Electronic Supplementary Information (ESI) for:

Deep tissue optical imaging of upconverting nanoparticles enabled by exploiting higher intrinsic quantum yield through using millisecond single pulse excitation with high peak power

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1. TEM image of the core NaYF₄:Yb³⁺,Tm³⁺ UCNPs



Figure S1. TEM image of core NaYF₄:Yb³⁺,Tm³⁺ nanoparticles. The mean diameter was determined to be approximately 32 nm.

2. The calculation of ion concentrations

The core-shell UCNPs used in this study are composed of two portions, *i.e.*, the optical active core NaY_{1-x-y}Yb_xTm_yF₄ (x and y denoting the molar concentrations in rare earth elements, respectively) and the shielding layer NaYF₄. The cores have an average diameter of D_1 =32 nm, and the core-shell particles have an average diameter of D_2 = 42 nm. In the rate equation model, N_0 and N_{Yb0} denote the number densities of Tm³⁺ and Yb³⁺ ions in the core portion, respectively.

The overall molar ratio between the three rare earth elements in the core-shell UCNPs was determined to be Y:Yb:Tm = 94.858%:5.100%:0.042% by ICP-OES analysis on a PerkinElmer Optima 8300. The molar weight of NaY_{1-x-y}Yb_xTm_yF₄ is thus obtained

$$Mr_1 = 22.99 + 88.91 \cdot (1 - x - y) + 173.04 \cdot x + 168.93 \cdot y + 19.00 \times 4$$
$$= 187.9 + 84.13 \cdot x + 80.02 \cdot y \text{ g/mol},$$

while the molar weight of NaYF₄ is

$$Mr_2 = 22.99 + 88.91 + 19.00 \times 4 = 187.9 \text{ g/mol}$$
.

The density of NaYF₄ crystal was determined to be $\rho = 4.2 \text{ g/cm}^3$ by X-ray energy dispersive spectroscopy (XEDS) analysis. The mass of the core portion in a single core-shell UCNP can be calculated by

$$m_{\text{core}} = V_{\text{core}} \cdot \rho = \frac{4}{3} \pi \left(\frac{D_1}{2}\right)^3 \cdot \rho = 2.3 \times 10^{-17} \text{g},$$

while the mass of the shell portion can be calculated by

$$m_{\rm shell} = V_{\rm shell} \cdot \rho = \left[\frac{4}{3}\pi \left(\frac{D_2}{2}\right)^3 - \frac{4}{3}\pi \left(\frac{D_1}{2}\right)^3\right] \cdot \rho = 2.9 \times 10^{-17} {\rm g}.$$

The overall molar ratio between the rare earth elements can be calculated by

Y: Yb: Tm =
$$\left[\frac{m_{\text{core}}}{Mr_1} \cdot (1 - x - y) + \frac{m_{\text{shell}}}{Mr_2}\right] : \frac{m_{\text{core}}}{Mr_1} \cdot x: \frac{m_{\text{core}}}{Mr_1} \cdot y$$

= 94.858%: 5.100%: 0.042%

x and y are thus determined to be

and

y=0.098%

The concentration of Yb³⁺ ions is obtained by

$$N_{\rm Yb0} = \frac{\rho \cdot 1}{Mr_1} \cdot x \cdot N_{\rm A} = 1.52 \times 10^{21} {\rm cm}^{-3}$$

involving N_A the Avogadro constant. Similarly, the concentration of Tm^{3+} ions is obtained by

$$N_0 = \frac{\rho \cdot 1}{Mr_1} \cdot y \cdot N_A = 1.25 \times 10^{19} \text{cm}^{-3}$$

3. Estimated number of ions in a single nanoparticle

In a single nanoparticle, the numbers of ions can be obtained by

$$n_{\rm Yb} = N_{\rm Yb0} \cdot V_{\rm core} = 26080$$

and

$$n_{\rm Tm} = N_0 \cdot V_{\rm core} = 215$$

for Yb³⁺ and Tm³⁺ ions, respectively.

4. The selection of the ETU rates

a. The selection of C_0 and C_1

The power density dependent steady-state quantum yield of the used core-shell UCNPs have been measured and reported recently in our previous work [Liu *et al.*, Nanoscale, 2013, **5**, 4770-4775]. In this study, the ETU rates C_0 and C_1 were selected on the principle of giving the best fitting between the simulated and experimental results, which ends up with

$$C_0 = 1.6 \times 10^{-18} \text{ cm}^3/\text{s},$$

and

$$C_1 = 6.2 \times 10^{-16} \text{ cm}^3/\text{s}$$

Since NaYF₄ is poorly characterized in bulk form, we compared such values with those reported in Ref. 28 for LiYF₄ crystal, which has similar phonon energies to NaYF₄. The selected value for C_0 here is several times smaller than that for sample no. 5 (C_0 =9.1×10⁻¹⁸ cm³/s) in Ref. 28 which has the most similar molar concentrations with the UCNPs investigated. This is reasonable considering the ~10 times less doping level of Tm³⁺ ions in this study, as C_0 decreases with the decrease of the molar concentration of Tm³⁺ ions, referring to the energy-transfer parameters for samples no. 2 and no. 4 in TABLE III reported in Ref. 28. The value for C_1 is on the same order of that reported in Ref. 28 for sample no. 5, exhibiting its reasonability.

b. The selection of C_2

The energy gap between states ${}^{3}H_{5}$ (Tm³⁺) and ${}^{3}H_{6}$ (Tm³⁺) (~8500 cm⁻¹) is similar to that between ${}^{1}G_{4}$ and ${}^{3}H_{4}$ ((~8620 cm⁻¹)²⁷. Thus, the energy mismatches of ETU0 and ETU1 are very similar. Using the energy-gap law formula [Diening *et al.*, J. Appl. Phys. 1998, **84**, 5900-5904], C_{2} was estimated to be similar to C_{0} . In this study, we take the same value for C_{2} with C_{0} .





Figure S2. Decay profile of (a) the NIR UC emission at 800 nm from Tm³⁺ ions and (b) the emission at 1035 nm from Yb³⁺ ions under a square-wave excitation at 975 nm.

6. Simulated cumulative QY at the average power density of 1 W/cm²



Figure S3. The temporally cumulative QY of NIR UC emission under CW excitation and under pulsed excitation with a fixed duty cycle of 4% and with various repetition rates. All the excitation approaches have the same average power density of 1 W/cm²



7. The influence of ETU rates on the simulated signal gain



8. The calculation of maximum permissible power-density for exposure to human skin

a. For CW excitation

The Maximum permissible exposure (MPE) for skin exposure to a CW laser beam is given by

$$MPE_{CW} = 0.2C_A W \cdot cm^{-2} \tag{1}$$

where C_A is a wavelength dependent parameter,

$$C_{\rm A} = 10^{2(\lambda - 0.700)} \tag{2}$$

with λ [µm] the wavelength. For the wavelength of 975 nm,

$$C_{\rm A} = 3.548$$
 (3)

and

$$MPE_{CW} = 709.6 \text{ mW} \cdot \text{cm}^{-2} \tag{4}$$

b. For repetitive pulse excitation

For repetitive-pulse lasers, two rules apply to determine the MPEs for skin exposure.

Rule 1. Single-pulse limit.

Exposure of the skin shall not exceed the MPE based upon a single-pulse exposure.

Rule 2. Average-power limit

The average irradiance of the pulse train shall not exceed the MPE applicable for the total pulse train, duration T.

The MPE found using Rule 1 for a 975 nm laser where T = 10 s, pulse width t = 20 ms and F = 2 Hz is:

$$MPE_{sp} = 1.1C_{A}t^{0.25} \text{ J} \cdot \text{cm}^{-2}$$

= 1.1 × 3.548 × 0.02^{0.25} J · cm⁻²
= 1.47 J · cm⁻² (5)

This MPE in terms of average power for Rule 1 is:

$$MPE: E = MPE_{sp} \times F = 2.9 \text{ W} \cdot \text{cm}^{-2}$$
(6)

The MPE found using Rule 2 is:

MPE:
$$H_{\text{group}} = 1.1C_A t^{0.25} \text{ J} \cdot \text{cm}^{-2}$$

= $1.1 \times 3.548 \times 10^{0.25} \text{ J} \cdot \text{cm}^{-2}$
= $6.94 \text{ J} \cdot \text{cm}^{-2}$ (7)

In terms of average irradiance, the MPE is

MPE:
$$E = \frac{MPE: H_{group}}{10 \text{ s}} = 694 \text{ mW} \cdot \text{cm}^{-2}$$
 (8)

Resultant MPE:

The MPE found using Rule 2 is the correct MPE to apply, since it is the smallest. The corresponding peak power-density in the pulse duration is

$$\rho = 694 \text{ mW} \cdot \text{cm}^{-2} \times \frac{\frac{1}{2 \text{ Hz}}}{20 \text{ ms}} = 17.4 \text{ W} \cdot \text{cm}^{-2}$$
(9)

c. For single pulse excitation

The MPE for the single pulse excitation can be calculated by

$$MPE_{sp} = \frac{1.1C_A t^{0.25}}{t} W \cdot cm^{-2}$$
(10)

For a 50 ms single pulse,

$$MPE_{sp} = 36.9 \ W \cdot cm^{-2}$$
(11)

Supporting references:

1. H. Liu, C. T. Xu, D. Lindgren, H. Xie, D. Thomas, C. Gundlach and S. Andersson-Engels, *Nanoscale*, 2013, **5**, 4770-4775.

2. A. Diening, P. E.-A. Möbert and G. Huber, J. Appl. Phys., 1998, 84, 5900-5904.