

Supporting Information for:

**A Theoretical Chemical Bonding, Vibronic Coupling and
Magnetic Anisotropy in Linear Iron(II) Complexes With
Single Molecule Magnet Behavior**

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SI.1 The Hamiltonian $\hat{H}_{BO} + \hat{H}_{SOC} + \hat{H}_Z$ in the basis of the $|M_L = \pm 2, M_S = \pm 2, \pm 1, 0\rangle$ basis functions, derivation of eigenvalues and g-tensor expressions.

The matrix of the Born-Oppenheimer Hamiltonian plus spin-orbit coupling plus the Zeeman operator $\hat{H}_{BO} + \hat{H}_{SOC} + \hat{H}_Z$ within the basis of products of orbital (M_L) and spin (M_S) functions $|M_L = \pm 2, M_S = \pm 2, \pm 1, 0\rangle$ for the 5E ground states is:

$$\begin{array}{cccccccccc}
 | +2, +2 \rangle & | -2, +2 \rangle & | -2, -2 \rangle & | +2, -2 \rangle & | +2, +1 \rangle & | -2, +1 \rangle & | -2, -1 \rangle & | +2, -1 \rangle & | +2, 0 \rangle & | -2, 0 \rangle \\
 \left[\begin{array}{cccccccccc}
 -\zeta + 6\beta H_z & -\delta_1 - i\delta_2 & 0 & 0 & 2\beta(H_x - iH_y) & 0 & 0 & 0 & 0 & 0 \\
 -\delta_1 + i\delta_2 & \zeta + 2\beta H_z & 0 & 0 & 0 & 2\beta(H_x - iH_y) & 0 & 0 & 0 & 0 \\
 0 & 0 & -\zeta - 6\beta H_z & -\delta_1 + i\delta_2 & 0 & 0 & 2\beta(H_x + iH_y) & 0 & 0 & 0 \\
 0 & 0 & -\delta_1 - i\delta_2 & -\zeta - 2\beta H_z & 0 & 0 & 0 & 2\beta(H_x + iH_y) & 0 & 0 \\
 2\beta(H_x + iH_y) & 0 & 0 & 0 & -\frac{\zeta}{2} + 4\beta H_z & -\delta_1 - i\delta_2 & 0 & 0 & \sqrt{\frac{3}{2}}\beta(H_x - iH_y) & 0 \\
 0 & 2\beta(H_x + iH_y) & 0 & 0 & -\delta_1 + i\delta_2 & \frac{\zeta}{2} & 0 & 0 & 0 & \sqrt{\frac{3}{2}}\beta(H_x - iH_y) \\
 0 & 0 & 2\beta(H_x - iH_y) & 0 & 0 & 0 & -\frac{\zeta}{2} - 4\beta H_z & -\delta_1 + i\delta_2 & 0 & \sqrt{\frac{3}{2}}\beta(H_x + iH_y) \\
 0 & 0 & 0 & 2\beta(H_x - iH_y) & 0 & 0 & -\delta_1 - i\delta_2 & \frac{\zeta}{2} & \sqrt{\frac{3}{2}}\beta(H_x + iH_y) & 0 \\
 0 & 0 & 0 & 0 & \sqrt{\frac{3}{2}}\beta(H_x + iH_y) & 0 & 0 & \sqrt{\frac{3}{2}}\beta(H_x - iH_y) & 2\beta H_z & -\delta_1 - i\delta_2 \\
 0 & 0 & 0 & 0 & 0 & \sqrt{\frac{3}{2}}\beta(H_x + iH_y) & \sqrt{\frac{3}{2}}\beta(H_x - iH_y) & 0 & -\delta_1 + i\delta_2 & -2\beta H_z
 \end{array} \right]
 \end{array}$$

(S1.1)

Without a magnetic field this matrix decomposes into five 2x2 matrices. For example focusing on the states $|2, 2\rangle$ and $|-2, 2\rangle$ we have :

$$\begin{bmatrix} -\zeta & -\delta_1 - i\delta_2 \\ -\delta_1 + i\delta_2 & \zeta \end{bmatrix}$$

Diagonalization yields:

$$E_{\pm} = \pm\sqrt{\zeta^2 + \delta^2}$$

In a similar way eigenvalues for the other pairs of states are obtained. They are listed in Table S1. Taking the component $|\pm 2, \pm 2\rangle$ of the ground state magnetic pair we get the following matrices of relevance for the calculation of the g-tensor components g_z and g_x , respectively ($g_y = g_x$);

g_z :

$$\begin{array}{cccccc} |2, 2\rangle & & |-2, 2\rangle & & |2, 2\rangle & & |-2, 2\rangle \\ \begin{bmatrix} -\zeta + 6\beta H_z & -\delta_1 - i\delta_2 \\ -\delta_1 + i\delta_2 & \zeta + 2\beta H_z \end{bmatrix} & = & 4\beta H_z \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} & + & \begin{bmatrix} -\zeta + 2\beta H_z & -\delta_1 - i\delta_2 \\ -\delta_1 + i\delta_2 & -\zeta - 2\beta H_z \end{bmatrix} & & \end{array} \quad (\text{S1.2})$$

$g_y = g_x$:

$$\begin{array}{cccccc} |2, 2\rangle & & |2, 1\rangle & & |2, 2\rangle & & |2, 1\rangle \\ \begin{bmatrix} -\zeta & 2\beta H_x \\ 2\beta H_x & -\frac{\zeta}{2} \end{bmatrix} & = & -(3/4)\zeta \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} & + & \begin{bmatrix} -\frac{\zeta}{4} & 2\beta H_x \\ 2\beta H_x & \frac{\zeta}{4} \end{bmatrix} & & \end{array} \quad (\text{S1.3})$$

Matrices S1.2 and S1.3 are represented in the form of a sum of a diagonal term and a baricentered term and lead to the following expressions for the dependence of the ground state eigenvalue on the H_z and H_x magnetic field directions:

$$E(H_z) = 4\beta H_z - \sqrt{(\zeta - 2\beta H_z)^2 + \delta^2} \quad (\text{S1.4})$$

$$E(H_x) = -(3/4)\zeta - \sqrt{\frac{\zeta^2}{16} + 4\beta^2 H_x^2} \quad (\text{S1.5})$$

where $\delta^2 = \delta_1^2 + \delta_2^2$ (see definitions for δ_1 and δ_2 in Section II.2). Following the definitions of g_z and g_x as the first derivatives of the energy with respect to the respective components of the magnetic field:

$$g_z = \frac{2}{\beta} (\partial E / H_z)_{H_z=0} \quad (\text{S1.6})$$

$$g_x = \frac{2}{\beta} (\partial E / H_x)_{H_x=0} \quad (\text{S1.7})$$

appropriate to a $s=1/2$ pseudospin-definition for the g -tensor. Using eqs. S1.6 and S1.7 and differentiating we obtain:

$$g_z = 8 + \frac{4\zeta}{\sqrt{\zeta^2 + \delta^2}} \quad (\text{S1.8})$$

$$g_x=0 \quad (\text{S1.9})$$

Following the same recipe expressions for the g -tensors of the other states are obtained and listed in Table S1.

Table S1. Expressions for the Energies and the Anisotropic g_z -tensor Component of the Magnetic Sublevels split out from 5E Ground State Energies by Spin-Orbit Coupling^a and Off-Axial Ligand Field.^b

$ M_L, M_S\rangle$	$M_J=M_L+M_S$	Energy	g_z^c
E, $ \pm 2, \pm 2\rangle$	± 4	$-\sqrt{\zeta^2 + \delta^2}$	$\frac{4\zeta}{\sqrt{\zeta^2 + \delta^2}} + 8$
A ₁ ,A ₂ $ \pm 2, \pm 1\rangle$	± 3	$-\sqrt{\frac{\zeta^2}{4} + \delta^2}$	$\frac{2\zeta}{\sqrt{\frac{\zeta^2}{4} + \delta^2}} + 4$
E $ \pm 2, 0\rangle$	± 2	$\pm \delta$	0
E $ \pm 2, \mp 1\rangle$	± 1	$\sqrt{\frac{\zeta^2}{4} + \delta^2}$	$-\frac{2\zeta}{\sqrt{\frac{\zeta^2}{4} + \delta^2}} + 4$
A ₁ ,A ₂ $ \pm 2, \mp 2\rangle$	0	$\sqrt{\zeta^2 + \delta^2}$	$-\frac{4\zeta}{\sqrt{\zeta^2 + \delta^2}} + 8$

^a quantified by the spin-orbit coupling constant ζ ; ^b quantified by the parameter $\delta = \sqrt{\delta_1^2 + \delta_2^2}$, see eqs.3-4 for definitions of δ_1 and δ_2 ; ^c g -tensor values are defined for pseudo spins $\pm 1/2$; non-listed values for $g_{x,y}$ are close to zero ($\delta \ll \zeta$) or $8(\delta \gg \zeta)$.

SI.2 Renner-Teller effect and solution of the vibronic energy eigenvalue problem.

The expression for the ground state potential energy for the ${}^5E\otimes\epsilon$ RT effect is formally equivalent to quadratic ${}^5E\otimes\epsilon$ Jahn-Teller coupling:⁶³⁻⁶⁹

$$\begin{array}{cccc} |{}^5E(d_{x^2-y^2})\rangle & |{}^5E(d_{xy})\rangle & |{}^5E(d_{x^2-y^2})\rangle & |{}^5E(d_{xy})\rangle \\ \hat{H}_{RT} + \hat{H}_{vib} = & \begin{bmatrix} \frac{1}{2}g(Q_x^2 - Q_y^2) & gQ_xQ_y \\ gQ_xQ_y & -\frac{1}{2}g(Q_x^2 - Q_y^2) \end{bmatrix} + & \begin{bmatrix} \frac{1}{2}K_\epsilon(Q_x^2 + Q_y^2) & 0 \\ 0 & \frac{1}{2}K_\epsilon(Q_x^2 + Q_y^2) \end{bmatrix} & \end{array} \quad (S2.1)$$

with g and K_ϵ –representing the quadratic RT and force constants, respectively, and Q_x and Q_y –the components of the ϵ -normal mode. This splits the ${}^5E(d_{x^2-y^2,xy})^1$ ground state surface into two sheets that are described by the simple expression:

$$E_{\pm} = \frac{1}{2}(K_\epsilon \pm g)(Q_x^2 + Q_y^2) \quad (S2.2)$$

Depending on whether $K_\epsilon > g$ or $K_\epsilon < g$, two different situations are encountered (Figure 2). In the first case, no new minima in the adiabatic ground state potential energy surface (APES) arise; however, the adiabatic approximation is violated by the vibronic mixing (via the term gQ_xQ_y) between the two sublevels. The other extreme, $K_\epsilon < g$, leads to a negative curvature for the lower surface, and therefore to dynamical instability. As a consequence, new minima on the APES corresponding to bent structures arise. To determine the distorted geometries and energy stabilizations at these new minima, higher-order vibronic terms need to be introduced. Traditionally this is most easily accomplished by introducing one diagonal fourth-order term which leads to eq.5 in the text. In eq.S2.3 the polar coordinate representations

$$E_{\pm} = \frac{1}{2}K_\epsilon(Q_x^2 + Q_y^2) + j(Q_x^2 + Q_y^2)^2 \pm \frac{g}{2}(Q_x^2 + Q_y^2)^2 = \frac{1}{2}(K_\epsilon \pm g)\rho^2 + j\rho^4 \quad (S2.3)$$

of Q_x and Q_y (eq.6) have been introduced; here ρ is the RT radius (eq. 7), while φ defines the direction of the distortion. Being cyclic, the coordinate φ disappears in the energy expressions (eq.S2.3, last term). Minimization of the ground state energy E . (eq.S2.3) with respect to ρ yields expressions for ρ_{min} , the stabilization energy E_{RT} , and the energy of the vertical Franck-Condon transition E_{FC} from the lower to the upper sheet of the APES at the position of the minima in terms of the parameters of the model. In the

present work, these parameters were adjusted by utilizing the following expressions:

$$\rho_{\min} = \sqrt{\frac{(g - K_{\varepsilon})}{4j}} \quad g > K_{\varepsilon} \quad (\text{S2.4})$$

$$E_{RT} = -E_{-}(\rho = \rho_{\min}) = \frac{1}{16} \frac{(g - K_{\varepsilon})^2}{j} \quad (\text{S2.5})$$

$$E_{FC} = E_{+}(\rho = \rho_{\min}) - E_{-}(\rho = \rho_{\min}) = g\rho_{\min} = g \frac{(g - K_{\varepsilon})}{4j} \quad (\text{S2.6})$$

Temperature-dependent isothermal magnetic susceptibilities and magnetization curves were calculated using a Hamiltonian of the form:

$$\hat{H} = \hat{H}_{vib} + \hat{H}_{SOC} + \hat{H}_{RT} + \hat{H}_Z \quad (\text{S2.7})$$

where, in addition to the Harmonic force term \hat{H}_{vib} (eq.S2.8) the kinetic energy operator

$\frac{1}{2} \hbar \omega_{\varepsilon} (\hat{P}_x'^2 + \hat{P}_y'^2)$ has been introduced:

$$\hat{H}_{vib} = \frac{1}{2} \hbar \omega_{\varepsilon} (\hat{P}_x'^2 + \hat{P}_y'^2 + \hat{Q}_x'^2 + \hat{Q}_y'^2) \quad (\text{S2.8})$$

Here, $\hbar \omega_{\varepsilon}$ is the energy of the two-dimensional harmonic oscillator and \hat{P}_i' and \hat{Q}_i' are dimensionless operators related to the observables for momentum and position by:

$$\hat{P}_i' = \frac{1}{\sqrt{\mu \hbar \omega_{\varepsilon}}} \hat{P}_i; \hat{Q}_i' = \sqrt{\frac{\mu \omega_{\varepsilon}}{\hbar}} \hat{Q}_i; i = x, y \quad (\text{S2.9})$$

The Hamiltonian of eq.S2.7 has been set up in the basis of the products of electronic wavefunctions $|\phi_i\rangle = |M_L = \pm 2, M_S = \pm 2, \pm 1, 0\rangle, i = 1:10$ and harmonic vibrational $|\chi_j(Q_x')\rangle, |\chi_k(Q_y')\rangle$ wavefunctions up to level n_{vib} .

$$|\Psi\rangle = \sum_{i=1}^{10} \sum_j \sum_k \sum_{j+k=0}^{n_{vib}} c_{ijk} |\phi_i \chi_j(Q_x) \chi_k(Q_y)\rangle \quad (\text{S2.10})$$

The total basis size N_v , without exploiting symmetries of vibronic wavefunctions, is given by:

$$N_v = 10 \left[\frac{n_{vib}(n_{vib} + 1)}{2} \right] \quad (\text{S2.11})$$

With the parameter set $j=8.2$, $g=461$ and $\hbar\omega_\varepsilon = 131 \text{ cm}^{-1}$, the five lowest thermally accessible levels are reproduced with good accuracy by taking vibrational levels up $n_{vib}=16$ into account (corresponding to a $N_v \times N_v = 1360 \times 1360$ matrix representation of \hat{H}). Isothermal magnetic susceptibilities and magnetizations were calculated employing the representation of \hat{H}_Z in the basis $|\Psi\rangle$ (eq.S2.10) and making explicit use of the calculated eigenvectors following procedures outlined previously.⁴⁷

Vibronic coupling in the presence of static distortions was taken into account following the theory of Ham.⁸⁰⁻⁸¹ To this end, an extra term, the strain Hamiltonian H_{Strain} was added to the Hamiltonian of eq.S2.7. We approximate this term with the LF matrix for the 25 quintet states (including the spin variables) deduced from the 1-to-1 mapping of the NEVPT2 results for the quintet states manifold onto ligand field theory:

$$(|M_L=2\rangle |M_L=1\rangle |M_L=0\rangle |M_L=-1\rangle |M_L=-2\rangle) \otimes |S=2, M_S=\pm 2, \pm 1, 0\rangle$$

$$H_{Strain} = \begin{bmatrix} -3534 & 325 & 1602 & 304 & 475 \\ 325 & 1196 & 700 & 3162 & 304 \\ 1602 & 700 & 4676 & 700 & 1602 \\ 304 & 3162 & 700 & 1196 & 325 \\ 475 & 304 & 1602 & 325 & -3534 \end{bmatrix} \otimes \begin{bmatrix} 1 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 1 \end{bmatrix} \quad (\text{S2.12})$$

This resulted in the (3400 × 3400) matrix of the Hamiltonian (including the effect of the magnetic field via the matrix of the Zeeman operator) which was solved in the same way as the Hamiltonian from eq.S2.7. In these calculations we have restricted vibronic activity to the ⁵E ground state and considered vibronic basis functions in the other three excited S=2 states as simple products of electronic and vibrational functions (eq.15).

SI.3. Angular Overlap Model expressions for C_{2v} distorted FeX_2 complexes.

In the angular overlap model the general matrix elements of the 5×5 LF one electron matrix for d-electrons is expressed in terms of parameters defined for well aligned d-orbitals of σ and π -type, e_λ ($\lambda=\sigma, \pi c, \pi s$) multiplied by factors $F_{\lambda\alpha}(\theta, \varphi, \psi)$ that depend solely on the angular geometry of the ligands (l, described by the Euler angles (θ, φ, ψ)).

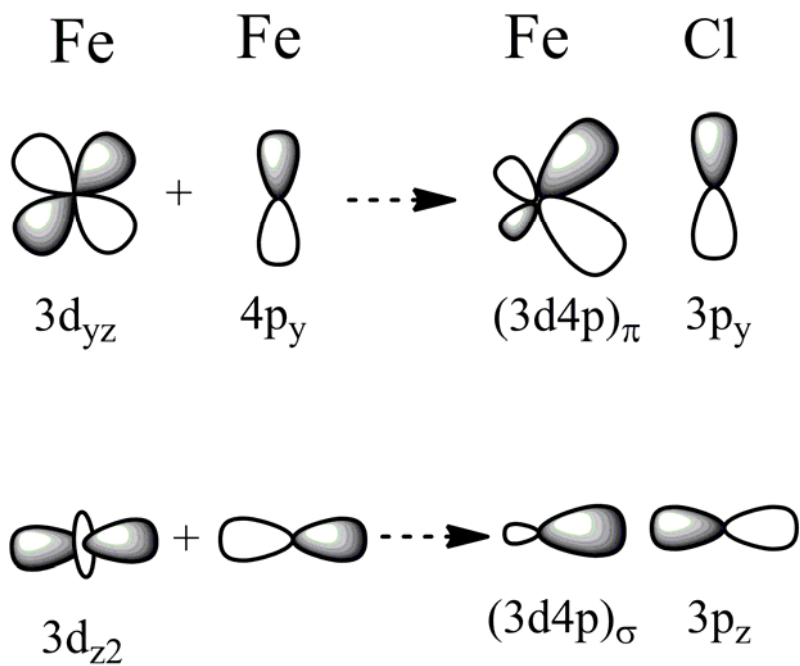
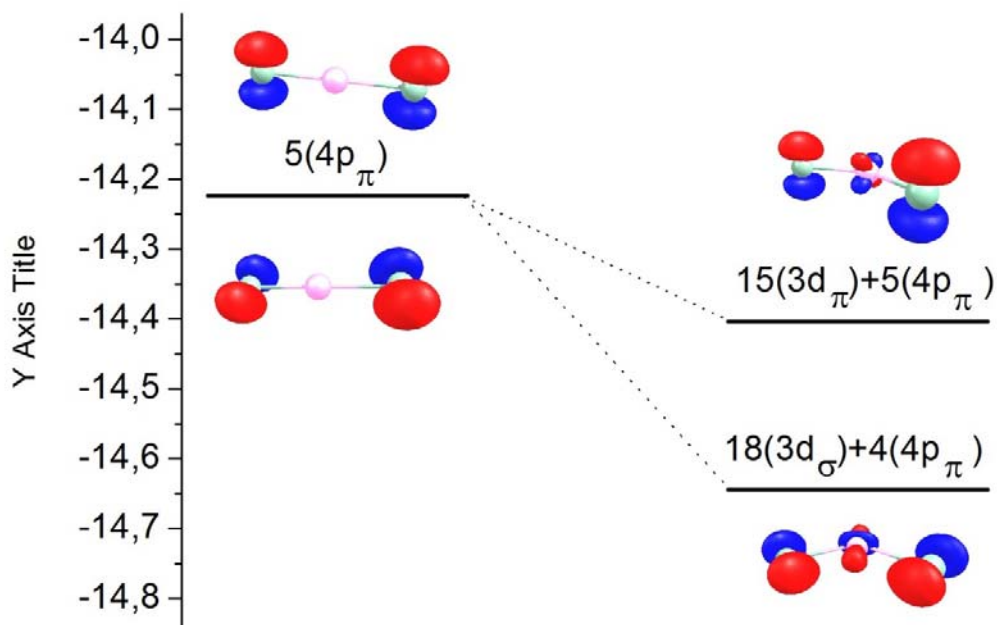
$$\mathbf{H}^{AOM}(p_\lambda) = \{H_{ij}^{AOM} = \sum_L \sum_\lambda F_{\lambda i}(\theta_l, \varphi_l, \psi_l) F_{\lambda j}(\theta_l, \varphi_l, \psi_l) e_{\lambda, l}; \lambda = \sigma, \pi c, \pi s\} \quad (S3.1)$$

The derived AOM energy expressions for a FeX_2 complex with C_{2v} symmetry (with a choice of Cartesian axes as shown in Figure 3) are given in eqs.S3.2.

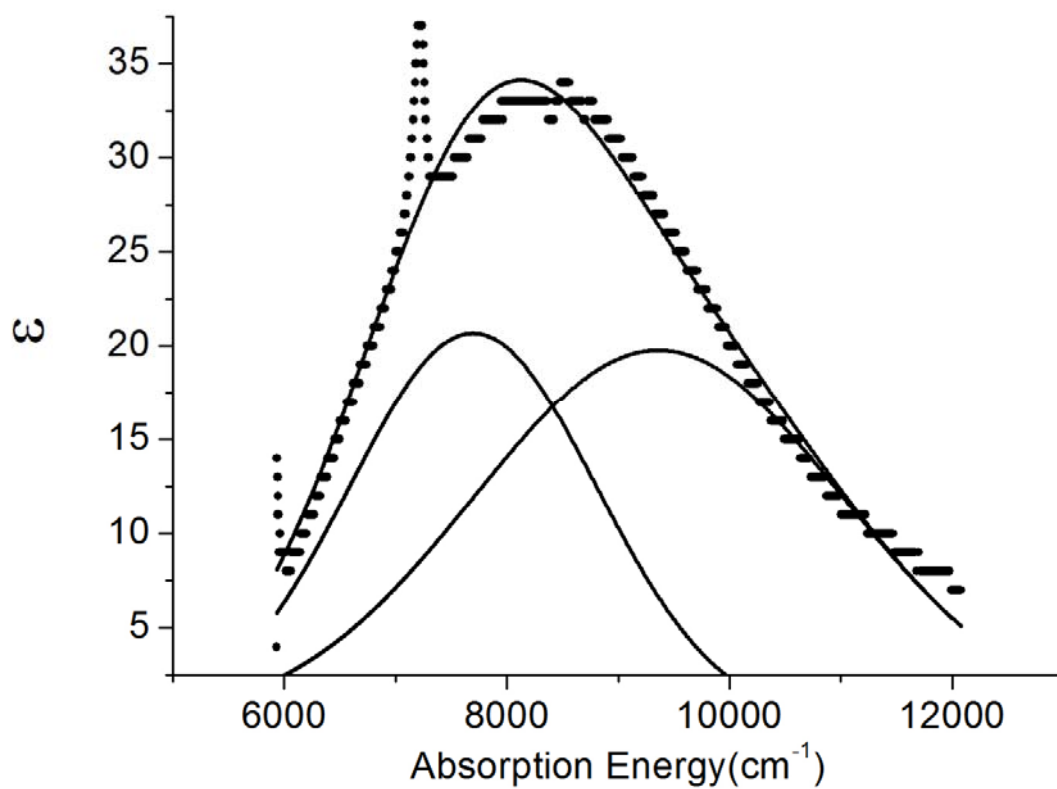
$$\begin{array}{l} \mathbf{a}_1: \quad \left| 3d_{z^2} \right\rangle \quad \left| 3d_{x^2-y^2} \right\rangle \\ \left[\begin{array}{cc} (1/8)(1+3\cos 2\theta)^2 e_\sigma + (3/2)\sin^2 2\theta e_{\pi c} & (\sqrt{3}/4)(1+3\cos 2\theta)\sin^2 \theta e_\sigma - (\sqrt{3}/2)\sin^2 2\theta e_{\pi c} \\ (\sqrt{3}/4)(1+3\cos 2\theta)\sin^2 \theta e_\sigma - (\sqrt{3}/2)\sin^2 2\theta e_{\pi c} & (3/2)\sin^4 \theta e_\sigma + (1/2)\sin^2 2\theta e_{\pi c} \end{array} \right] \\ \mathbf{a}_2: \left| 3d_{yz} \right\rangle \quad 2\cos^2 \theta e_{\pi s} \\ \mathbf{b}_1: \left| 3d_{xy} \right\rangle \quad 2\sin^2 \theta e_{\pi s} \\ \mathbf{b}_2: \left| 3d_{xz} \right\rangle \quad (3/2)\sin^2 2\theta e_\sigma + 2\cos^2 2\theta e_{\pi c} \end{array} \quad (S3.2)$$

Substitution of θ for a linear FeX_2 complex (Figure 3) leads to eq.10.

SI.4. Pseudo-Jahn-Teller activity due to 3d-4p leading to off-centric distortion in hypothetical (d^5) FeCl_2^+ and $\text{Mn}(\text{CH}_3)_2$.



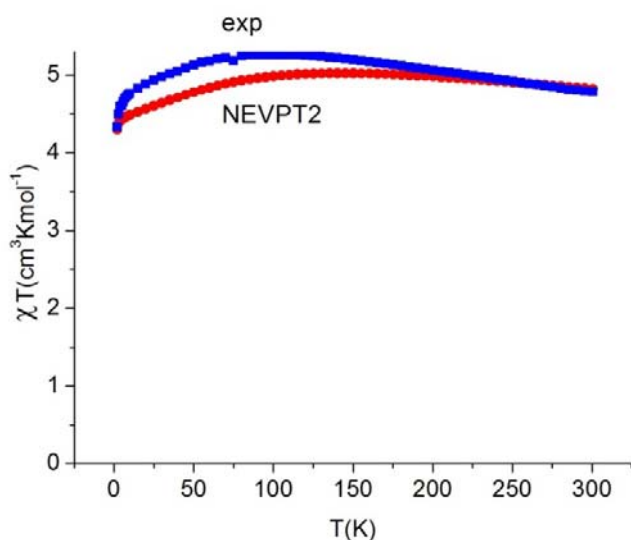
SI.5 Near infrared d-d absorption due to FeN₂ linear core in complex 1



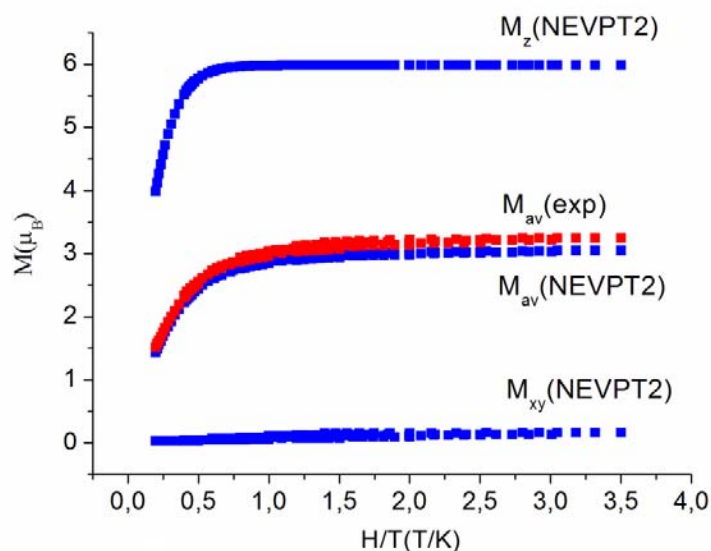
Near infrared d-d absorption due to the FeN₂ linear core in complex **1** (points) and its simulation with two Gaussian envelopes ($F=b_3 \cdot \exp\{-[(x-b_1)/b_2]^2\} + b_6 \cdot \exp\{-[(x-b_4)/b_5]^2\}$); $b_1=7696 \text{ cm}^{-1}$, $b_2=1563 \text{ cm}^{-1}$, $b_3=21$; $b_4=9360 \text{ cm}^{-1}$, $b_5=2336 \text{ cm}^{-1}$, $b_6=20$ and a standard deviation for ϵ 1.75.

SI.6. Experimental (Ref.55) and *ab initio* (NEVPT2, static geometry from x-ray data) of the molar susceptibility (χT) and molar isothermal magnetization (M) of complex 2.

Experimental (blue squares, from Ref.55) and theoretical (NEVPT2, red circles) magnetic susceptibility of complex 2 in the temperature range from 2K to 300K under a static applied dc field B of 1000 Oe:



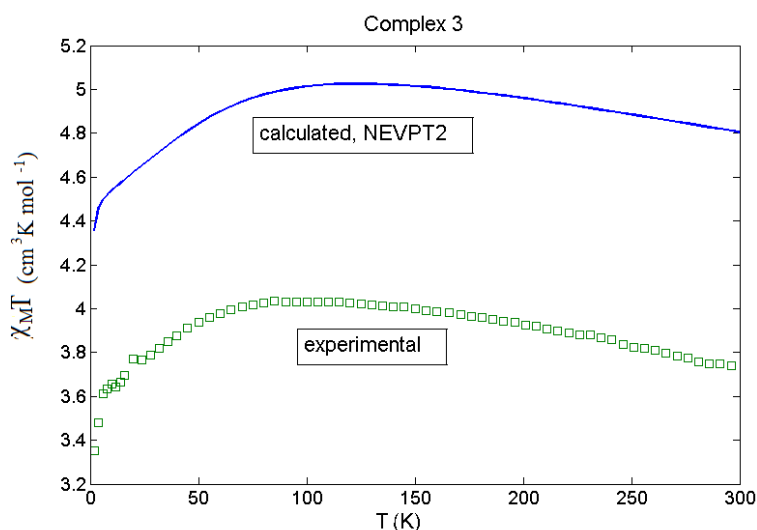
Experimental (Ref.55, red squares, from Ref.55) and calculated (spherical averaged and main values M_z and $M_{x,y}$, NEVPT2, blue squares) magnetization of complex



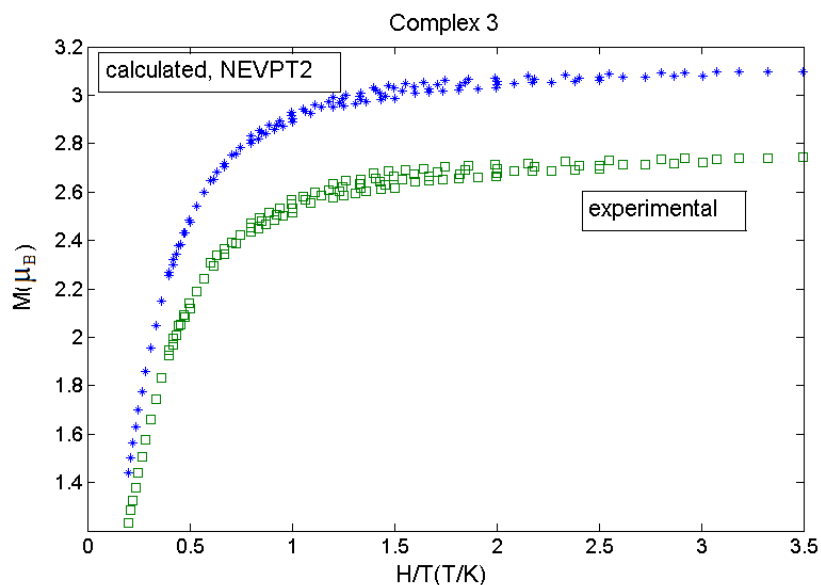
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SI.7. Experimental (Ref.55) and *ab initio* (NEVPT2, static geometry from X-ray data) of the molar susceptibility (χT) and molar isothermal magnetization (M) of complexes 3-6.

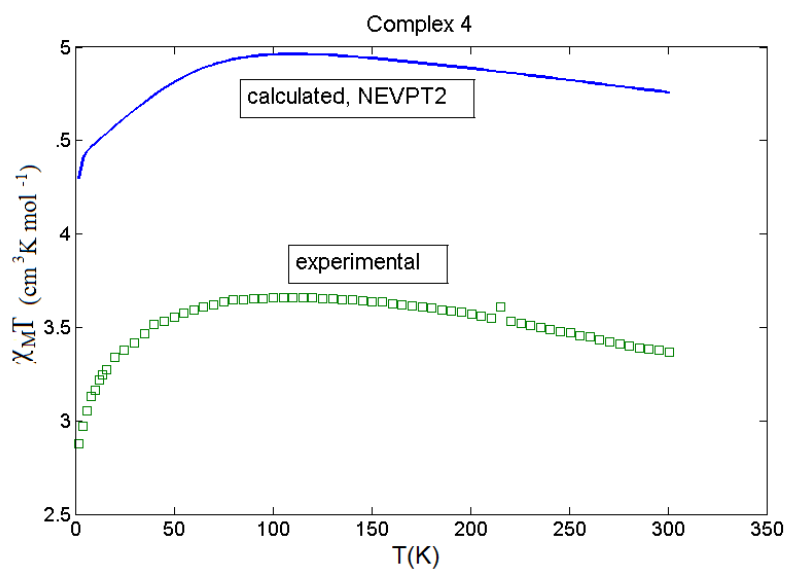
Experimental (Ref.55, green squares) and theoretical (NEVPT2, blue line) magnetic susceptibility of complex **3** in the temperature range from 2K to 300K under a static applied dc field B of 1000 Oe:



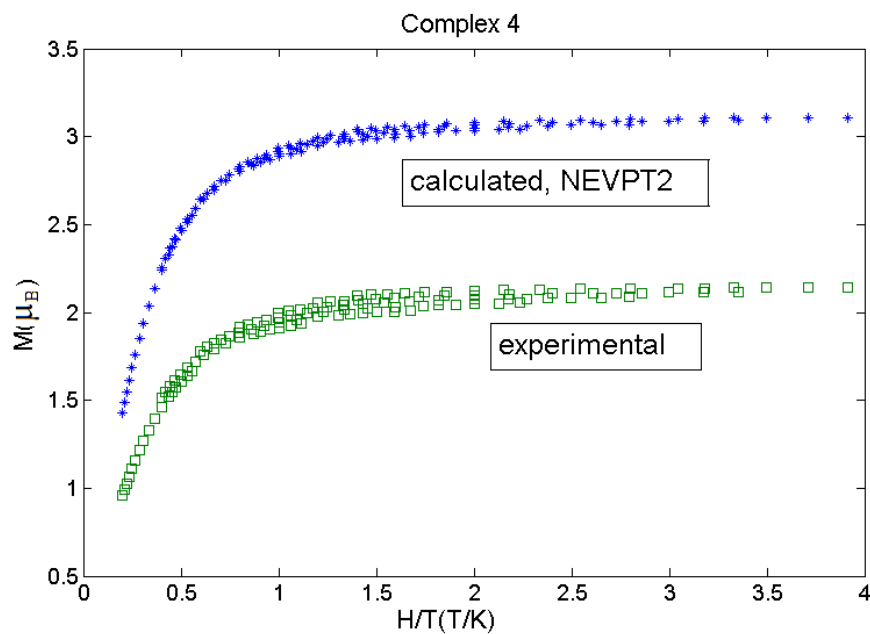
Experimental (Ref.55, green squares) and calculated (NEVPT2, blue stars) magnetization of complex **3**:



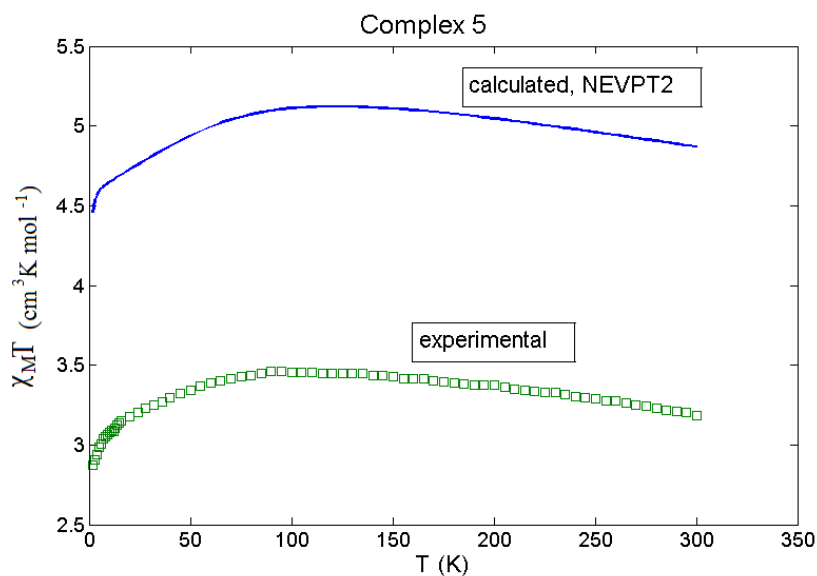
Experimental (Ref.55, green squares) and theoretical (NEVPT2, blue line) magnetic susceptibility of complex **4** in the temperature range from 2K to 300K under a static applied dc field B of 1000 Oe:



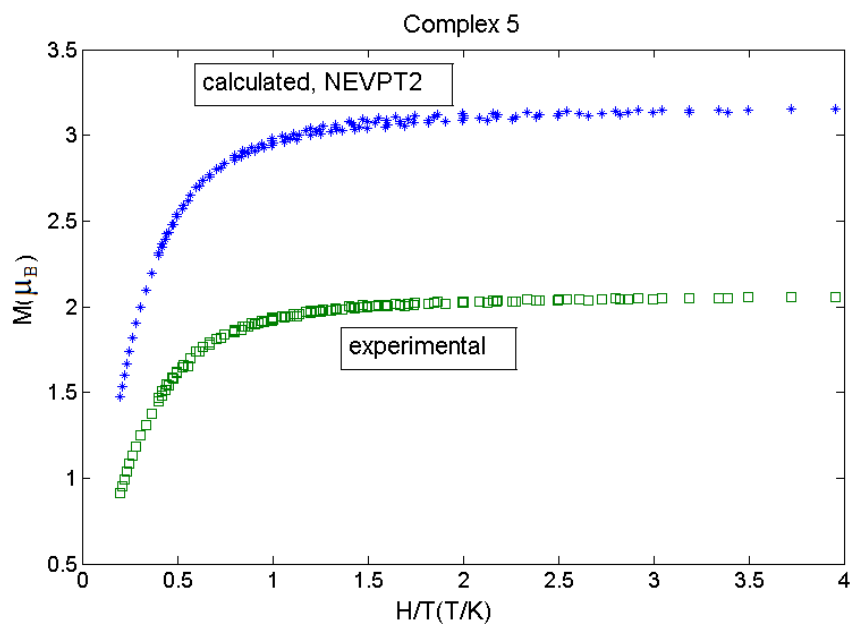
Experimental (Ref.55, green squares) and calculated (NEVPT2, blue stars) magnetization of complex **4**:



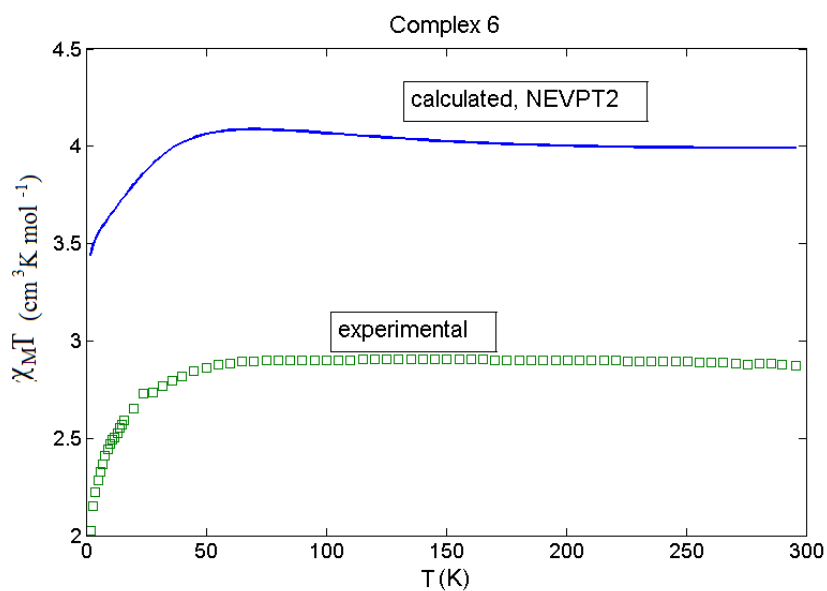
Experimental (Ref.55, green squares) and theoretical (NEVPT2, blue line) magnetic susceptibility of complex **5** in the temperature range from 2K to 300K under a static applied dc field B of 1000 Oe:



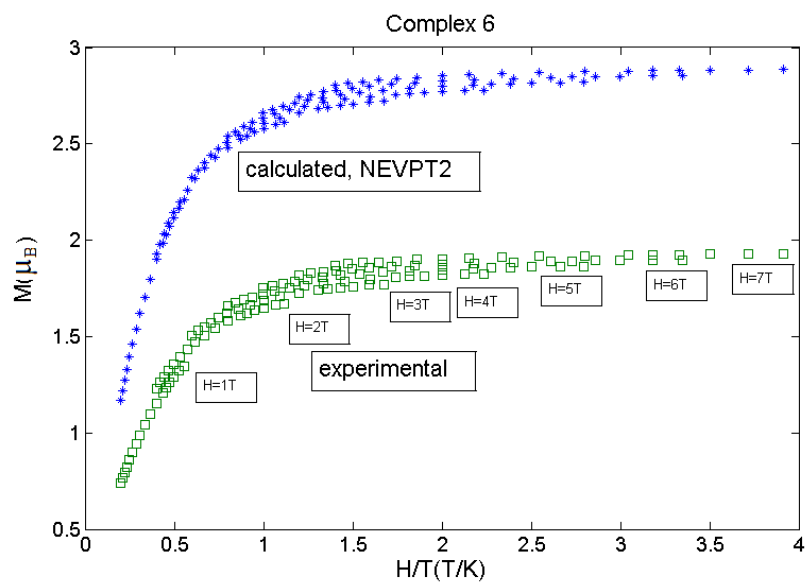
Experimental (Ref.55, green squares) and calculated (NEVPT2, blue stars) magnetization of complex **5**:



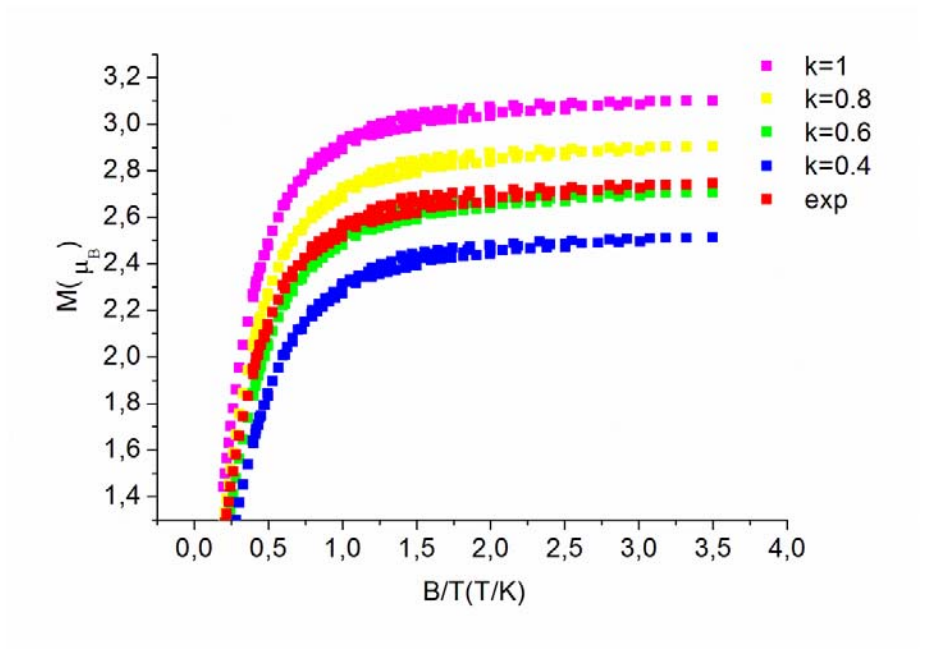
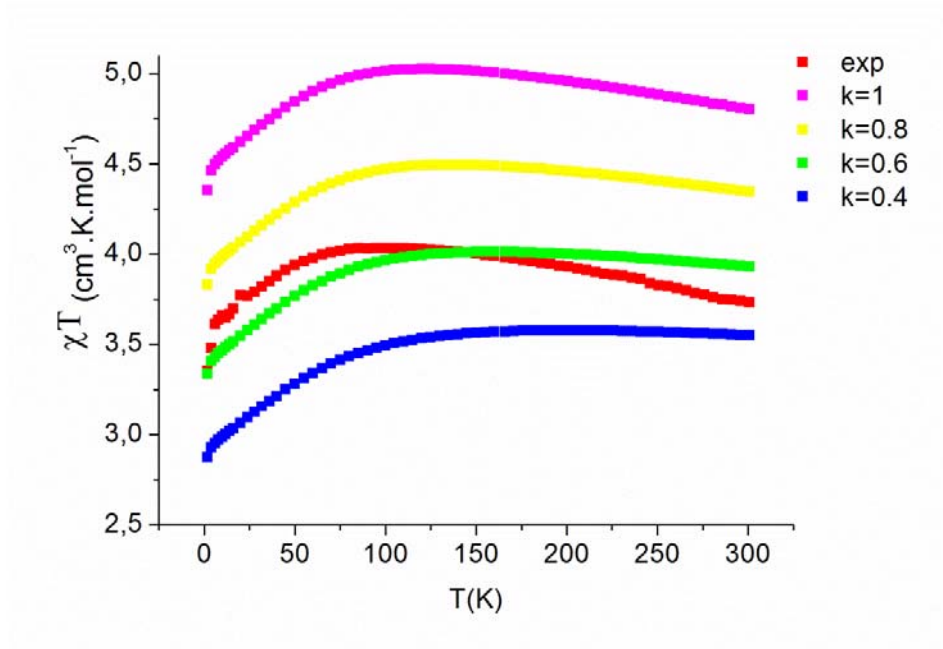
Experimental (Ref.55, green squares) and theoretical (NEVPT2, blue line) magnetic susceptibility of complex **6** in the temperature range from 2K to 300K under a static applied dc field B of 1000 Oe:



Experimental (Ref.55, green squares) and calculated (NEVPT2, blue stars) magnetization of complex **6**:

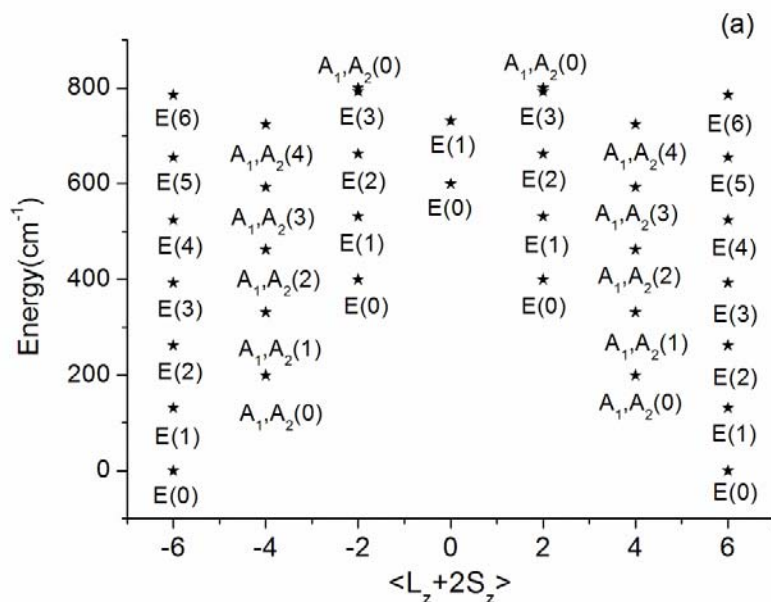


SI.8 Effect of the Stevens orbital reduction factor k introduced into the Zeeman operator $H_z = \beta H(2S + kL)$ on the magnetic susceptibility (top) and the isothermal magnetization for complex 3 (bottom)

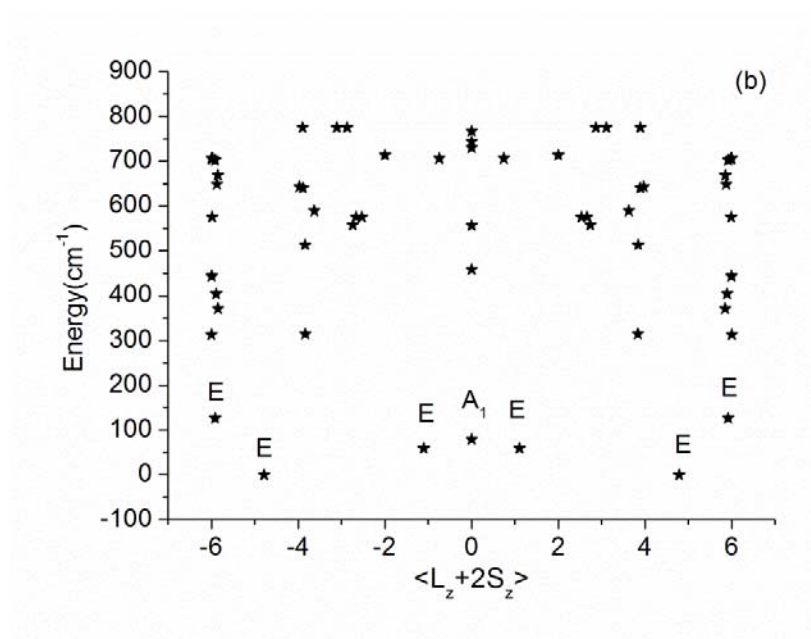


SI.9 Vibronic energy level diagrams for the limiting case of zero vibronic coupling $f=0$, static limit and $f=0.1$

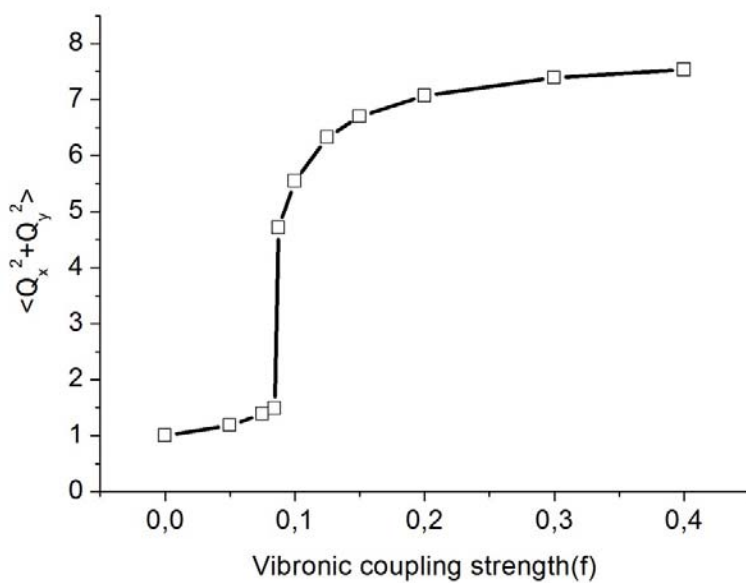
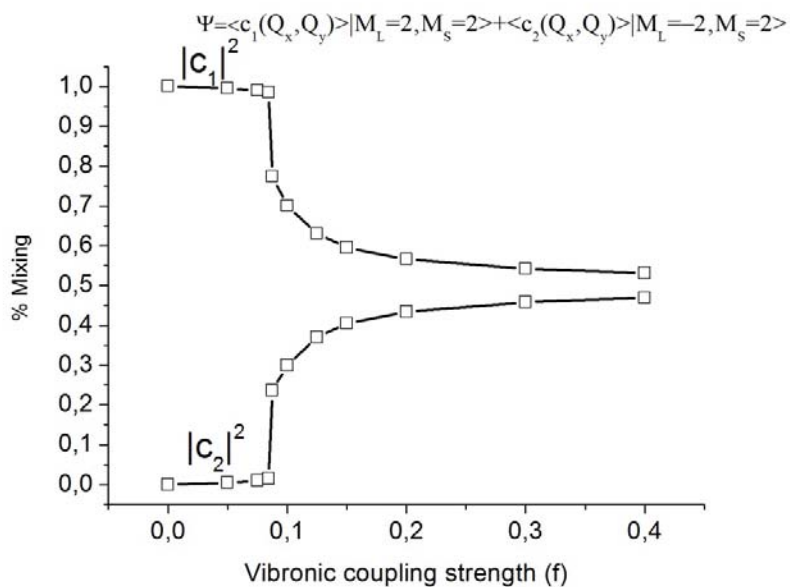
Combined electronic - vibrational energy level diagram for the limiting case of zero vibronic coupling $f=0$, static limit, (a) and $f=0.1$



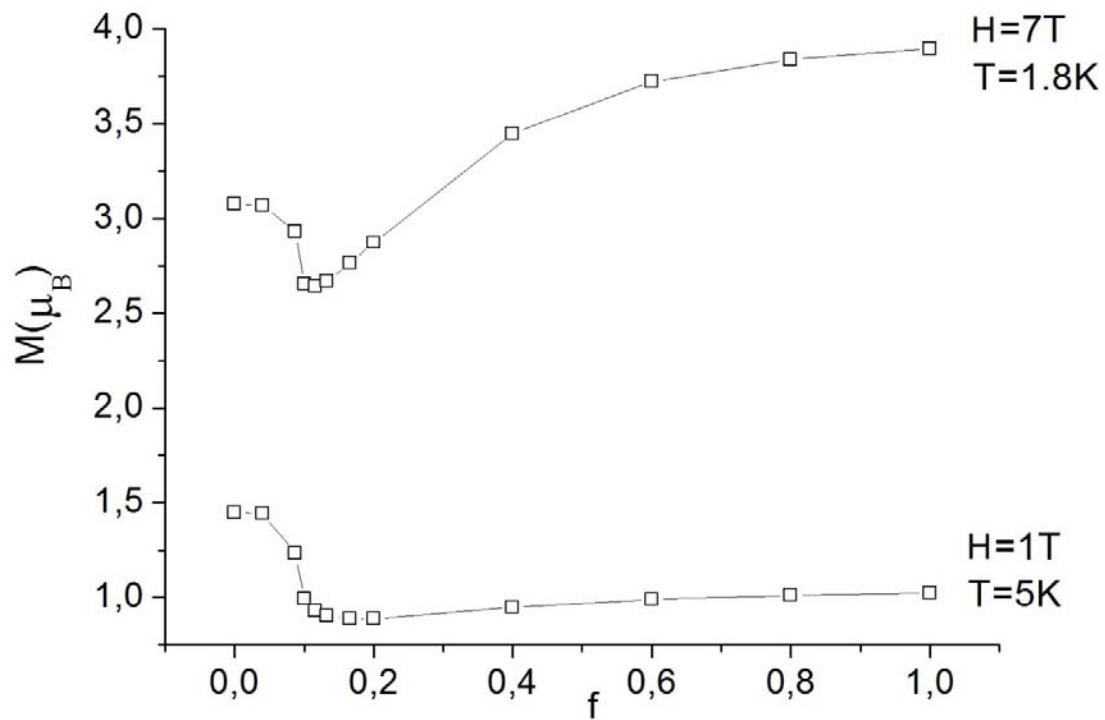
(b):



SI.10 Vibronically driven mixing between $|M_L=2, M_S=2\rangle$ and $|M_L=-2, M_S=2\rangle$ total angular momentum states (top) and large amplitude motions (bottom):

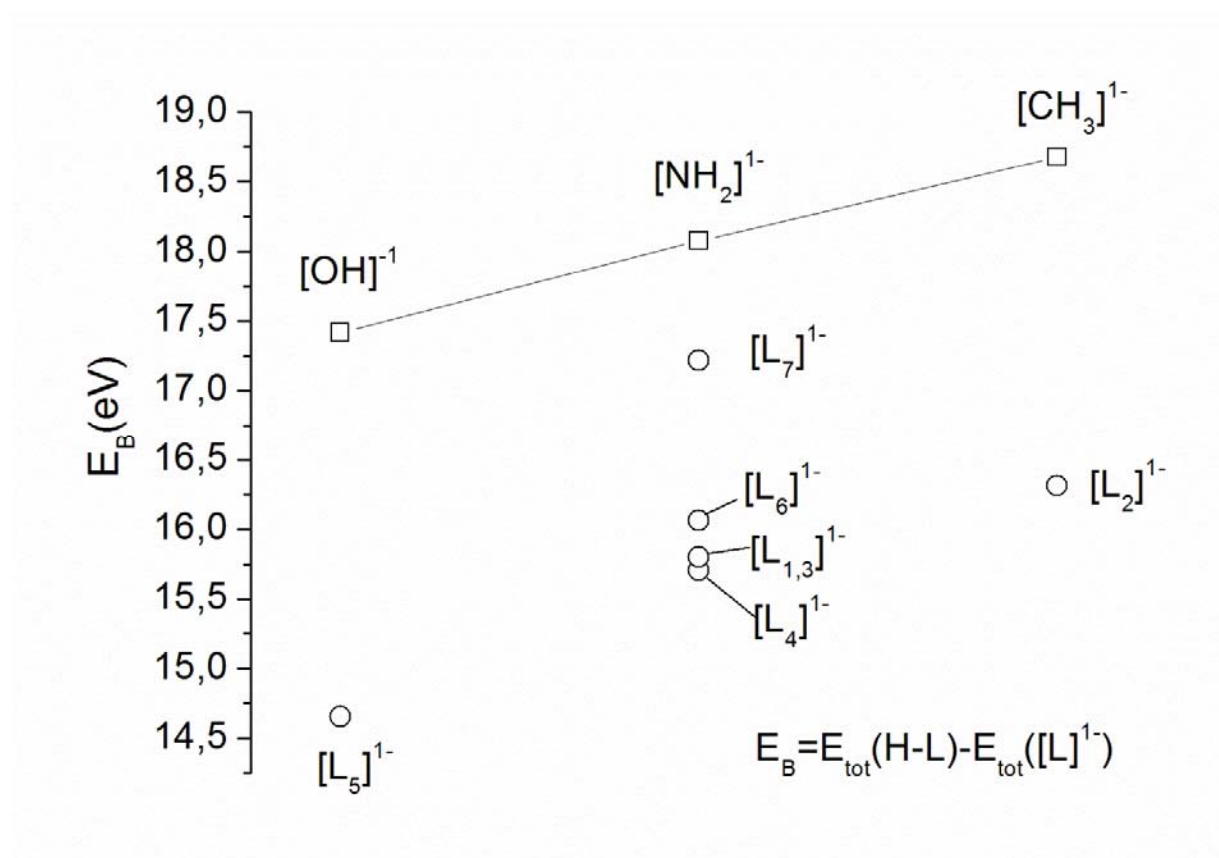


SI.11. Dependence of the magnetization on the vibronic coupling strength f .



SI.12. Variations of the Lewis basicity with the ligand substituents.

The affinity of OH^- , NH_2^- and CH_3^- , toward the proton as quantified by the energy E_B defined in the figure below reflects the expected increase of the Lewis basicity of these anions across this series. For comparison in this figure we also include the corresponding values of E_B for the ligand of complexes **1-7**. The comparison reflects a dramatic change of the Lewis basicity –an overall decrease with respect to the CH_3^- , NH_2^- and OH^- and additionally variations with the substituents (for the N-ligands).



SI.13. Complete list of LF matrices $\mathbf{H}_{\text{NEVPT2}}^{\text{LFT}}(S=2, M_s=2)$ for complexes 1-6 from a 1-to-1 mapping to NEVPT2 results

LF matrices for the $S=2, M_s=2$ states:

Complex 1

$$\mathbf{H}_{\text{NEVPT2}}^{\text{LFT}}(S=2, M_s=2) = \begin{array}{c} \begin{array}{ccccc} |d_{z^2}\rangle & |d_{xz}\rangle & |d_{yz}\rangle & |d_{x^2-y^2}\rangle & |d_{xy}\rangle \end{array} \\ \left[\begin{array}{ccccc} 2576 & -532 & -34 & 270 & 829 \\ -532 & -471 & -156 & 36 & -255 \\ -34 & -156 & 6147 & 44 & -563 \\ 270 & 36 & 44 & -4155 & 102 \\ 829 & -255 & -563 & 102 & -4097 \end{array} \right] \end{array}$$

Complex 2

$$\mathbf{H}_{\text{NEVPT2}}^{\text{LFT}}(S=2, M_s=2) = \begin{array}{c} \begin{array}{ccccc} |d_{z^2}\rangle & |d_{xz}\rangle & |d_{yz}\rangle & |d_{x^2-y^2}\rangle & |d_{xy}\rangle \end{array} \\ \left[\begin{array}{ccccc} 2130 & -43 & -4 & -9 & -116 \\ -43 & 1852 & 15 & 45 & 64 \\ -4 & 15 & 1810 & 24 & -177 \\ -9 & 45 & 24 & -2868 & 32 \\ -116 & 64 & -177 & 32 & -2924 \end{array} \right] \end{array}$$

Complex 3

$$\mathbf{H}_{\text{NEVPT2}}^{\text{LFT}}(S=2, M_s=2) = \begin{array}{c} \begin{array}{ccccc} |d_{z^2}\rangle & |d_{xz}\rangle & |d_{yz}\rangle & |d_{x^2-y^2}\rangle & |d_{xy}\rangle \end{array} \\ \left[\begin{array}{ccccc} 5107 & -137 & -4 & -897 & 690 \\ -137 & -1865 & 121 & 171 & -96 \\ -4 & 121 & 4217 & 36 & 169 \\ -897 & 171 & 36 & -3650 & -183 \\ 690 & -96 & 169 & -183 & -3809 \end{array} \right] \end{array}$$

Complex 4

$$\begin{array}{c}
 |d_{z^2}\rangle \quad |d_{xz}\rangle \quad |d_{yz}\rangle \quad |d_{x^2-y^2}\rangle \quad |d_{xy}\rangle \\
 \mathbf{H}_{\text{NEVPT2}}^{\text{LFT}}(S=2, M_s=2) = \begin{bmatrix}
 5267 & -111 & 29 & -1518 & -290 \\
 -111 & -1720 & 154 & 178 & -3 \\
 29 & 154 & 3965 & 9 & 121 \\
 -1518 & 178 & 9 & -3469 & 156 \\
 -290 & -3 & 121 & 156 & -4042
 \end{bmatrix}
 \end{array}$$

Complex 5

$$\begin{array}{c}
 |d_{z^2}\rangle \quad |d_{xz}\rangle \quad |d_{yz}\rangle \quad |d_{x^2-y^2}\rangle \quad |d_{xy}\rangle \\
 \mathbf{H}_{\text{NEVPT2}}^{\text{LFT}}(S=2, M_s=2) = \begin{bmatrix}
 6073 & 423 & 3 & -1320 & -836 \\
 423 & 442 & 432 & 153 & 90 \\
 3 & 432 & 2357 & 40 & 138 \\
 -1320 & 153 & 40 & -4366 & 220 \\
 -836 & 90 & 138 & 220 & -4506
 \end{bmatrix}
 \end{array}$$

Complex 6

$$\begin{array}{c}
 |d_{z^2}\rangle \quad |d_{xz}\rangle \quad |d_{yz}\rangle \quad |d_{x^2-y^2}\rangle \quad |d_{xy}\rangle \\
 \mathbf{H}_{\text{NEVPT2}}^{\text{LFT}}(S=2, M_s=2) = \begin{bmatrix}
 4676 & -84 & -987 & 967 & -2049 \\
 -84 & 4358 & -80 & 13 & 296 \\
 -987 & -80 & -1965 & -20 & 554 \\
 967 & 13 & -20 & -3105 & -203 \\
 -2049 & 296 & 554 & -203 & -3964
 \end{bmatrix}
 \end{array}$$

Matrices $\mathbf{H}_{\text{NEVPT2}}^{\text{LFT}}(S=2, M_s=2)$ have been written down in a form which allows a direct comparison with the corresponding AOM matrices (see SI.3). The former have been transformed into the basis the angular momentum eigenfunctions for $|L, M_L\rangle$, $L=2$, $M_L=\pm 2, \pm 1, 0$ given by the transformation matrix \mathbf{T}

$$\begin{bmatrix} |2,2\rangle \\ |2,1\rangle \\ |2,0\rangle \\ |2,-1\rangle \\ |2,-2\rangle \end{bmatrix} = \begin{bmatrix} 0 & 0 & 0 & \frac{1}{\sqrt{2}} & \frac{1}{i\sqrt{2}} \\ 0 & -\frac{1}{\sqrt{2}} & \frac{i}{\sqrt{2}} & 0 & 0 \\ 1 & 0 & 0 & 0 & 0 \\ 0 & \frac{1}{\sqrt{2}} & \frac{i}{\sqrt{2}} & 0 & 0 \\ 0 & 0 & 0 & \frac{1}{\sqrt{2}} & -\frac{1}{i\sqrt{2}} \end{bmatrix} \begin{bmatrix} |d_{z^2}\rangle \\ |d_{xz}\rangle \\ |d_{yz}\rangle \\ |d_{x^2-y^2}\rangle \\ |d_{xy}\rangle \end{bmatrix}$$

We identify these transformed matrices with the Strain energy Hamiltonian employed in calculations of vibronic levels in the presence of low-symmetry ligand field strain:

$$\mathbf{H}_{\text{Strain}} = \mathbf{T}\mathbf{H}_{\text{NEVPT2}}^{\text{LFT}}\mathbf{T}^\dagger$$

The matrices $\mathbf{H}_{\text{Strain}}$ are complex. Below we list the diagonal elements and the moduli

$\sqrt{H_{\text{Strain}}(i,j) \cdot H_{\text{Strain}}^*(i,j)}$ of the off-diagonal $H_{\text{Strain}}(i,j)$, $i \neq j$:

Complex 1

$$\begin{array}{ccccc} & |2,2\rangle & |2,1\rangle & |2,0\rangle & |2,-1\rangle & |2,-2\rangle \\ \mathbf{H}_{\text{Strain}} = & \begin{bmatrix} -4126 & 303 & 616 & 318 & 106 \\ 303 & 2838 & 377 & 3313 & 318 \\ 616 & 377 & 2576 & 377 & 616 \\ 318 & 3313 & 377 & 2838 & 303 \\ 106 & 318 & 616 & 303 & -4126 \end{bmatrix} \end{array}$$

Complex 2

$$\begin{array}{ccccc} & |2,2\rangle & |2,1\rangle & |2,0\rangle & |2,-1\rangle & |2,-2\rangle \\ \mathbf{H}_{\text{Strain}} = & \begin{bmatrix} -2896 & 69 & 82 & 119 & 43 \\ 69 & 1831 & 30 & 26 & 119 \\ 82 & 30 & 2130 & 30 & 82 \\ 119 & 26 & 30 & 1831 & 69 \\ 43 & 119 & 82 & 69 & -2896 \end{bmatrix} \end{array}$$

Complex 3

$$\begin{array}{ccccc} & |2,2\rangle & |2,1\rangle & |2,0\rangle & |2,-1\rangle & |2,-2\rangle \\ \mathbf{H}_{\text{Strain}} = & \begin{bmatrix} -3730 & 182 & 800 & 30 & 199 \\ 182 & 1176 & 97 & 3043 & 30 \\ 800 & 97 & 5107 & 97 & 800 \\ 30 & 3043 & 97 & 1176 & 182 \\ 199 & 30 & 800 & 182 & -3730 \end{bmatrix} \end{array}$$

Complex 4

$$\begin{array}{ccccc} & |2,2\rangle & |2,1\rangle & |2,0\rangle & |2,-1\rangle & |2,-2\rangle \\ \mathbf{H}_{\text{Strain}} = & \begin{bmatrix} -3756 & 149 & 1093 & 29 & 326 \\ 149 & 1122 & 81 & 2846 & 29 \\ 1093 & 81 & 5267 & 81 & 1093 \\ 29 & 2846 & 81 & 1122 & 149 \\ 326 & 29 & 1093 & 149 & -3756 \end{bmatrix} \end{array}$$

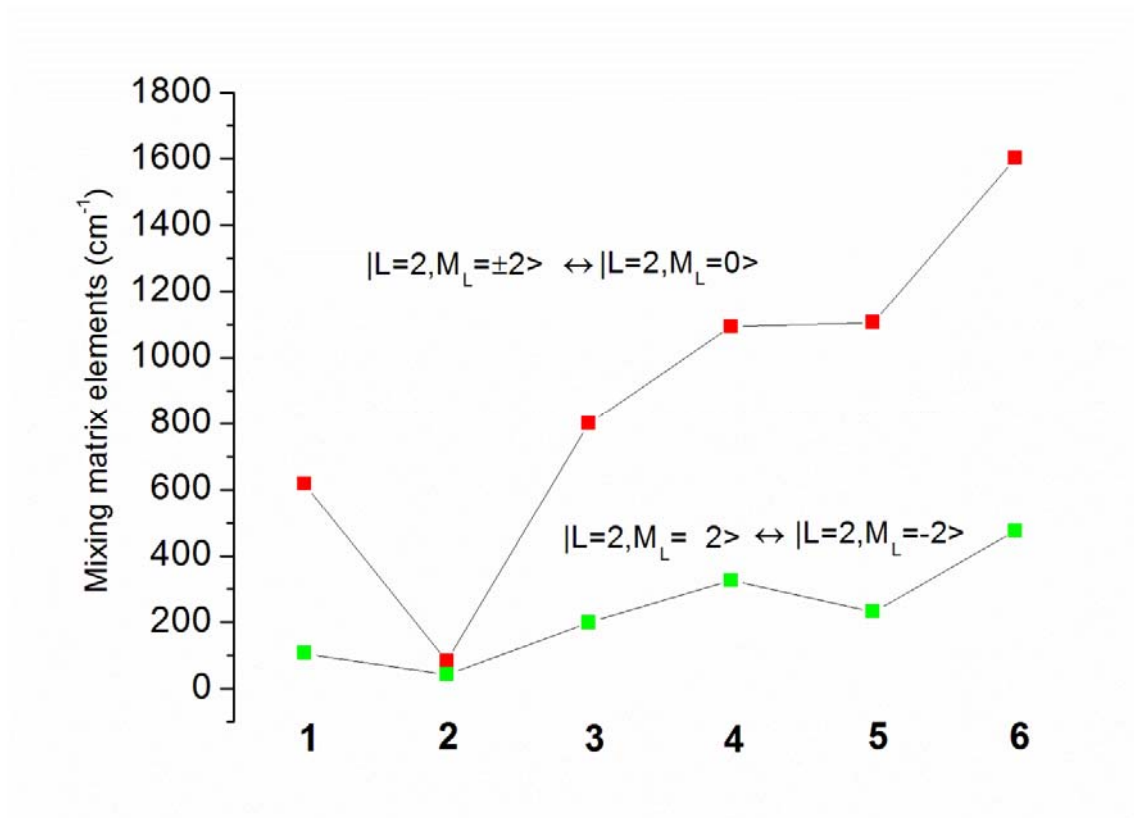
Complex 5

$$\begin{array}{ccccc} & |2,2\rangle & |2,1\rangle & |2,0\rangle & |2,-1\rangle & |2,-2\rangle \\ \mathbf{H}_{\text{Strain}} = & \begin{bmatrix} -4436 & 147 & 1105 & 66 & 231 \\ 147 & 1400 & 299 & 1051 & 66 \\ 1105 & 299 & 6073 & 299 & 1105 \\ 66 & 1051 & 299 & 1400 & 147 \\ 231 & 66 & 1105 & 147 & -4436 \end{bmatrix} \end{array}$$

Complex 6

$$\begin{array}{ccccc} & |2,2\rangle & |2,1\rangle & |2,0\rangle & |2,-1\rangle & |2,-2\rangle \\ \mathbf{H}_{\text{Strain}} = & \begin{bmatrix} -3535 & 324 & 1602 & 304 & 475 \\ 324 & 1197 & 700 & 3162 & 304 \\ 1602 & 700 & 4676 & 700 & 1602 \\ 304 & 3162 & 700 & 1197 & 324 \\ 475 & 304 & 1602 & 324 & -3535 \end{bmatrix} \end{array}$$

Matrix elements expressed in the $|L, M_L\rangle$, $L=2$, $M_L=\pm 2, \pm 1, 0$ basis are particularly suitable for an analysis of the effect of low-symmetry ligand field on the orbital moment reduction. Below we plot some matrix elements of relevance for the discussion in Section 3.2.3.



Matrix elements affecting the orbital moments of the magnetic ground states of the $\text{Fe}^{\text{II}}\text{X}_2$ SMM, responsible for the orbital moment reduction and quantum tunneling of the magnetization.

SI.14 Energies of the lowest ground state quintet to the lowest eleven triplet transitions from NEVPT2 and the spin-orbit split terms from the nominal 5E ground state calculated using full CI within the manifold of the electronic states due to the d^6 configuration of Fe(II) and deductions of B , C and ζ from a best fit to these data.

FeN2, Complex 1; B and C from a least squares fit to spin-forbidden transitions

	energy (NEVPT2)	energy(calc.)	difference	weight
31	18023.000	16942.277	-1080.723	1.000
31	18145.000	16967.543	-1177.457	1.000
31	19390.000	17898.579	-1491.421	1.000
31	20209.000	19504.469	-704.531	1.000
31	20328.000	20006.306	-321.694	1.000
31	20394.000	20747.511	353.511	1.000
31	20403.000	20770.940	367.940	1.000
31	20655.000	20883.799	228.799	1.000
31	20657.000	21413.907	756.907	1.000
31	21403.000	22556.880	1153.880	1.000
31	21421.000	23334.156	1913.156	1.000

Sigma, weighted = 1011.559

NEXT CALCULATION WITH OPTIMIZED PARAMETERS (AOM parameters from spin quintet transitions)

B = 1511.824021
C = 3245.444663
*
esig = 3351.000000
epis = 5137.000000
epic = 1828.000000

AOM matrix

	xy	xz	yz	x2y2	z2
	.000	.000	.000	.000	.000
	.000	3656.000	.000	.000	.000
	.000	.000	10274.000	.000	.000
	.000	.000	.000	.000	.000
	.000	.000	.000	.000	6702.000

zeta from a least squares fit to SOC split terms of 5E :

NEXT CALCULATION WITH OPTIMIZED PARAMETERS

zeta = 413.702538
*
esig = 3351.000000
epis = 5137.000000
epic = 1828.000000
B = 1512.000000
C = 3245.000000

2	.000	2
4	204.681	2
5	408.099	1
6	408.252	1
7	605.394	1
8	616.113	1
9	812.321	1
10	812.631	1

	energy (NEVPT2)	energy(calc.)	difference	weight
2A	189.8,190.4	204.681	14.681	1.000
1A	322.800	408.099	85.299	1.000
1A	485.500	408.252	-77.248	1.000
1A	612.800	605.394	-7.406	1.000
1A	625.000	616.113	-8.887	1.000
1A	810.100	812.321	2.221	1.000
1A	810.300	812.631	2.331	1.000

Sigma, weighted = 44.082

FeC2, Complex 2 ; B and C from a least squares fit to spin-forbidden transitions

NEXT CALCULATION WITH OPTIMIZED PARAMETERS

B = 1418.171058
C = 3563.521806
*
esig = 2513.000000
epis = 2364.000000
epic = 2364.000000

AOM matrix

	xy	xz	yz	x2y2	z2
	.000	.000	.000	.000	.000
	.000	4728.000	.000	.000	.000
	.000	.000	4728.000	.000	.000
	.000	.000	.000	.000	.000
	.000	.000	.000	.000	5026.000

	energy (NEVPT2)	energy(calc.)	difference	weight
32	21255.6,21320.2	21165.877	-122.123	1.000
32	21363.6,21475.4	21386.612	-33.388	1.000
32	21846.2, 21901.7	21408.693	-465.307	1.000
31	20622.500	21467.156	845.156	1.000
32	22185.5, 22201.5	22070.948	-123.052	1.000
32	23454.5, 23458.5	23361.753	-94.247	1.000

Sigma, weighted = 402.254

NEXT CALCULATION WITH OPTIMIZED PARAMETERS

zeta = 407.038120
*
esig = 2513.000000
epis = 2364.000000
epic = 2364.000000
B = 1418.000000
C = 3564.000000

	energy(NEVPT2)	energy(calc.)	difference	weight
2	195.2,195.3	200.043	4.843	1.000
2	353.0,444.6	400.051	1.051	1.000
2	602.0,603.8	600.084	-2.816	1.000
1	799.7	800.026	.026	1.000
1	800.0	800.350	.350	1.000

Sigma, weighted = 2.554

FeN2 Complex 3

NEXT CALCULATION WITH OPTIMIZED PARAMETERS

B = 970.316709
C = 3781.851172
*
esig = 4418.000000
epis = 3973.000000
epic = 932.000000

AOM matrix

	xy	xz	yz	x2y2	z2
	.000	.000	.000	.000	.000
	.000	1864.000	.000	.000	.000
	.000	.000	7946.000	.000	.000
	.000	.000	.000	.000	.000
	.000	.000	.000	.000	8836.000

	energy (exp.)	energy(calc.)	difference	weight
--	---------------	---------------	------------	--------

31	14944.000	16992.330	2048.330	1.000
31	15300.000	17364.524	2064.524	1.000
31	18596.000	17714.038	-881.962	1.000
31	19868.000	19209.173	-658.827	1.000
31	20031.000	19357.001	-673.999	1.000
31	20408.000	19957.752	-450.248	1.000
31	20545.000	20421.458	-123.542	1.000
31	20832.000	20422.267	-409.733	1.000
31	21284.000	20593.476	-690.524	1.000
31	21911.000	21682.768	-228.232	1.000
31	21928.000	21771.679	-156.321	1.000

Sigma, weighted = 1002.874

NEXT CALCULATION WITH OPTIMIZED PARAMETERS

zeta = 426.569658
*
esig = 4418.000000
epis = 3973.000000
epic = 932.000000
B = 970.000000
C = 3782.000000

AOM matrix

	xy	xz	yz	x2y2	z2
	.000	.000	.000	.000	.000
	.000	1864.000	.000	.000	.000
	.000	.000	7946.000	.000	.000
	.000	.000	.000	.000	.000
	.000	.000	.000	.000	8836.000

no.	Energy	Mult.(total)	Mult.(Spin)	Mult.(Bahn)
2	.000	2		
4	204.419	2		
5	410.245	1		
6	410.470	1		
7	604.452	1		
8	634.552	1		
9	831.912	1		
10	832.531	1		

	energy(NEVPT2)	energy(calc.)	difference	weight
2A	176.400,179.8	204.419	26.319	1.000
1A	286.600	410.245	123.645	1.000
1A	529.100	410.470	-118.630	1.000
1A	632.000	604.452	-27.548	1.000
1A	646.600	634.552	-12.048	1.000
1A	822.000	831.912	9.912	1.000
1A	822.500	832.531	10.031	1.000

Sigma, weighted = 66.715

Complex 4

NEXT CALCULATION WITH OPTIMIZED PARAMETERS

B = 822.514678
C = 3965.134070
*
esig = 4512.000000
epis = 3860.000000
epic = 1018.000000

AOM matrix

	xy	xz	yz	x2y2	z2
	.000	.000	.000	.000	.000
	.000	2036.000	.000	.000	.000
	.000	.000	7720.000	.000	.000
	.000	.000	.000	.000	.000
	.000	.000	.000	.000	9024.000

	energy(NEVPT2)	energy(calc.)	difference	weight
31	14826.000	17026.612	2200.612	1.000
31	15126.000	17320.869	2194.869	1.000
31	18603.000	18005.480	-597.520	1.000
31	19760.000	19276.012	-483.988	1.000
31	20157.000	19493.578	-663.422	1.000
31	20350.000	19899.430	-450.570	1.000
31	20588.000	20435.931	-152.069	1.000
31	21092.000	20439.872	-652.128	1.000
31	21699.000	20624.046	-1074.954	1.000
31	22026.000	21729.541	-296.459	1.000
31	22147.000	21919.030	-227.970	1.000

Sigma, weighted = 1071.886

NEXT CALCULATION WITH OPTIMIZED PARAMETERS

zeta = 459.340093
*
esig = 4512.000000
epis = 3860.000000
epic = 1018.000000
B = 822.000000
C = 3965.000000

2	.000	2
4	220.022	2
5	441.667	1
6	441.872	1
7	651.568	1
8	682.310	1
9	895.606	1
10	896.336	1

	energy(NEVPT2)	energy(calc.)	difference	weight
1A	242.700	441.667	198.967	1.000
1A	624.600	441.872	-182.728	1.000
1A	704.400	651.568	-52.832	1.000
1A	710.100	682.310	-27.790	1.000
1A	870.400	895.606	25.206	1.000
1A	870.700	896.336	25.636	1.000

Sigma, weighted = 113.896

FeO2 Complex 5

NEXT CALCULATION WITH OPTIMIZED PARAMETERS

B = 1089.634044
C = 3734.199544
*
esig = 5255.000000
epis = 3397.000000
epic = 2439.000000

AOM matrix

	xy	xz	yz	x2y2	z2
.000	.000	.000	.000	.000	.000
.000	4878.000	.000	.000	.000	.000
.000	.000	6794.000	.000	.000	.000
.000	.000	.000	.000	.000	.000
.000	.000	.000	.000	.000	10510.000

	energy(NEVPT2)	energy(calc.)	difference	weight
31	16813.000	18701.652	1888.652	1.000
31	17024.000	18752.620	1728.620	1.000
31	19213.000	19697.928	484.928	1.000
31	20376.000	20160.790	-215.210	1.000
31	21559.000	20530.490	-1028.510	1.000
31	21573.000	20896.217	-676.783	1.000
31	21789.000	20898.006	-890.994	1.000
31	21977.000	21615.522	-361.478	1.000
31	22434.000	21653.316	-780.684	1.000
31	22699.000	22438.641	-260.359	1.000
31	23231.000	23146.061	-84.939	1.000

Sigma, weighted = 951.624

EXT CALCULATION WITH OPTIMIZED PARAMETERS

zeta = 419.695627
*
esig = 5255.000000
epis = 3397.000000
epic = 2439.000000
B = 1090.000000
C = 3734.000000

2	.000	2
4	207.865	2
5	414.973	1
6	414.981	1
7	619.609	1
8	623.310	1
9	827.308	1
10	827.546	1

	energy(NEVPT2)	energy(calc.)	difference	weight
2A	184.800	207.865	23.065	1.000
1A	300.900	414.973	114.073	1.000
1A	521.400	414.981	-106.419	1.000
1A	632.800	619.609	-13.191	1.000
1A	641.500	623.310	-18.190	1.000
1A	820.400	827.308	6.908	1.000
1A	820.500	827.546	7.046	1.000

Sigma, weighted = 60.323

FeN2 Complex 6, bent

NEXT CALCULATION WITH OPTIMIZED PARAMETERS

B = 601.565204
C = 4363.667501
*
esig = 5291.000000
epis = 3414.000000
epic = 884.000000

AOM matrix

	xy	xz	yz	x2y2	z2
230.757	299.910	-39.257	-25.519	44.200	
299.910	7089.782	-179.343	32.969	-52.124	
-39.257	-179.343	1833.322	343.079	-594.230	
-25.519	32.969	343.079	747.916	-270.769	
44.200	-52.124	-594.230	-270.769	9276.223	

no.	Energy	Mult.(total)	Mult.(Spin)	Mult.(Bahn)
1	.000	5	5	1
2	431.088	5	5	1
3	1655.452	5	5	1
4	6892.644	5	5	1
5	9118.458	5	5	1
6	17166.822	3	3	1
7	17580.372	3	3	1
8	18965.957	3	3	1
9	19397.533	3	3	1
10	20209.499	3	3	1
11	20250.241	3	3	1
12	20778.353	3	3	1
13	21235.648	3	3	1
14	21316.116	3	3	1
15	22051.937	3	3	1
16	22221.843	3	3	1
17	22591.456	3	3	1
18	22661.917	3	3	1
19	23923.432	3	3	1
20	23941.724	3	3	1
21	24450.551	3	3	1
22	24644.504	3	3	1
23	24761.703	3	3	1
24	24962.617	3	3	1

	energy(NEVPT2)	energy(calc.)	difference	weight
31	15593.000	17166.822	1573.822	1.000
31	16340.000	17580.372	1240.372	1.000
31	18806.000	18965.957	159.957	1.000
31	18977.000	19397.533	420.533	1.000
31	20631.000	20209.499	-421.501	1.000
31	20631.000	20250.241	-380.759	1.000
31	21282.000	20778.353	-503.647	1.000
31	21538.000	21235.648	-302.352	1.000
31	22096.000	21316.116	-779.884	1.000
31	22632.000	22051.937	-580.063	1.000
31	22768.000	22221.843	-546.157	1.000

Sigma, weighted = 746.420

NEXT CALCULATION WITH OPTIMIZED PARAMETERS

zeta = 821.954276
*
esig = 5291.000000
epis = 3414.000000
epic = 884.000000
B = 602.000000
C = 4364.000000

1	.000	1
2	.072	1
3	325.774	1
4	330.940	1
5	537.469	1
6	925.496	1
7	1067.246	1
8	1252.741	1
9	1523.451	1
10	1524.610	1

	energy (exp.)	energy(calc.)	difference	weight	
1		78.100	325.774	247.674	1.000
1		84.900	330.940	246.040	1.000
1		111.200	537.469	426.269	1.000
1		1347.500	925.496	-422.004	1.000
1		1373.200	1067.246	-305.954	1.000
1		1384.400	1252.741	-131.659	1.000
1		1464.700	1523.451	58.751	1.000
1		1464.800	1524.610	59.810	1.000

Sigma, weighted = 273.779

Complex 7 Fe[N(CCH3)]2

NEXT CALCULATION WITH OPTIMIZED PARAMETERS

B = 1833.533996
C = 3030.579870
*
esig = 3977.000000
epis = 4530.000000
epic = 2321.000000

AOM matrix

	xy	xz	yz	x2y2	z2
	.365	34.276	18.561	.093	-.297
34.276		6774.630	402.192	26.327	-10.362
18.561	402.192		6926.796	-14.393	14.879
.093	26.327	-14.393		.275	-.148
-.297	-10.362	14.879	-.148		7953.935

	energy(NEVPT2)	energy(calc.)	difference	weight
31	20176.000	20443.505	267.505	1.000
31	20309.000	20724.530	415.530	1.000
31	21542.000	21214.879	-327.121	1.000
31	21597.000	21508.945	-88.055	1.000
31	21688.000	21541.698	-146.302	1.000
31	21771.000	21652.027	-118.973	1.000
31	21858.000	21672.623	-185.377	1.000
31	22204.000	22541.143	337.143	1.000
31	22270.000	22543.898	273.898	1.000
31	22323.000	22565.895	242.895	1.000
31	23376.000	22679.234	-696.766	1.000

Sigma, weighted = 325.028

NEXT CALCULATION WITH OPTIMIZED PARAMETERS

zeta = 591.085659
*
esig = 3977.000000
epis = 4530.000000
epic = 2321.000000
B = 1834.000000
C = 3030.000000

	energy(NEVPT2)	energy(calc.)	difference	weight
1	122.800	294.768	171.968	1.000
1	172.200	585.461	413.261	1.000
1	937.600	585.473	-352.127	1.000
1	985.700	871.189	-114.511	1.000
1	986.800	873.318	-113.482	1.000
1	1106.700	1154.798	48.098	1.000
1	1106.800	1155.799	48.999	1.000

Sigma, weighted = 225.216

SI.15. Optimized geometries of FeX₂ model complexes (Fig.1) and FeCl₂¹⁺

```
NAME = fech32_pbe_opt_qro.inp  
| 1> ! PBE VDW10 TZVP TZVP/J  
| 2> !noautostart noiter moread normalprint  
| 3> %moinp "fech32_pbe_opt.qro"  
| 4>  
| 5> *xyzfile 0 5 fech32_pbe_opt.xyz
```

CARTESIAN COORDINATES (ANGSTROEM)

```
Fe -0.000000 0.000000 0.000000  
C -0.000183 0.000000 2.025151  
C 0.000183 0.000000 -2.025151  
H -1.024491 -0.000002 2.437754  
H 1.024491 -0.000002 -2.437754  
H 0.512820 -0.886256 2.440171  
H 0.512817 0.886258 2.440171  
H -0.512820 -0.886256 -2.440171  
H -0.512817 0.886258 -2.440171
```

INTERNAL COORDINATES (ANGSTROEM)

```
Fe 0 0 0 0.000000 0.000 0.000  
C 1 0 0 2.025151 0.000 0.000  
C 1 2 0 2.025151 180.000 0.000  
H 2 1 3 1.104286 111.945 0.000  
H 3 1 2 1.104286 111.945 0.000  
H 2 1 3 1.104927 112.059 0.000  
H 2 1 3 1.104927 112.059 0.000  
H 3 1 2 1.104927 112.059 0.000  
H 3 1 2 1.104927 112.059 0.000
```

```
NAME = feoh2_pbe_opt_qro.inp
| 1> # Model complex Fe(OH)2
| 2> #
| 3> !PBE TZVP TZV/J
| 4> !noautostart noiter moread normalprint
| 5> %moinp "feoh2_pbe_opt.qro"
| 6>
| 7>
| 8> *xyzfile 0 5 feoh2_pbe_opt.xyz
| 9>
| 10>
| 11>          ****END OF INPUT****
```

CARTESIAN COORDINATES (ANGSTROEM)

```
-----
Fe  -0.000000  0.000030  0.000000
N   -0.000000  0.000029  1.856424
N   -0.000000  0.000031 -1.856424
H    0.828499 -0.000265  2.448922
H   -0.828499  0.000220  2.448922
H    0.828499  0.000220 -2.448922
H   -0.828499 -0.000265 -2.448922
```

INTERNAL COORDINATES (ANGSTROEM)

```
-----
Fe  0  0  0  0.000000  0.000  0.000
N   1  0  0  1.856424  0.000  0.000
N   1  2  0  1.856424 180.000  0.000
H   2  1  3  1.018560 125.570  0.000
H   2  1  3  1.018560 125.570  0.000
H   3  1  2  1.018560 125.570  0.000
H   3  1  2  1.018560 125.570  0.000
```

INPUT FILE

```
=====
NAME = feoh2_pbe_opt_qro.inp
| 1> # Model complex Fe(OH)2
| 2> #
| 3> ! TZVP TZV/J
| 4> !noautostart noiter moread normalprint
| 5> %moinp "feoh2_pbe_opt.qro"
| 6>
| 7>
| 8>
| 9>
| 10> *xyzfile 0 5 feoh2_pbe_opt.xyz
| 11>
| 12>
| 13>
| 14>
| 15>          ****END OF INPUT****
```

CARTESIAN COORDINATES (ANGSTROEM)

```
-----  
Fe -0.000000  0.000356  0.000000  
O  0.047556  0.000357  1.792348  
O -0.047556  0.000357 -1.792348  
H  0.832091 -0.000535  2.359706  
H -0.832091 -0.000535 -2.359706
```

INTERNAL COORDINATES (ANGSTROEM)

```
-----  
Fe  0  0  0  0.000000  0.000  0.000  
O   1  0  0  1.792979  0.000  0.000  
O   1  2  0  1.792979 180.000  0.000  
H   2  1  3  0.968190 127.393 269.934  
H   3  1  2  0.968190 127.393 269.934
```

=====

INPUT FILE

=====

```
NAME = fe3cl2_bent_opt_qro.inp  
| 1> ! PBE TZVP TZV/J  
| 2> !noautostart noiter moread normalprint  
| 3> %moinp "fe3cl2_bent_opt.qro"  
| 4>  
| 5> *xyzfile 1 6 fe3cl2_bent_opt.xyz  
| 6>  
| 7>  
| 8>  
| 9>  
| 10>          ****END OF INPUT****  
=====
```

CARTESIAN COORDINATES (ANGSTROEM)

```
-----  
Fe  0.000000  0.000000 -0.425386  
Cl  1.910929  0.000000 -1.271307  
Cl -1.910930  0.000000 -1.271307
```

INTERNAL COORDINATES (ANGSTROEM)

```
-----  
Fe  0  0  0  0.000000  0.000  0.000  
Cl  1  0  0  2.089792  0.000  0.000  
Cl  1  2  0  2.089793 132.244  0.000
```