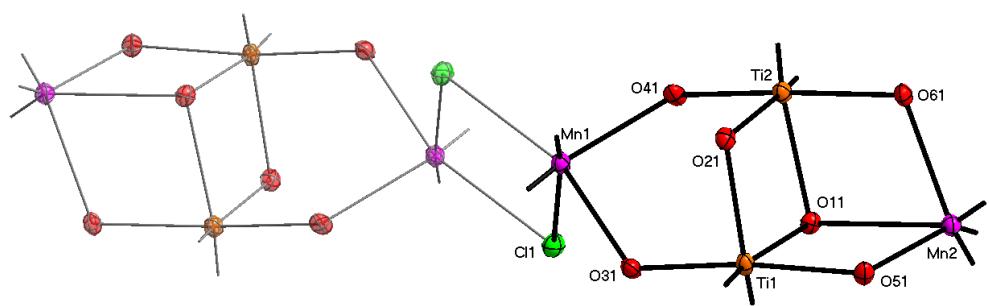
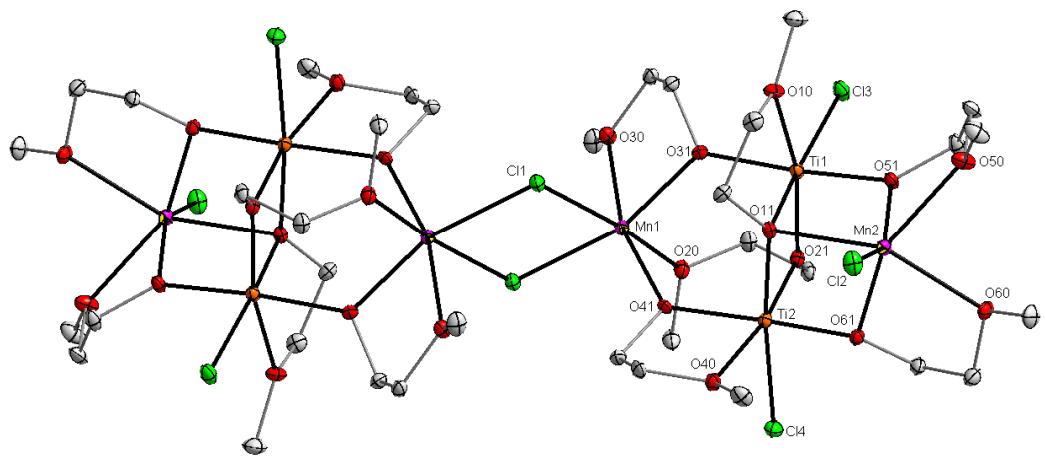


Figure. S1 The structure of **1** (top); its metal–oxygen core (bottom). H atoms omitted for clarity.

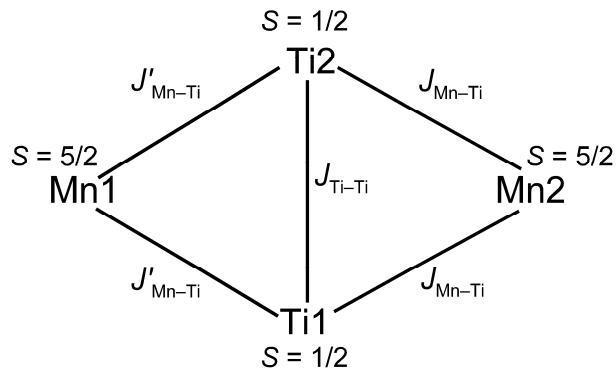
1.3 Magnetic Measurements.

Magnetic susceptibility down to 1.8 K at 500 mT and magnetization up to 5 T at 2.0 K were measured with a Quantum Design SQUID magnetometer. Diamagnetic correction was calculated using Pascal's constants.

Magnetic properties

The experimental data were fit using the Kambe vector coupling method. The exact symmetry of Mn_2Ti_2 cluster is low, but the reasonable simplifying approximation is possible by neglecting $\text{Mn}(1)\cdots\text{Mn}(2)$ ($6.190(2)$ Å) and taking equal $\text{Mn}(1)\cdots\text{Ti}(1)$ ($3.706(2)$ Å) and $\text{Mn}(1)\cdots\text{Ti}(2)$ ($3.631(2)$ Å) as well as $\text{Mn}(2)\cdots\text{Ti}(1)$ ($3.315(2)$ Å) and $\text{Mn}(2)\cdots\text{Ti}(2)$ ($3.350(2)$ Å) magnetic interactions (Scheme S1).

Scheme S1



Assuming the wingtip-wingtip interactions to be zero because of the large distance between Mn1 and Mn2 ions ($6.190(2)$ Å), the Heisenberg spin Hamiltonian can be written as:

$$\begin{aligned} \mathbf{H} = & -2J_{\text{Mn}-\text{Ti}} (\mathbf{S}_{\text{Mn}2}\cdot\mathbf{S}_{\text{Ti}1} + \mathbf{S}_{\text{Mn}2}\cdot\mathbf{S}_{\text{Ti}2}) - 2J'_{\text{Mn}-\text{Ti}} (\mathbf{S}_{\text{Mn}1}\cdot\mathbf{S}_{\text{Ti}1} + \mathbf{S}_{\text{Mn}1}\cdot\mathbf{S}_{\text{Ti}2}) \\ & - 2J_{\text{Ti}-\text{Ti}} \mathbf{S}_{\text{Ti}1}\cdot\mathbf{S}_{\text{Ti}2}. \end{aligned} \quad (1)$$

Unfortunately, the Kambe method can not be used to solve the system with C_s symmetry (Eq. 1). Thus an approximation was introduced, setting $J_{\text{Mn}-\text{Ti}} = J'_{\text{Mn}-\text{Ti}}$. Similar assumptions have been usually applied to the analysis of magnetic properties of tetramers.⁴⁻¹⁴ The spin Hamiltonian of Eq. 1 simplifies to:

$$\mathbf{H} = -2J_{\text{Mn}-\text{Ti}} (\mathbf{S}_{\text{Mn}1}\cdot\mathbf{S}_{\text{Ti}1} + \mathbf{S}_{\text{Mn}1}\cdot\mathbf{S}_{\text{Ti}2} + \mathbf{S}_{\text{Mn}2}\cdot\mathbf{S}_{\text{Ti}1} + \mathbf{S}_{\text{Mn}2}\cdot\mathbf{S}_{\text{Ti}2}) - 2J_{\text{Ti}-\text{Ti}} \mathbf{S}_{\text{Ti}1}\cdot\mathbf{S}_{\text{Ti}2}, \quad (2)$$

with the idealized C_{2v} symmetry. Defining $\mathbf{S}_{\text{Ti}} = \mathbf{S}_{\text{Ti}1} + \mathbf{S}_{\text{Ti}2}$, $\mathbf{S}_{\text{Mn}} = \mathbf{S}_{\text{Mn}1} + \mathbf{S}_{\text{Mn}2}$, and $\mathbf{S}_{\text{T}} = \mathbf{S}_{\text{Ti}} + \mathbf{S}_{\text{Mn}}$, one obtains 22 coupled spin states, which energies are given by:

$$\begin{aligned} E(S_{\text{T}}, S_{\text{Ti}}, S_{\text{Mn}}) = & -J_{\text{Ti}-\text{Mn}} [S_{\text{T}}(S_{\text{T}} + 1) - S_{\text{Ti}}(S_{\text{Ti}} + 1) - S_{\text{Mn}}(S_{\text{Mn}} + 1)] \\ & - J_{\text{Ti}-\text{Ti}} [S_{\text{Ti}}(S_{\text{Ti}} + 1) - 3/2]. \end{aligned} \quad (3)$$

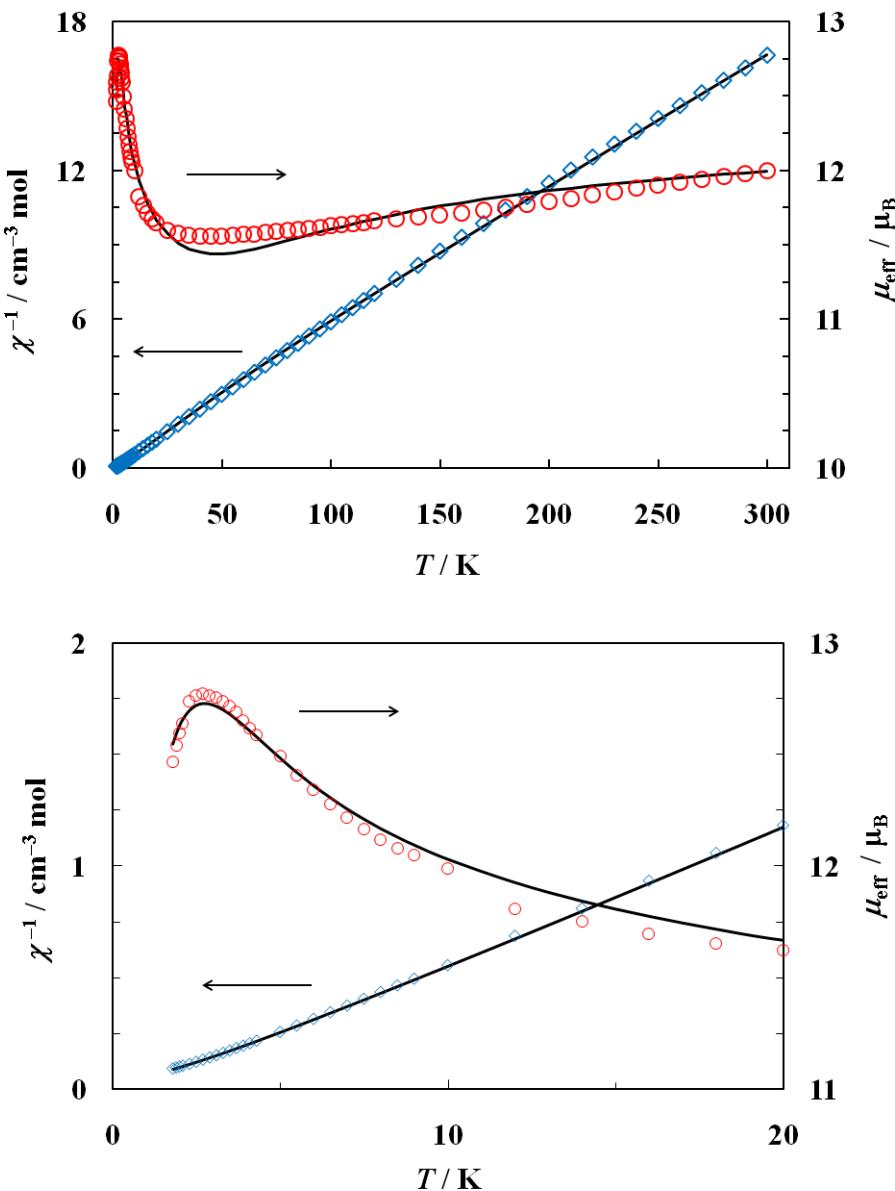


Figure S2. Variation of χ^{-1} (diamonds) and μ_{eff} (triangles) per $\text{Mn}^{\text{II}}_4\text{Ti}^{\text{III}}_4$ octamer with temperature. The solid lines represent the fit to the Kambe model (see text for details).

Magnetic exchange interactions between the both Mn_2Ti_2 subunits of **1** was estimated in molecular field approximation,¹⁵ and the expression for the magnetic susceptibility was $\chi_{\text{mf}} = \chi / (1 - 2zJ'\chi / Ng^2\beta^2)$, where χ is the susceptibility of Mn_2Ti_2 tetramer.

The magnetization of complex **1** was calculated in magnetic field of 500 mT using the method described by Belorizky.^{16,17} Least-squares fitting of the data gave $J_{\text{Ti-Ti}} = -28(1) \text{ cm}^{-1}$, $J_{\text{Ti-Mn}} = -5.2(2) \text{ cm}^{-1}$, $g = 1.982(2)$, and $zJ' = 0.0166(4) \text{ cm}^{-1}$. Temperature independent paramagnetism (TIP) was set at $130 \cdot 10^{-6} \text{ emu} \cdot \text{mole}^{-1}$ for Ti(III) and 0 for Mn(II) ions.¹⁸ No paramagnetic impurity was needed for the simulation. The agreement factor $R = \Sigma[(\chi T)_{\text{exp}} - (\chi T)_{\text{calc}}]^2 / \Sigma[(\chi T)_{\text{exp}}]^2$ was $8.7 \cdot 10^{-5}$ (68 points). The ground state was found to be $|S_{\text{T}}, S_{\text{Ti}}, S_{\text{Mn}}\rangle$

$= |4, 1, 5\rangle$ with the sixfold degenerate $|n, 0, n\rangle$ state ($n = 0-5$) at 7.3 cm^{-1} above the ground state.

To make sure if the global minimum had been obtained in the fitting procedure, the relative error surface diagram^{4,6,11,13} was calculated (Figure S3). The error surface was generated as a function of $J_{\text{Ti-Ti}}$ and $J_{\text{Ti-Mn}}$ with g and zJ' parameters changing freely ($g = 1.97-2.00$ and $zJ' = 0.014-0.028 \text{ cm}^{-1}$ for the fits comprised in the Figure). The minimum achieved on the error surface is rather swallow and only a little worse fit is obtained for $J_{\text{Ti-Ti}}$ ranging from -20 to -40 cm^{-1} and $J_{\text{Ti-Mn}}$ from -4 to -7 cm^{-1} . However, all results gave the $|4, 1, 5\rangle$ ground state and the $|n, 0, n\rangle$ sextet as the first excited state.

Magnetic interactions between Mn_2Ti_2 clusters are weak. $J_{\text{Mn-Mn}}$ between the bridging Mn^{II} ions may be estimated from zJ' employing the relation:¹⁹ $J_{\text{Mn-Mn}} S_{\text{Mn}}^2 = zJ' S_{\text{T}}^2$, with $S_{\text{Mn}} = 5/2$ and $S_{\text{T}} = 4$. The value of $J_{\text{Mn-Mn}} = 0.02 \text{ cm}^{-1}$ supports the assumption that both halves of the Mn_4Ti_4 cluster may be treated independently.

The low-temperature magnetization of **1** is almost saturated at 5 T at a value of *ca* $17 \mu_{\text{B}}$ (Figure S4). It is a little more than the value of $16 \mu_{\text{B}}$ expected for weakly interacting $S = 4$ dimer but below $20 \mu_{\text{B}}$ for $S = 5$ dimer ($g = 2.00$ in both cases). The experimental magnetization may be simulated using the exchange parameters obtained from fitting $\mu_{\text{eff}}(T)$ data, but with g value equal to 2.13 . A source of that discrepancy is not clear, but the lowest-temperature magnetization may be influenced by ZFS and other anisotropic effects. Assuming that only the ground state is populated (a first excited level at 7.3 cm^{-1}), we have tried to calculate magnetization by the full matrix diagonalization of the spin Hamiltonian:

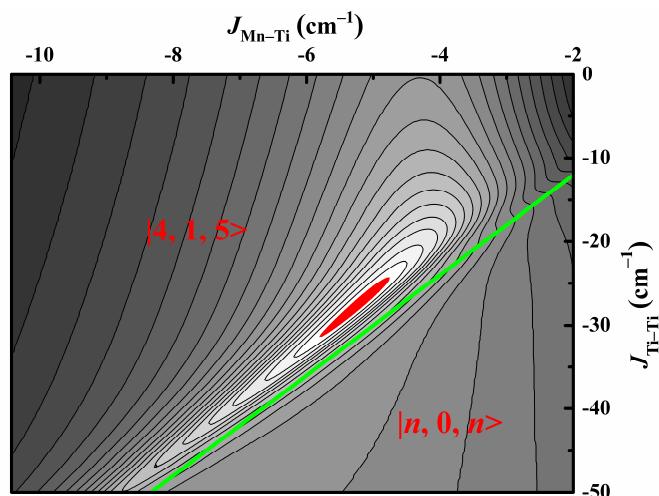


Figure S3. Contour projection of the error surface of the magnetic susceptibility fit. The green line separates the parts of $\{J_{\text{Mn-Ti}}, J_{\text{Ti-Ti}}\}$ space with different ground states.

$$\mathbf{H} = \mu_B \mathbf{B} \cdot \mathbf{g} \cdot \mathbf{S} + D[\mathbf{S}_z^2 - 1/3 \cdot \mathbf{S}(\mathbf{S} + 1)] + E(\mathbf{S}_x^2 - \mathbf{S}_y^2), \quad (4)$$

where $S = 4$, D and E are the second order ZFS parameters. Good fit was obtained for $D = -0.12 \text{ cm}^{-1}$, $E = 0$, but $g = 2.16$ was still too large.

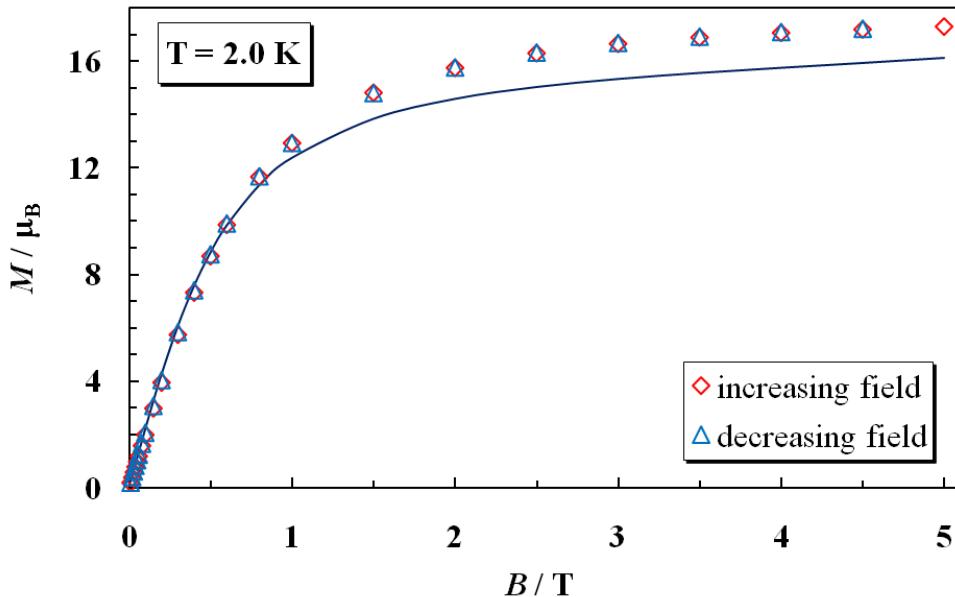


Figure S4. Field dependence of magnetization for **1**. The solid curve is calculated with the parameters obtained from fitting the susceptibility data.

Compounds contained magnetically coupled d¹–d⁵ system are very rare. We have found only one example of manganese(II)–titanium(III) compound, [(Cp)₂TiCl]₂MnCl₂.¹⁸ It is a trimer containing the Ti^{III}(μ-Cl)₂Mn^{II}(μ-Cl)₂Ti^{III} core. The magnetic parameters were estimated from fitting the magnetic susceptibility: $J_{\text{Ti-Mn}} = -8.24 \text{ cm}^{-1}$, $g = 2.00$ and $J_{\text{Ti-Ti}} \approx 0$. d¹–d⁵ magnetic exchange was also observed in vanadium(IV)–manganese(II) compounds. Oxalate-bridged [MnVO(ox)₂(phen)₂]·H₂O (phen = 1,10-phenanthroline) contains antiferromagnetically coupled Mn^{II}–V^{IV}O pairs and the fitting of the susceptibility data gave $J = -30.9 \text{ cm}^{-1}$, $g_{\text{Mn}} = 2.01$, $g_{\text{VO}} = 2.06$.²⁰ The magnetic behavior of [(VO)LMn(*m*-OAc)(H₂O)][ClO₄]·H₂O, where L is a symmetric tetraaminodiphenol macrocyclic ligand, is characteristic of an antiferromagnetically coupled system with $J = -14.2 \text{ cm}^{-1}$, $g_{\text{Mn}} = 2.00$, $g_{\text{VO}} = 1.86$.²¹ Tetranuclear compound {[Tp^{*}V(O)(CN)₂]₂[Mn(bipy)₂]₂[OTf]₂}·MeCN (Tp^{*} = hydridotris(3,5-dimethylpyrazol-1-yl)borate, OTf = trifluoromethanesulfonate, bipy = 2,2'-bipyridine) includes rectangular, cyanido-bridged Mn^{II}₂V^{IV}₂ clusters.²² Antiferromagnetic exchange between the Mn^{II} and V^{IV}O magnetic centers has been characterized by $J_{\text{V-Mn}}$ (sides) = -7.85 cm^{-1} , $J_{\text{V-V}}$ (diagonal) = -0.010 cm^{-1} , $g = 2.07$ ($S = 4$ ground state). On the

contrary to the above-mentioned Mn^{II}-V^{IV} compounds, [LV(μ -O)(μ -OAc)₂MnL] (L = hydridotris(pyrazolyl)borate) contains V=O...Mn bridge and exhibits ferromagnetic behavior ($J = +22.8 \text{ cm}^{-1}$, $g = 1.979$, $S = 3$ ground state).²³

There exist a lot of magnetically coupled compounds with Ti^{III}-Ti^{III} core, most of them are variations of [Cp₂Ti(μ -X)₂TiCp₂] system (Cp = cyclopentadienyl or its derivative). The halogeno-bridged complexes exhibit strong antiferromagnetic behavior with exchange integrals $-100 \div -160 \text{ cm}^{-1}$.^{24,25} Stronger antiferromagnetic interactions were found in dimers containing (μ -Cl)₃ bridges: $J = -171 \text{ cm}^{-1}$ for [NEt₄]₃[Ti₂Cl₉] (with unusually low g value 1.61 and $N\alpha = 880 \cdot 10^{-6} \text{ cm}^{-1}$),²⁶ $J = -262 \text{ cm}^{-1}$ for [PPh₄][Ti₂(μ -Cl)₃Cl₄(PEt₃)₂] \cdot CH₂Cl₂.²⁷

Oxo-bridged titanium(III) dimers are more diversified. Linear or nearly linear molecules are weakly antiferromagnetic or even ferromagnetic: $J = -7.8 \text{ cm}^{-1}$ in the linear [L₂Ti₂(NCO)₄(μ -O)] (L = 1,4,7-trimethyl-1,4,7-triazacyclononane)²⁸ and $J = +11 \text{ cm}^{-1}$ in [(Cp₂Ti)₂(μ -O)]²⁹ (Ti-O-Ti angle equal to 170.9°); both values of exchange constants were corroborated with theoretical calculations.^{30,31} Strong antiferromagnetic coupling was discovered in metoxo-bridged [Cp₂Ti(μ -OCH₃)]₂. The Ti-O-Ti bond angles are equal to 108.8° and the exchange constant J , determined from the temperature dependence of the EPR intensity, is equal to -268 cm^{-1} .²⁴

Small Ti-O-Ti bond angles in **1** (90.87° and 95.60°) and, consequently, very short Ti-Ti distance of 2.967 Å, could facilitate direct metal-metal antiferromagnetic exchange. As pointed out by Kahn,³² an indirect Ti-O-Ti exchange is strongly antiferromagnetic for the bond angle much larger than 90° (assuming titanium(III) magnetic orbitals of x^2-y^2 type), but an accidental orthogonality of the orbitals (and ferromagnetism) may takes place for bridging angle close to 90°. The low symmetry of the [Ti₂O₂] moiety excludes pure ferromagnetic behavior of **1**, but the exchange constant $J_{\text{Ti-Ti}}$ is relatively small.

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