

ARTICLE

## Plasmonic coupling with most of the transition metals: a new family of broad band and near infrared nanoantennas

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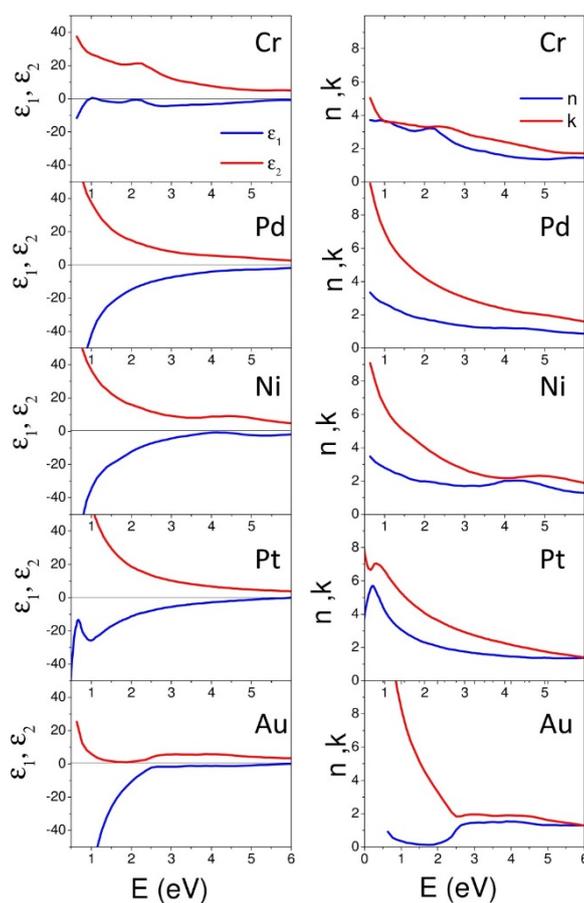


Figure S11: dielectric functions  $\epsilon=\epsilon_1+i\epsilon_2$  (left) and corresponding optical indexes  $\tilde{n}=n+ik$  (right) of various metals (Cr, Pd, Ni, Pt, Au) (from references <sup>1,2</sup> and <sup>3</sup> for Pt)

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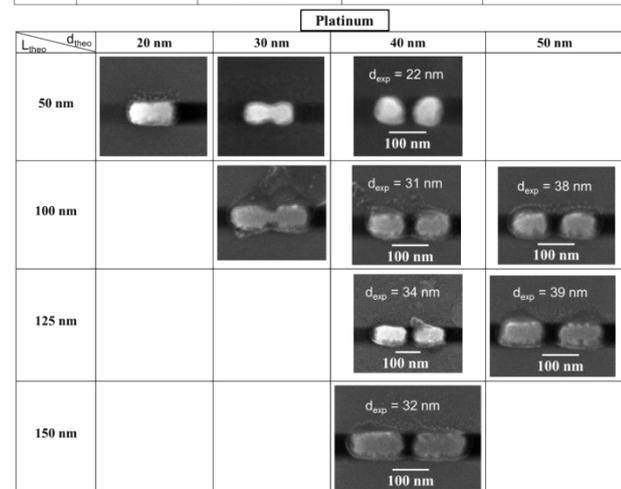
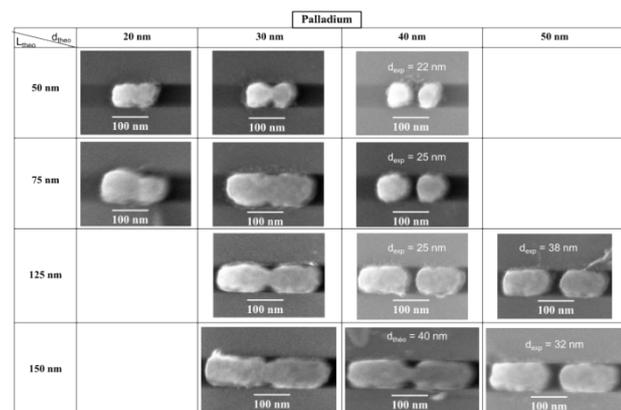
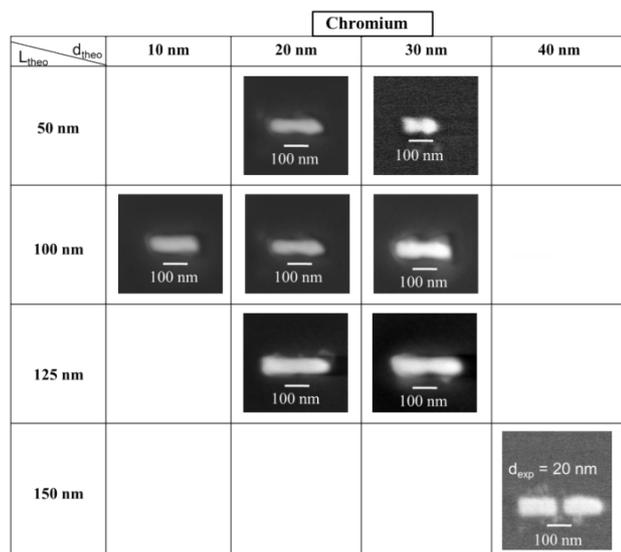


Figure S12: SEM images of parallelepipedal dimers elaborated by nanolithography. For theoretical interparticle distances  $d_{theo}$  lower than 30 nm, the parallelepipeds are in contact, whereas for theoretical interparticle distances larger than 40 nm, the parallelepipeds do not touch.

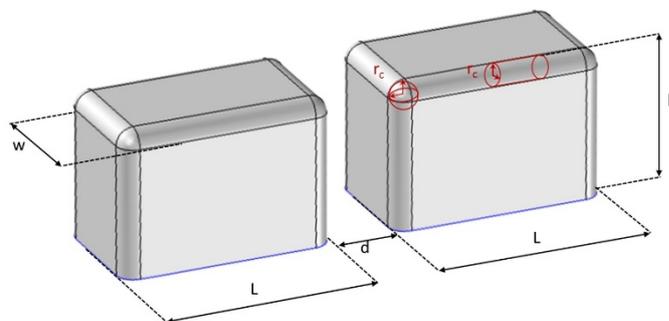


Figure S13: parameters characterizing the parallelepipedal monomers and dimers involved in calculations: Length  $L$ , width  $w$ , height  $h$ , rounding parameter  $r_c$ , interparticle distance  $d$ .

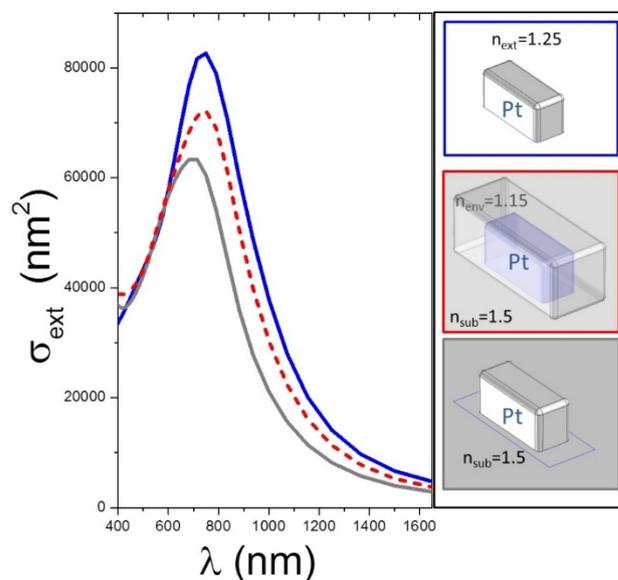


Figure S14: calculated extinction cross-sections for a platinum parallelepiped ( $L=150$  nm,  $w=60$  nm,  $h=70$  nm,  $r_c=8$  nm) :

(blue): NA in a homogeneous medium ( $n_{ext}=1.25$ ); (red): NA on a substrate ( $n_{sub}=1.5$ ) surrounded from its bottom by a local medium defined as a parallelepiped of dimensions:  $L_e=250$  nm,  $w_e=100$  nm,  $h_e=100$  nm,  $r_{ec}=8$  nm and of optical index  $n_{env}=1.15$ ; (gray) : NA on a substrate of optical index  $n_{sub}=1.5$

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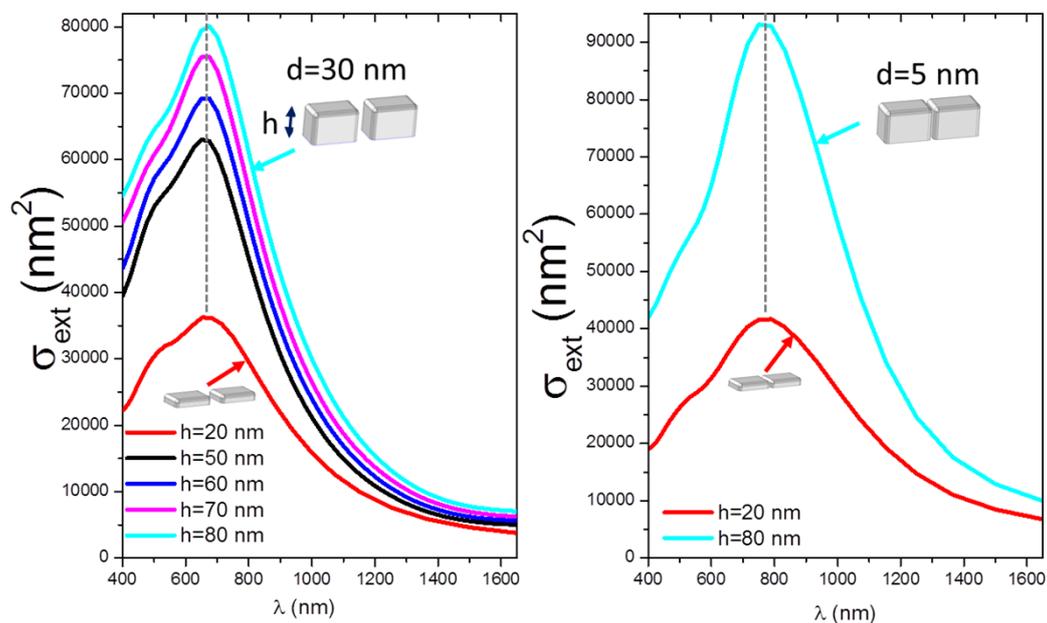


Figure S15: influence of the height  $h$  of the parallelepipeds ( $L=100$  nm,  $w=60$  nm,  $r_c=8$  nm) on the extinction cross-sections (left:  $d=30$  nm, right:  $d=5$  nm). The reducing of the height of the parallelepipeds induces a damping of the LSPR, but no spectral shift is observed, even for small interparticle distances

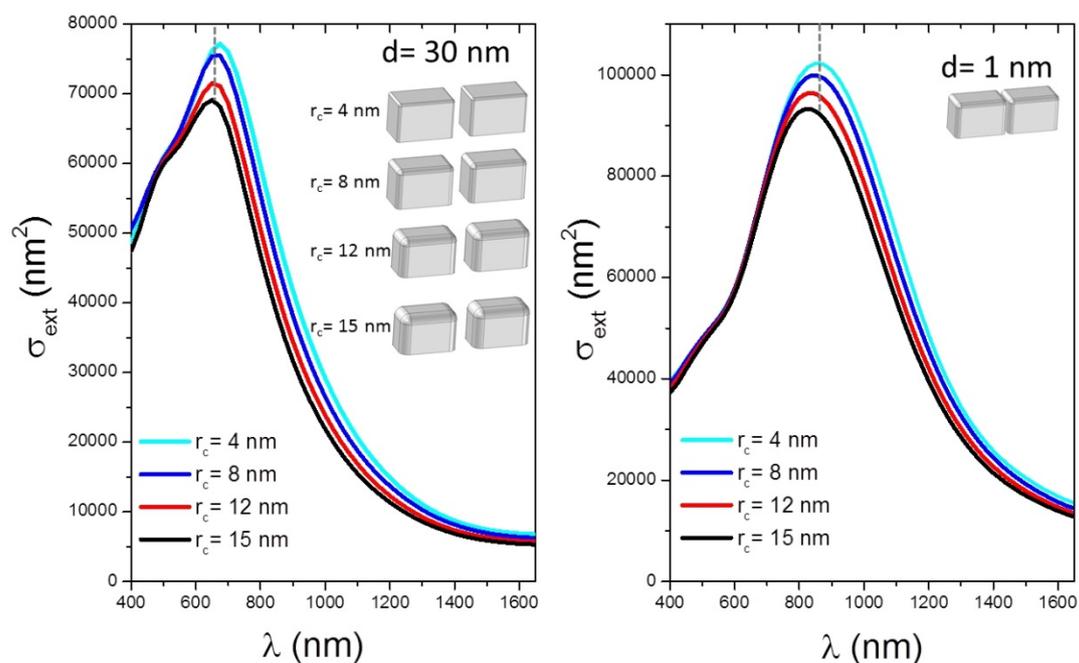


Figure S16: influence of the rounding parameter  $r_c$  characterizing the parallelepipeds ( $L=100$  nm,  $w=60$  nm,  $h=70$  nm) on the extinction cross-sections (left:  $d=30$  nm, right:  $d=1$  nm)

For large interparticle distances, the influence of the rounding is very small. It induces a slight decreasing and redshift of the LSPR maximum. It is only for very small interparticle distances that its influence cannot be neglected (out of range of the experimental interparticle distances)

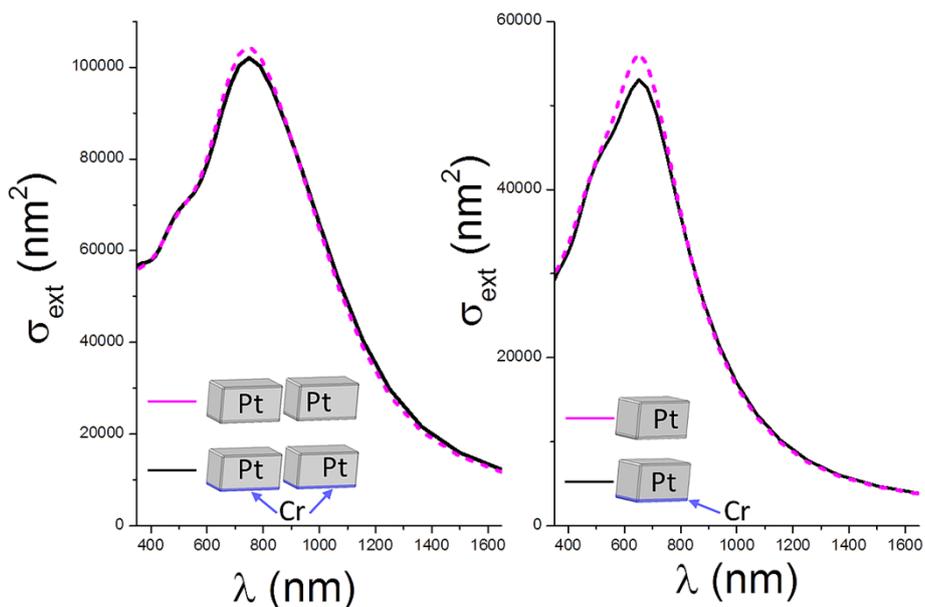


Figure S17: calculated optical extinction cross-sections of parallelepipedal dimers and monomers ( $L=125$  nm,  $w=60$  nm,  $h=70$  nm,  $r_c=3$  nm) made of pure platinum (magenta dotted line,  $h=70$  nm) or platinum over chromium (black line,  $h_{Pt}=64$  nm,  $e_{Cr}=6$  nm). For the dimers, the interparticle distance is  $d=35$  nm. The optical index of the surrounding medium is  $n_{ext}=1.35$ .

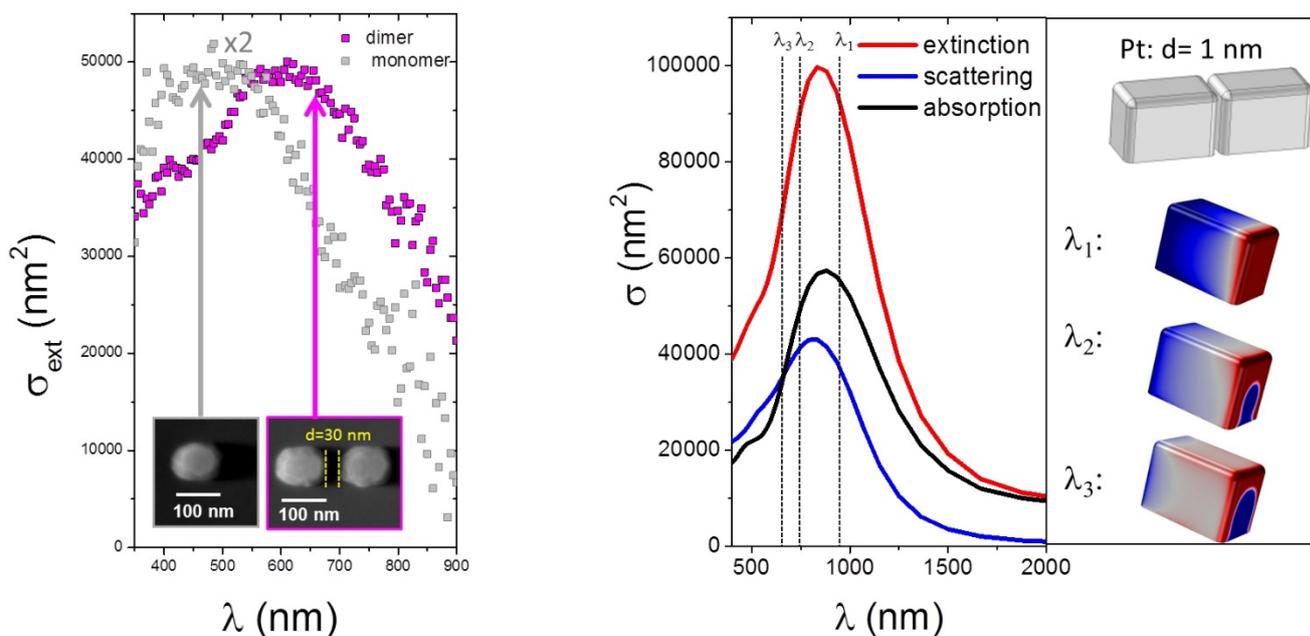


Figure S18: comparison between platinum monomer and dimer extinction cross-sections for  $L=75$  nm. The monomer cross-section (gray squares) is multiplied by 2. The shift in the LSPR of the dimer as compared to the monomer clearly evidences the coupling effect.

Figure S19: (left) Calculated extinction (red), scattering (blue) and absorption (black) cross-sections for a dimer of parallelepipeds ( $L=100$  nm,  $w=60$  nm,  $h=70$  nm,  $d=1$  nm); (right): corresponding charge distribution on one of the parallelepiped surface for the three wavelengths  $\lambda_1$ ,  $\lambda_2$ ,  $\lambda_3$  reported on the spectra.

**Notes and references**

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