

Supplement

The grand potential of the present model is expressed in eq.(2) of the manuscript as a functional of Helmholtz energy for density profile of certain species, in which the ideal gas contribution is written as

$$F^{id} [\{\rho_i\}] = k_B T \sum_i \int d\mathbf{r} \rho_i(\mathbf{r}) \left[\ln(\rho_i(\mathbf{r}) \lambda_i^3 - 1) \right] \quad (\text{S1})$$

where λ_i is the de Broglie wavelength. The direct Coulomb contribution is given by

$$F_C^{ex} [\{\rho_i\}] = \frac{1}{2} \iint d\mathbf{r}_1 d\mathbf{r}_2 \sum_{i,j} \frac{z_i z_j e^2 \rho_i(\mathbf{r}_1) \rho_j(\mathbf{r}_2)}{\epsilon \epsilon_0 |\mathbf{r}_1 - \mathbf{r}_2|} \quad (\text{S2})$$

The excluded volume term, F_{hs}^{ex} , is calculated according to MFMT,

$$\begin{aligned} F_{hs}^{ex} = k_B T \int d\mathbf{r} & \left[-n_0 \ln(1 - n_3) + \frac{n_1 n_2 - \mathbf{n}_{V1} \cdot \mathbf{n}_{V2}}{1 - n_3} \right. \\ & \left. + \frac{n_2^3 \ln(1 - n_3) - 3n_2 \mathbf{n}_{V2} \cdot \mathbf{n}_{V2} \ln(1 - n_3)}{36\pi n_3^2} + \frac{n_2^3 - 3n_2 \mathbf{n}_{V2} \cdot \mathbf{n}_{V2}}{36\pi n_3 (1 - n_3)^2} \right] \quad (\text{S3}) \end{aligned}$$

where n_i , $i = 0, 1, 2, 3$ and \mathbf{n}_{Vj} ($j = 1, 2$) are scalar and vector total weighted densities respectively.

The detail definitions for the weighted densities are given in Ref.17. F_{ele}^{ex} is derived from a truncated second-order functional Taylor expansion around the corresponding bulk fluid, which is given by

$$\begin{aligned} F_{ele}^{ex} = F_{ele}^{ex} (\{\rho_i^b\}) & + \sum_i \int d\mathbf{r} \frac{\delta F_{ele}^{ex}}{\delta \rho_i(\mathbf{r})} \Delta \rho_i(\mathbf{r}) \\ & + \frac{1}{2} \sum_{i,j} \iint d\mathbf{r} d\mathbf{r}' \frac{\delta^2 F_{ele}^{ex}}{\delta \rho_i(\mathbf{r}) \delta \rho_j(\mathbf{r})} \Delta \rho_i(\mathbf{r}) \Delta \rho_j(\mathbf{r}') \quad (\text{S4}) \end{aligned}$$

where $\Delta \rho_i(\mathbf{r}) = \rho_i(\mathbf{r}) - \rho_i^b$, and $\{\rho_i^b\}$ is the bulk density of species i . The first-order and second-order direct correlation functions are defined as,

$$k_B T \Delta C_i^{(1)}(\mathbf{r}) = -\delta F_{ele}^{ex} / \delta \rho_i(\mathbf{r}) \quad (\text{S5})$$

$$k_B T \Delta C_{ij}^{(2)}(|\mathbf{r} - \mathbf{r}'|) = -\delta^2 F_{ele}^{ex} / \delta \rho_i(\mathbf{r}) \delta \rho_j(\mathbf{r}') \quad (\text{S6})$$

$\Delta C_{ij}^{(2)}(r)$ is obtained from the Ornstein–Zernike integral equation with the MSA [1,2],

$$\Delta C_{ij}^{(2)} = \frac{z_1 z_2 e^2}{\epsilon \epsilon_0} \left[\frac{2B}{d} - \left(\frac{B}{d} \right)^2 s - \frac{1}{s} \right], \quad s = |\mathbf{r} - \mathbf{r}'| \leq d \quad (\text{S7})$$

in which $B = [y^2 + y - y(1+2y)^{1/2}] / y^2$, $y^2 = (4\pi d^2 / \epsilon \epsilon_0 k_B T) \sum_i \rho_i^b z_i^2 e^2$.

- [1] Waisman, E. and J. L. Lebowitz., *J. Chem. Phys.*, 1972, **56**(6): 3086.
- [2] Waisman, E. and J. L. Lebowitz., *J. Chem. Phys.*, 1972, **56**(6): 3093.