# **Electronic Supplementary Information**

# Predicting and rationalizing the effect of surface charge distribution and orientation on nano-wire based FET bio-sensors

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#### Conductance sensitivity with SMSCSL

The base conductance  $G_0$  of a cylindrical nano wire can be expressed as:

$$G_0 = \frac{\pi R^2 e \left( n_0 \mu_n + p_0 \mu_p \right)}{L}$$
(S1)

where *R* is the radius of the wire, *e* the elementary charge,  $n_0$  the wire electron density,  $\mu_n$  the wire electron mobility,  $p_0$  the wire hole density,  $\mu_p$  the wire hole mobility and *L* the length of the wire. For a nano-wire material with high enough doping, Equation (S1) can be approximated by considering only the electron or the hole density (only  $n_0$  or  $p_0$ ) and mobility (indicated simply as  $\mu$ ). For example, for a p-type doped nano-wire, it is possible to rewrite Equation (S1) as:

$$G_0 = \frac{\pi R^2 e p_0 \mu}{L} \tag{S2}$$

Following Sørensen et al. equations,<sup>1</sup> one can define the generic sensitivity of a cylindric nano wire as:

$$\frac{\Delta G}{G_0} = -\frac{2}{Re\,p_0}\Gamma\sigma_s\tag{S3}$$

where  $\sigma_s$  is the "sensed" charge density at the surface of the wire and  $\Gamma$  a dimensionless function quantifying the actual sensitivity of the nano-wire.

In the special case of a cylindrical nano-wire it is possible to express  $\Gamma$  analytically as:<sup>2</sup>

$$\Gamma = \frac{\varepsilon_1 K_0 \left(\frac{R+\Delta R}{\lambda_D}\right) \left(\frac{\lambda_D}{\lambda_{TF}}\right) \left(\frac{R}{\lambda_D}\right)}{\left[K_0 \left(\frac{R+\Delta R}{\lambda_D}\right) \left(\frac{\lambda_D}{R+\Delta R}\right) + \ln\left(\frac{R+\Delta R}{R}\right) K_1 \left(\frac{R+\Delta R}{\lambda_D}\right) \frac{\varepsilon_3}{\varepsilon_2}\right] \varepsilon_1 \left(\frac{R}{\lambda_{TF}}\right) I_1 \left(\frac{R}{\lambda_{TF}}\right) + \varepsilon_3 K_1 \left(\frac{R+\Delta R}{\lambda_D}\right) I_0 \left(\frac{R}{\lambda_{TF}}\right)}$$
(S4)

where  $\lambda_D$  is the Debye length of the buffer,  $\lambda_{TF}$  the Thomas-Fermi distance typical of the wire material,  $\varepsilon_1$ ,  $\varepsilon_2$  and  $\varepsilon_3$  are the relative permittivities of the wire, the oxide layer and the buffer, respectively,  $\Delta R$  is the oxide layer thickness,  $I_0$ ,  $I_1$ ,  $K_0$  and  $K_1$  are the modified Bessel functions of first and second kind, respectively.<sup>3</sup> Equation (S4) differs from the one reported in Ref. 1 since it considers explicitly a finite thickness for the oxide layer.

When considering the presence of additional charges at a distance *l* from the wire that gives rise to a charge density  $\sigma_b$ , it is possible to rewrite Equation (S3) as:

$$\frac{\Delta G}{G_0} = -\frac{2}{Re\,p_0} \Gamma\left(\Gamma_l \,\boldsymbol{\sigma}_b + \boldsymbol{\sigma}_s\right) \tag{S5}$$

where  $\Gamma_l$  is defined as in Equation (7) of Ref. 1:

$$\Gamma_l = 2\frac{R}{R+l} \left( 1 + \sqrt{\frac{R}{R+l}} \exp(l/\lambda_D) \right)^{-1}$$
(S6)

In the multiple charges model a series of *m* charges  $Q_i$  with distance  $l_i$  from the nano-wire surface gives rise to  $\sigma_{b_i}$  charge densities. The complete expression for the sensitivity is then given by:

$$\frac{\Delta G}{G_0} = -\frac{2}{Re\,p_0} \Gamma \left[ \sigma_s + \sum_i^m \left( \Gamma_{l_i} \,\sigma_{b_i} \right) \right] \tag{S7}$$

In the current contribution we assumed that the binding event does not perturb sensibly the charges already present on the surface of the nano-wire.  $\sigma_s$  can then be omitted, giving Equation (1).  $\sigma_{b_i}$  is computed as reported in Equation (2).

For a n-type nano-wire based BIOFET it is necessary to remove the factor -1 in Equations (1), (S3) and (S5), and use the appropriate value for the charge carrier density ( $n_0$  instead of  $p_0$ ) and mobility where needed.

The various parameters and their interpretation inside the multiple charges model are depicted in Fig. 1.

#### **Buffer solution Debye length**

Ions dissolved in solution lower the effectiveness of analyte charge(s) on the BIOFET signal. This effect is known as Debye screening and its characteristic Debye length  $\lambda_D$  can be evaluated as:

$$\lambda_D = \sqrt{\frac{\varepsilon_0 \varepsilon_3 K T}{2N_A e^2 I}} \tag{S8}$$

$$I = \frac{1}{2} \sum_{B=1}^{n} C_B Z_B^2$$
(S9)

where  $\varepsilon_0$  is the vacuum permittivity, *K* is the Boltzmann constant, *T* is the absolute temperature of the experiment (298.15 K),  $N_A$  is the Avogadro number and *I* is the ionic strength of the solution, which is based on the concentration *C* and charge *Z* of each ionic species present in the buffer.

#### Simulated nano-wire characteristics

To evaluate the change in conductance sensitivity (signal) depending on the different charge systems, we need to define a nanowire based BIOFET device for all our determinations. We consider a cylindric nano-wire. For p-type silicon nano-wire, we base our choice of parameters on the data reported by the Lieber group as described in Ref. 4, with the characteristics reported in Table 1S<sup>†</sup>. By using the average values<sup>4</sup> for the nano-wire radius, length, mobility and base transconductance, it is possible to retrieve an average value for the hole density  $p_0$  following Equation (S10), inverse of Equation (S2):

$$p_0 = \frac{G_0 L}{\pi R^2 e \mu} \tag{S10}$$

In this way we can retrieve the average value for  $p_0$  reported in Table 1S<sup>†</sup>. The value for the Thomas-Fermi length is computed using the following equation:

$$\lambda_{TF} = \sqrt{\frac{\hbar^2 \varepsilon_r \pi^{4/3}}{m^* e^2 p_0^{1/3}}}$$
(S11)

where  $\varepsilon_r = \varepsilon_0 \varepsilon_1$  and the effective hole mass  $m^*$  value was set to 0.54 times the mass of an electron at rest.<sup>5</sup> The value for the mobility  $\mu$  reported in Table 1S<sup>†</sup> differs from the value (0.056 m<sup>2</sup>V<sup>-1</sup>s<sup>-1</sup>) reported as average in Ref. 4. In fact we decide to lower such value to 0.01 in order to simulate the effects of the chemical treatments to which the wires are subjected when transformed into BIOFETs, since found appropriate in previous studies.<sup>6</sup>

For n-type  $In_2O_3$  nano-wire, we base our parameters on different sources.<sup>7–11</sup> Its characteristics are reported in Table 2S<sup>†</sup>. In order to apply Equation (S11) the value for  $m^*$  was set to 0.35 times the mass of an electron at rest.<sup>10</sup> In Equations (S10) and (S11)  $p_0$  has to be substituted with  $n_0$ .

#### Charge of a ionizable amino acid

For the ionizable amino acids HIS, LYS, ARG and the N terminal, the charge at a given pH value can be expressed as:

$$q = \frac{10^{pK_a - pH}}{1 + 10^{pK_a - pH}}$$
(S12)

For the ionizable amino acids ASP, GLU, CYS, TYR and the C terminal the charge is computed as q' = q - 1.

Parameter	Symbol	Value	(units)
Length	L	2000.0	(nm)
Radius	R	10.0	(nm)
Oxide layer thickness	$\Delta R$	2.0	(nm)
Thomas-Fermi Length	$\lambda_{TF}$	2.04	(nm)
Charge carrier density	$p_0$	$1.11E^{24}$	$(m^{-3})$
Charge carrier mobility	μ	0.01	$(m^2 V^{-1} s^{-1})$
Wire permittivity	$\epsilon_1$	12.0	$(\varepsilon_0)$
Oxide layer permittivity	$\epsilon_2$	3.9	$(\varepsilon_0)$
Solvent permittivity	ε3	78.0	$(\varepsilon_0)$

Table 1S<sup>†</sup> Parameters defining the simulated silicon p-type doped nano-wire based BIOFET.

 $\textbf{Table 2S} \dagger \text{ Parameters defining the simulated In}_2O_3 \text{ n-type semiconducting nano-wire based device, as retrieved from different sources.}^{7-11}$ 

Parameter	Symbol	Value	(units)	Reference
Length	L	2000.0	(nm)	9
Radius	R	5.0	(nm)	9
Oxide layer thickness	$\Delta R$	0.0	(nm)	11
Thomas-Fermi Length	$\lambda_{TF}$	1.179	(nm)	а
Charge carrier density	$n_0$	$4.6E^{25}$	$(m^{-3})$	7
Charge carrier mobility	μ	0.0078	$(m^2 V^{-1} s^{-1})$	7
Wire permittivity	$\epsilon_1$	9.0	$(\varepsilon_0)$	7,8
Solvent permittivity	ε3	78.0	$(\varepsilon_0)$	

<sup>*a*</sup> Value obtained with Equation (S11) and the data from Ref. 10.

## Other biomolecule mimicking charge distributions

**Fig. 1S**<sup>†</sup> Dependence of the conductance sensitivity (signal) of a Si p-type doped nano-BIOFET on the Debye length  $\lambda_D$ . Different biomolecule models are reported, each constituted by five charges (m = 5) at constant internal distances. The signal is evaluated at different distances D (in nm) from the nano-wire surface. a) Linear disposition, system charge -2, and b) their predicted signal. c) Compact disposition, system charge -3, and d) their predicted signal.



**Fig. 2S**<sup>†</sup> Dependence of the conductance sensitivity (signal) of a Si p-type doped nano-BIOFET on the Debye length  $\lambda_D$ . Different biomolecule models are reported, each constituted by three charges (m = 3, system charge -1) at constant internal distances, as reported. The signal is evaluated at different distances D (in nm) from the nano-wire surface. a) 0.5 nm inter-charge distance and b) the predicted signal. c) 1.0 nm inter-charge distance and d) the predicted signal with e) a detailed view of the **ZSP** region. f) 1.5 nm inter-charge distance and g) the predicted signal. The c and d panels are the same as in Fig. 2e,f.



**Fig. 3S**<sup>†</sup> Dependence of the conductance sensitivity (signal) of a Si p-type doped nano-BIOFET on the Debye length  $\lambda_D$ . The reported biomolecule model is constituted by m = 3 charges (system charge -1) at constant internal distances. The signal is evaluated at D = 1 nm for different Thomas-Fermi lengths. a) The system and b) its predicted signal for different values of  $\lambda_{TF}$ .



**Fig. 4S**<sup>†</sup> Dependence of the conductance sensitivity (signal) of a Si p-type doped nano-BIOFET on the Debye length  $\lambda_D$ . The reported biomolecule model is constituted by m = 2 charges (system charge -1) at constant internal distances. The signal is evaluated at different distances D (in nm) from the nano-wire surface. a) The system and b) its predicted signal.



Fig. 5S<sup>†</sup> Dependence of the conductance sensitivity (signal) of a Si p-type doped nano-BIOFET on the Debye length  $\lambda_D$ . The reported biomolecule model is constituted by m = 3 charges (system charge +1) at constant internal distances. The signal is evaluated at different distances D (in nm) from the nano-wire surface. a) The system and b) its predicted signal.



Fig. 6S<sup>†</sup> Dependence of the conductance sensitivity (signal) of a Si p-type doped nano-BIOFET on the Debye length  $\lambda_D$ . The reported biomolecule model is constituted by m = 3 charges (system charge -1) at constant internal distances. The signal is evaluated at different distances D (in nm) from the nano-wire surface. a) The system is the same as in Fig. 2e and Fig. 2S<sup>†</sup>c but approaches the surface from the opposite direction. b) The predicted signal.



**Fig. 7S**<sup>†</sup> Dependence of the conductance sensitivity (signal) of a Si p-type doped nano-BIOFET on the Debye length  $\lambda_D$ . The reported biomolecule model is constituted by m = 3 charges (system charge -1) at constant internal distances. a) The system is the same as in Fig. 2e and is anchored to the surface at a distance of 1 nm. (D = 2 nm) b) The same system is tilted by  $45^{\circ}$  around the anchor point. c) The predicted signals.



## **Protein treatment**

The crystal structure of each protein, as pdb file, is cleaned of any water molecules and extra residues. In order to ensure that every amino-acid has its proper lateral chain, we use the PDB2PQR program.<sup>12</sup> Using the program VMD<sup>13</sup> we locate the center of mass of each apoprotein, center the protein structure on these coordinates and orient it along its main rotational axis. The minimum and maximum coordinate values along each axes are found, and assumed as the coordinates of the vertices of the rectangular parallelepiped enclosing the protein, which we use as an approximation of the protein itself. We consider the wire surface as completely coated with proteins, that is the ideal situation when the linking of proteins to the wire surface is perfectly efficient. Once the face upon which the protein is residing on the wire surface is defined, by dividing the total wire surface by the area of such face gives the number of residing proteins, and consequently the charge density.

## Streptavidin

Tetramer:

Fig. 8S<sup>†</sup> reports the X-ray structure<sup>14</sup> of the tetrameric form of Streptavidin, with the dimensions of the encompassing parallelepiped. The cavities for the recognition of biotin are on the yz face, which has an area of 31.35 nm<sup>2</sup>. Supposing full coverage, the surface of the nano-wire described in Table 1S<sup>†</sup> can accomodate 4727 proteins, while the surface of the nano-wire described in Table 2S<sup>†</sup> can accomodate 2004 proteins.

Fig. 8S<sup>†</sup> The structure of the tetrameric form of Streptavidin.



Single chain:

Fig. 9S<sup> $\dagger$ </sup> reports the X-ray structure<sup>14</sup> of a single chain of Streptavidin, with the dimensions of the encompassing parallelepiped. The cavity for the recognition of biotin is on the *xy* face, which has an area of 10.20 nm<sup>2</sup>. Supposing full coverage, the nano-wire surface (Table 2S<sup> $\dagger$ </sup>) can accomodate 6160 proteins.

Fig. 9S<sup>†</sup> The structure of a single chain of Streptavidin.



Fig. 10S<sup>†</sup> Dependence of the conductance sensitivity (signal) of the  $In_2O_3$  n-type semiconducting nano-BIOFET from the Debye length  $\lambda_D$  for the Streptavidin protein (4 interdigitated chains, natural form) at different pH values.



**Fig. 11S**<sup>†</sup> Dependence of the conductance sensitivity (signal) of a  $In_2O_3$  n-type semiconducting based nano-BIOFET on the Debye length  $\lambda_D$  for the Streptavidin protein (1 single chain, N = 6160) linking to a biotinylated surface. a) Structure of the chain A of the Streptavidin protein (in gray) with the position of the biotin molecule in green. Superimposed, the position of the ionizable amino acids (m = 26, system charge -2.18 as computed at pH 7.4) as spheres differently colored, as in Fig. 4. The structure is aligned with the depicted axes. The nano-wire surface is supposed to be 2.65 nm distant from the protein COM (yellow sphere) along the negative direction of the *z* axis. b) The computed signal.



#### N protein

Fig. 12S<sup>†</sup> reports the NMR structure<sup>15</sup> of N protein, with the dimensions of the encompassing parallelepiped. In their experiments, Ishikawa et al.<sup>11</sup> functionalize their nano-wires with engineered Fibronectin, to which a linker molecule connects the N protein. In their supporting information is reported that the Fibronectin constitutes a layer ca. 4.3 nm thick. Fig. 13S<sup>†</sup> reports the structure of the linker molecule structure after a simple energy minimization, and the distance between the Fibronectin anchor atom (P) and the anchor for the N protein (S). Overall, the bio-functionalization layer can be thought as ca. 6 nm thick. See Ref. 11 for further details on the bio-functionalization layer.

Depending on the direction of approach to the nano-wire surface (Table  $2S^{\dagger}$ ) and supposing full coverage, a different number of proteins can be accomodated on the surface, as reported in Table  $3S^{\dagger}$ .

Fig. 12S<sup>†</sup> The structure of N protein.



Fig. 13S<sup>†</sup> Linker molecule structure.



**Table 3S**<sup> $\dagger$ </sup> N protein COM - nano-wire surface distance, when the protein approaches from different directions. For each distance, 6 nm for the bio-functionalization layer are included. By changing the approaching direction, the number of proteins *N* considered for full coverage changes accordingly.

Direction	Distance	N	Direction	Distance	N	Direction	Distance	N
	(nm)			(nm)			(nm)	
+x	7.8	3381	+y	9.3	5118	+z	8.6	6671
-x	7.9	3381	-y	8.3	5118	-z	//	//

# **Notes and References**

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