

## Supplementary Information

# Metal-Insulator Transition in Variably Doped $(\text{Bi}_{1-x}\text{Sb}_x)_2\text{Se}_3$ Nanosheets

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## Evidence of multi-band transport in doped $(\text{Bi}_{1-x}\text{Sb}_x)_2\text{Se}_3$ nanosheets in Hall effect data

In un-doped  $\text{Bi}_2\text{Se}_3$  nanosheet, due the high bulk carrier density ( $>10^{19}/\text{cm}^3$ ) from Se vacancies, the transport is dominated by the bulk carriers and Hall resistance shows a nearly perfect linear relationship with perpendicular magnetic field,  $B$  (solid black line, Figure S1). The extracted Hall slope,  $R_H=R_{xy}/B$ , from linear fitting of  $R_{xy}(B)$  is expected to relate to the bulk carrier density  $n_b = -1/(eR_H)$ . However, as Sb doping reduces the bulk carrier density and conductivity contribution, the effect of surface states conduction in topological insulators is expected to be readily observable through non-linear Hall slope data.<sup>1,2</sup> In a two band transport system, the  $R_{xy}$  is given by<sup>2</sup>

$$R_{xy}(B) = -(B/e)[(n_s\mu_s^2 + n_b\mu_b^2) + B^2\mu_s^2\mu_b^2(n_s+n_b)] / [(n_s\mu_s + n_b\mu_b)^2 + B^2\mu_s^2\mu_b^2(n_s+n_b)^2] \quad (1)$$

where  $n_s$ ,  $\mu_s$  and  $n_b$ ,  $\mu_b$  represent the carrier density and mobility for the surface and bulk channel.

Eq. 1 approaches the asymptotic behavior  $R_{xy}/B(\infty) = -(1/e)(n_s+n_b)^{-1}$  in the strong field limit, and

approaches  $R_{xy}/B(0) = -(1/e)(n_s\mu_s^2 + n_b\mu_b^2) / (n_s\mu_s + n_b\mu_b)^2$  in the small field limit. If  $\mu_b \neq \mu_s$ , then

one has a field dependent  $R_{xy}/B$  and the field Hall slope is larger than the high field slope:

$|R_{xy}/B(0)| > R_{xy}/B(\infty)$ . This effect was indeed observed in our doped  $(\text{Bi}_{1-x}\text{Sb}_x)_2\text{Se}_3$  nanosheet

samples. Figure S1 A compares the  $R_{xy}$  data for a pure sample and a  $(\text{Bi}_{1-x}\text{Sb}_x)_2\text{Se}_3$  nanosheet

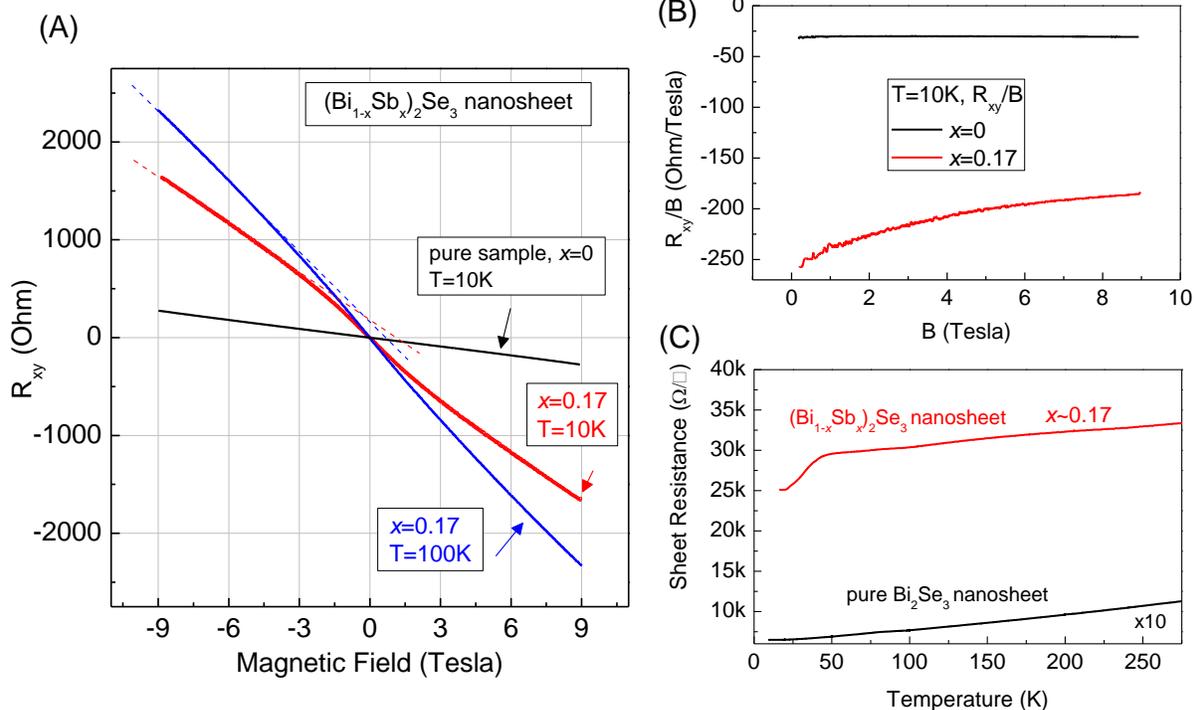
sample with  $x \sim 0.17$ . The non-linear  $R_{xy}(B)$  data in doped sample is illustrated by the linear

extrapolations of high field data not passing origin (dashed lines). The difference in Hall effect

between un-doped and doped samples is further highlighted in plotting  $R_{xy}/B$  vs. in Figure S1B.

In contrast to the constant Hall slope for un-doped sample, the magnitude of  $R_{xy}/B$  gradually

decreases and approaches a constant value at high  $B$ , as Eq.1 predicts.



**Fig. S1.** (A) Field dependent Hall resistance for pure  $\text{Bi}_2\text{Se}_3$  nanosheet and  $(\text{Bi}_{1-x}\text{Sb}_x)_2\text{Se}_3$  nanosheet with  $x=0.17$  at various temperatures, to show the nonlinear  $R_{xy}$  vs  $B$  in doped sample. (B) The Hall slope  $R_{xy}/B$  plotted as a function of  $B$  for pure and doped samples, highlighting the constant slope in pure sample vs. a  $B$ -dependent slope in doped sample. The non-linear  $R_{xy}$  vs  $B$  in doped sample is attributed to the increased contribution from the surface states channel in a two band (bulk plus surface) transport system. (C) Sheet resistance vs. temperature for the two  $(\text{Bi}_{1-x}\text{Sb}_x)_2\text{Se}_3$  nanosheet samples in (A) and (B).

## REFERENCES

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