# Supplementary Information for

# Photoluminescence of Monolayer MoS<sub>2</sub> on LaAlO<sub>3</sub> and SrTiO<sub>3</sub> substrates

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# 1. Calculation of the enhancement factor for Raman and Photoluminescence (PL) intensity normalization.

The enhancement factor is defined as  $\Gamma^{-1} = \frac{I_{MoS_2}^{freesanding}}{I_{MoS_2}^{on-substrate}}$ . The calculation method is based on

the model including optical interference and absorption effects due to the substrate geometry, which is described in Ref.<sup>1</sup> By multiplying  $\Gamma^{-1}$  for each measured Raman and PL spectrum, the optical interference effect attributed to the different samples is eliminated and intrinsic difference between each system can be compared. Fig.S1(a) shows the calculated enhancement factor  $\Gamma^{-1}$  of monolayer to four-layer MoS<sub>2</sub> on SiO<sub>2</sub>, LaAlO<sub>3</sub>, Gel-film and SrTiO<sub>3</sub> substrate. Fig.S1(b) and (c) show the comparison of the raw and normalized Raman and PL spectra of monolayer MoS<sub>2</sub> on different substrates.

## 2. Thickness-dependence of Raman spectra for MoS<sub>2</sub> on different substrates

Fig.S2 shows the normalized Raman spectra of 1-4L MoS<sub>2</sub> on different types of substrates.

From the Raman spectra, for each kind of substrate, both  $E_{2g}^1$  and  $A_{1g}$  modes of MoS<sub>2</sub> are clearly observed. The dependence of these two mode frequencies on layer thickness exhibits similar trend for four selected substrates, i.e.  $E_{2g}^1$  mode moves to lower frequency end while  $A_{1g}$  mode shifts to higher frequency end when the number of layers increases.

#### 3. Thickness-dependence of PL spectra for MoS<sub>2</sub> on different substrates

Fig.S3 shows the normalized PL spectra of 1-4L MoS<sub>2</sub> on different substrates. It is known that, for monolayer MoS<sub>2</sub>, PL spectrum is composed with two excitonic peaks A and B, associated with direct optical transitions from the lowest conduction bands to the highest spin-split valence bands.<sup>2</sup> From Fig.S3, both excitonic peaks can be clearly observed in all samples. Meanwhile, all the samples show very similar thickness-dependence, i.e. when the number of layers decreases, the intensity of peak A dramatically increases and reaches to the maximum for monolayer thin film, due to the indirect-to-direct band transition. For the bilayer samples, we also observed an emission peak around 1.55 eV (Peak I) which is ascribed to indirect band transition. Besides the change of peak intensity, the position and shape of peak also changes with thin film thickness. Peak A moves to higher energy end when thin film thickness decreases.

# 4. The fitting of PL spectra of monolayer MoS<sub>2</sub> on different substrates.

We fit the experimental PL data of monolayer  $MoS_2$  on  $SiO_2$ , LaAlO<sub>3</sub>, Gel-film and SrTiO<sub>3</sub> substrate. The fits are composed by three Lorenz functions corresponding neutral excitons (A<sup>0</sup>), trions (A<sup>-</sup>) and the direct band transition (peak B). The fitting results are shown in Fig.S4.

#### 5. The electron density calculation

We estimate the electron density in monolayer  $MoS_2$  from the analysis of PL intensity of neutral excitons and trions emissions. According to mass action model which is based on the dynamic equilibrium between neutral excitons (A<sup>0</sup>), free electrons and trions (A<sup>-</sup>), the following relationship is obtained:<sup>3-5</sup>

$$\frac{N_{A^0} n_{el}}{N_{A^-}} = (\frac{4m_{A^0} m_e}{\pi h^2 m_{A^-}}) k_B T \exp(-\frac{E_b}{k_B T})$$

Where  $k_B$  is the Boltzmann constant,  $E_b$  is the trion binding energy (~20 meV),<sup>6</sup> T is the temperature.  $m_e$  (0.35  $m_0$ ),  $m_{A^0}$  (0.8  $m_0$ ) and  $m_{A^-}$  (1.15  $m_0$ ) are the effective mass of neutral excitons, trions and electrons, respectively, where  $m_0$  is the mass of free electrons.<sup>7</sup>  $N_{A^0}$  and  $N_{A^-}$  are the population of neutral exciton and trion, repectively.  $n_{el}$  is the electron density.

To establish the relationship between PL intensity and the population of neutral exciton and trion, we consider a three level model that includes a trion, an exciton and the ground state.<sup>8</sup> Based on this model, the PL intensity weight can be related to the population of neutral exciton and trion as:

$$\frac{I_{A^{-}}}{I_{total}} = \frac{I_{A^{-}}}{I_{A^{-}} + I_{A^{0}}} = \frac{\frac{\gamma_{A^{-}}}{\gamma_{A^{0}}} \frac{N_{A^{-}}}{N_{A^{0}}}}{1 + \frac{\gamma_{A^{-}}}{\gamma_{A^{0}}} \frac{N_{A^{-}}}{N_{A^{0}}}}$$

Where  $I_{A^0}$  and  $I_{A^-}$  are integrated PL intensity of neutral exciton and trion, respectively.  $\gamma_{A^0}$  and  $\gamma_{A^-}$  are the relative decay rate of the neutral exciton and trion, respectively. The value of  $\frac{\gamma_{A^-}}{\gamma_{A^0}}$  is ~ 0.15 according to Mouri et al.'s study.<sup>8</sup>

Combining the three level model and mass action model, the  $n_{el}$  can be expressed as:

$$n_{el} = \frac{\frac{I_{A^-}}{I_{total}}}{\frac{\gamma_{A^-}}{\gamma_{A^0}} (1 - \frac{I_{A^-}}{I_{total}})} \left[ (\frac{4m_{A^0}m_e}{\pi h^2 m_{A^-}}) k_B T \exp(-\frac{E_b}{k_B T}) \right]$$

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Figure.S1. (a).The calculated enhancement factor of monolayer to 4-layer  $MoS_2$  on  $SiO_2$ , LAO,Gel-film and STO substrates. (b). The raw and normalized Raman spectra of monolayer  $MoS_2$  on different substrates . (c). The raw and normalized PL spectra of monolayer  $MoS_2$  on different substrates.



Figure S1. Normalized Raman spectra, frequencies of the  $A_{1g}$  and  $E^{1}_{2g}$  modes and the difference between the two modes of monolayer to 4-layer MoS<sub>2</sub> on SiO<sub>2</sub>, LAO, STO and Gel-film substrates.



Figure S3. Thickness-dependent of the normalized PL spectra of  $MoS_2$  on  $SiO_2$ , LAO, Gel-film and STO substrates. The inserts are PL spectra of substrates.



Figure.S4. The fitting of PL spectra of monolayer  $MoS_2$  on  $SiO_2$  (a), LAO (b), Gel-film (c), STO (d) substrates. The open circle symbols are experimental data. The thin solid lines are the Lorentz functions that compose the fit.