# Lone Pair Electrons Minimize the Lattice Thermal Conductivity Supplementary information

Michele D. Nielsen<sup>1</sup>, Vidvuds Ozolins<sup>2,\*</sup> and Joseph P. Heremans<sup>1,3,\*</sup>

1. Department of Mechanical and Aerospace Engineering, Ohio State University, Columbus, Ohio, USA

2. Department of Materials Science and Engineering, University of California, Los Angeles, USA

3. Department of Physics, Ohio State University, Columbus, Ohio, USA

## 1. X-ray diffraction (XRD) data

XRD data on all samples are reported in Fig. S1; all index to rocksalt structures,

except for trigonal AgBiSe<sub>2</sub>. No reference data were found for NaBiTe<sub>2</sub>.



**Figure S1:** X-ray diffraction data on the compounds prepared for this study. Except for AgBiSe<sub>2</sub>, which was prepared in both its trigonal and rocksalt form by varying the heat treatment, all samples index in the NaCl structure. The dots are reference peaks from the International Centre for Diffraction Data PDF database.

The lattice parameters of the  $Ag_{1-x}Na_xSbTe_2$  alloys, shown in Fig. S2, follow a linear dependence on x, consistent with Vegard's law.

Electronic Supplementary Material (ESI) for Energy & Environmental Science This journal is O The Royal Society of Chemistry 2013



*Figure S2:* The lattice parameter (in the rocksalt unit cell) of  $Ag_{1-x}Na_xSbTe_2$  alloys as a function of composition.

We show typical temperature-dependent X-ray diffraction data in Fig. S3.



Figure S3 Typical temperature-dependent X-ray spectra of measured cubic compounds.

From these, the temperature-dependent lattice parameter can be calculated. The results are plotted Fig. S4: the lattice parameter increases linearly with *T*, so a single isotropic linear thermal expansion  $\alpha$  can be derived and is reported in Table 1 (main text). The data reported for the nominal composition AgSbTe<sub>2</sub> were in fact obtained on Ag<sub>0.73</sub>Sb<sub>1.12</sub>Te<sub>2</sub>.



**Figure S4** Temperature-dependence of the lattice parameter calculated from data like Fig. S3. The linear thermal expansion coefficients  $\alpha$  are calculated from the slope of these curves.

## 2. Electrical properties

The electrical resistivity of most compounds studied experimentally is shown in Fig. S5. NaSbSe<sub>2</sub> is too electrically insulating to be measured. Except for NaSbTe<sub>2</sub>, all compounds behave like semiconductors; NaBiTe<sub>2</sub> behaves like a doped one. Consistently with previous results on AgSbTe<sub>2</sub>,<sup>1</sup> Ag<sub>0.73</sub>Sb<sub>1.12</sub>Te<sub>2</sub> behaves like a very narrow-gap semiconductor. All this implies that the density of native defects, to the extent that they act as donors or acceptors, is limited to sub-percentage levels.



*Figure S5: Electrical conductivity of NaSbTe*<sub>2</sub>, *Ag*<sub>0.73</sub>*Sb*<sub>1.12</sub>*Te*<sub>2</sub>, *NaBiTe*<sub>2</sub>, *AgSbSe*<sub>2</sub>, *AgBiSe*<sub>2</sub> (*trigonal*), *and AgBiSe*<sub>2</sub> (*cubic*). *NaSbSe*<sub>2</sub> was electrically insulating.

AgSbTe<sub>2</sub>, Phys. Rev. B 77 245204 (2008)

<sup>&</sup>lt;sup>1</sup> V. Jovovic and J. P. Heremans, Energy Band Gap and Valence Band Structure of

#### 3. Calculated phonon spectra: summary

The calculated phonon dispersions are reported in sections 5 and 6. From them, modeaveraged Grüneisen parameters  $\gamma$ , sound velocities and the zone edge energies, expressed as Debye temperatures  $\Theta_i$  can be derived and tabulated in Table S1. Select data were reproduced in Table 1 of the main text. Lattice parameters *a* and bulk moduli *B* were calculated from the variation of the total energy vs. volume *E* (*V*) by fitting the calculated points to 3rd-order polynomials in *V*.

**Table S1.** The interatomic spacing (a/2), calculated bulk modulus B, mode-and-frequency averaged Grüneisen parameter $\gamma$ , sound velocities, and zone-boundary Debye energy  $k_B\Theta_i$  for phonons along the axes indicated.

Theory	a/2	В	γ	Sound velocities(m/s)			Zone-boundary energy $\Theta_i$ (K)		
	(nm)	(GPa)		Gamma-to-X [x 0 -x]			X [¾ 0 -¾]		
				TA1	TA2	LA	TA1	TA2	LA
NaSbSe <sub>2</sub>	0.2906	44	1.7	2132	2770	4329	76	89	114
NaSbTe <sub>2</sub>	0.3103	35	1.6	1664	2430	3630	63	85	94
NaBiTe <sub>2</sub>	0.3138	34	1.5	1295	2055	3131	40	56	56
				Gamma-to-K [0 x x]			K [0 ¾ ¾ ]		
				TA1	TA2	LA	TA1	TA2	LA
AgSbSe <sub>2</sub>	0.2819	78	3.5	1362	2105	3433	47	52	62
AgSbTe <sub>2</sub>	0.2972	67	2.3	1325	2469	3526	53	56	65
AgBiSe <sub>2</sub>	0.2893	74	2.5						
AgBiTe <sub>2</sub>	0.3018	66	2.5	1278	2116	3223	47	47	55

Note that the Debye temperature values  $\Theta_D$  fitted to the specific heat (Table 1) are nearly twice the Debye temperatures  $\Theta_i$  calculated from the zone-boundary acoustic phonon energies  $k_B\Theta_i$  in Table S1. Two factors contribute to this. Firstly, there is a contribution of the optical modes to the high-temperature value of *C*. Secondly, Fig. 2 of the main text shows that many acoustic modes do not have monotonic energy-momentum relations, so that the maximum phonon energy is higher than the energy  $k_B\Theta_i$  at the zone boundary. The calculated mode-averaged Grüneisen parameters for all 72 compounds (8 Group I cations  $\times$  3 Group V cations  $\times$  3 Group VI anions) are given in Table S2 below.

**Table S2:** Calculated mode-averaged Grüneisen parameters  $\gamma$  for all 72 compounds grouped by the group I cation. Higher  $\gamma$  values (stronger anharmonicity) correspond to warmer colors. Red squares indicate that the compounds have unstable phonon modes at the calculated equilibrium lattice parameter. AgSbSe<sub>2</sub> is predicted to have a weakly unstable L-point TA mode, and is found experimentally to have the lowest thermal conductivity of all materials considered in our study.



Electronic Supplementary Material (ESI) for Energy & Environmental Science This journal is O The Royal Society of Chemistry 2013



## 4. Literature with structural information of known *ABX*<sub>2</sub> compounds.

In agreement with our calculations, the crystal structures of several of the  $ABX_2$  compounds have been determined experimentally and found to be either not rocksalt based or resemble highly distorted variants of rocksalt. The references are:

CuSbS<sub>2</sub>: Razmara, M.F.; Henderson, C.M.B.; Pattrick, R.A.D. Mineralogical Magazine **61**, 79-88 (1997)

**CuBiS<sub>2</sub>:** Portheine, J.C.; Nowacki, W. Zeitschrift fuer Kristallographie, Kristallgeometrie, Kristallphysik, Kristallchemie **141**, 387-402 (1975);

CuSbSe<sub>2</sub>: Imamov, R.M.; Pinsker, Z.G.; Ivchenko, A.I. Kristallografiya **9**, 853-856 (1964);

AgSbS<sub>2</sub>: Effenberger, H.; Paar, W.H.; Topa, D.; Criddle, A.J.; Fleck, M. American Mineralogist **87**, 753-764 (2002);

NaAsS<sub>2</sub>: Palazzi, M.; Jaulmes, S. Acta Cryst. B 33, 908-910 (1977);

KSbS<sub>2</sub>: Graf, H.A.; Schaefer, H. Zeitschrift fuer Anorganische und Allgemeine Chemie **414**, 211-219 (1975);

CsSbS<sub>2</sub>: Kanishcheva, A.S.; Mikhailov, Yu.N.; Kuznetsov, V.G.; Batog, B.N. Doklady Akademii Nauk SSSR **251**, 603-605 (1980);

**RbAsSe<sub>2</sub>:** Sheldrick, W.S.; Haeusler, H.J. Zeitschrift fuer Anorganische und Allgemeine Chemie **561**, 139-148 (1988);

**RbSbS<sub>2</sub>:** Kanishcheva, A.S.; Kuznetsov, V.G.; Lazarev, V.B.; Tarasova, T.G. Zhurnal Strukturnoi Khimii **18**, 1069-1072 (1977)

Rocksalt CuBiSe<sub>2</sub>, AgAsSe<sub>2</sub>, AgSbSe<sub>2</sub>: Zhuze, V.P.; Sergeeva, V.M.; Shtrum, E.L. Soviet physics - Technical physics **3**, (10) 1925-1938 (1958); Tomaszewski, K. Phase Transition 38, 127-220 (1992); Geller, S.; Wernick, J.H. Acta Crystallographica 12, 46-54 (1959).

## 5. Calculated phonon spectra of the rocksalt compounds

The Brillouin zone of the rocksalt compounds, with the atomic arrangement of Fig. 1a, is shown in Fig. S6. The calculated phonon dispersions of rocksalt compounds follow with the compound name indicated above each panel. Compounds with spectra that contain imaginary calculated frequencies are not shown, with the exception of AgSbSe<sub>2</sub>, which has been synthesized in the cubic structure in this study and in previous ones (Ref. 13 of the main text).



*Figure S6.* Brillouin zone of cubic D4 (AF-IIb), and phonon dispersions of the stable noble metal D4 rocksalt-based  $ABX_2$  compounds with A=noble metal (see below)



CuSbTe2



CuBiTe2





Electronic Supplementary Material (ESI) for Energy & Environmental Science This journal is C The Royal Society of Chemistry 2013



AgBiS2





AgBiTe2



AuSbTe2



AuBiTe2

#### 6. Calculated phonon spectra of the rhombohedral rocksalt-based compounds

The Brillouin zone of the compounds with rhombohedral L11 crystal structures, the atomic arrangement shown in Fig. 1b, is shown in Fig. S7. The calculated phonon dispersions of those compounds follow, with the compound name indicated above each panel. Compounds with spectra that contain negative calculated frequencies are unstable and do not exist in the rhombohedral L11 crystal structure; accordingly, their phonons are not shown.



*Figure S7:* Brillouin zone of rhombohedral R3m (AF-II). Phonon dispersions of the stable R3m rocksalt-based ABX<sub>2</sub> compounds with A=alkali metal (see below).

Electronic Supplementary Material (ESI) for Energy & Environmental Science This journal is C The Royal Society of Chemistry 2013



NaAsSe2

Electronic Supplementary Material (ESI) for Energy & Environmental Science This journal is C The Royal Society of Chemistry 2013



NaAsTe2



NaSbS2

![](_page_22_Figure_1.jpeg)

NaSbTe2 250 200 150 Frequency (1/cm) 100 50 0 -50 B1IB XIQ FP1 ZIL P G Ζ G L Phonon wave vector

![](_page_24_Figure_1.jpeg)

NaBiS2

![](_page_25_Figure_1.jpeg)

![](_page_26_Figure_1.jpeg)

NaBiTe2

Electronic Supplementary Material (ESI) for Energy & Environmental Science This journal is C The Royal Society of Chemistry 2013

![](_page_27_Figure_1.jpeg)

KAsSe2

![](_page_28_Figure_1.jpeg)

![](_page_29_Figure_1.jpeg)

![](_page_30_Figure_1.jpeg)

KSbTe2

![](_page_31_Figure_1.jpeg)

KBiS2

![](_page_32_Figure_1.jpeg)

![](_page_33_Figure_1.jpeg)

KBiTe2

![](_page_34_Figure_1.jpeg)

CsSbSe2 250 200 150 Frequency (1/cm) 100 50 0 -50 B1IB Ζ XIQ ZIL P G G FP1 L Phonon wave vector

250 200 150 Frequency (1/cm) 100 50 0 -50 B1IB Ζ XIQ ZIL P G G FP1 L Phonon wave vector

CsSbTe2

![](_page_37_Figure_1.jpeg)

250 200 150 Frequency (1/cm) 100 50 0 -50 B1IB XIQ FP1 ZIL P G Ζ G L Phonon wave vector

CsBiSe2

![](_page_39_Figure_1.jpeg)

CsBiTe2

![](_page_40_Figure_1.jpeg)

RbAsTe2

250 200 150 Frequency (1/cm) 100 50 0 -50 B1IB XIQ FP1 ZIL P G Ζ G L Phonon wave vector

RbSbSe2

250 200 150 Frequency (1/cm) 100 50 0 -50 B1IB XIQ FP1 ZIL P G Ζ G L Phonon wave vector

RbSbTe2

![](_page_43_Figure_1.jpeg)

RbBiS2

250 200 150 Frequency (1/cm) 100 50 0 -50 B1IB XIQ FP1 ZIL P G Ζ G L Phonon wave vector

RbBiSe2

![](_page_45_Figure_1.jpeg)

RbBiTe2

Electronic Supplementary Material (ESI) for Energy & Environmental Science This journal is C The Royal Society of Chemistry 2013

![](_page_46_Figure_1.jpeg)

![](_page_47_Figure_1.jpeg)

TIAsTe2

![](_page_48_Figure_1.jpeg)

TISbTe2

![](_page_49_Figure_1.jpeg)

TISbSe2

Electronic Supplementary Material (ESI) for Energy & Environmental Science This journal is C The Royal Society of Chemistry 2013

TISbS2

![](_page_50_Figure_2.jpeg)

250 200 150 Frequency (1/cm) 100 50 0 -50 B1IB XIQ FP1 ZIL P G Ζ G L Phonon wave vector

TIBiTe2

250 200 150 Frequency (1/cm) 100 50 0 -50 B1IB XIQ FP1 ZIL P G Ζ G L Phonon wave vector

TIBiSe2

Electronic Supplementary Material (ESI) for Energy & Environmental Science This journal is C The Royal Society of Chemistry 2013

TIBiS2

![](_page_53_Figure_2.jpeg)