

Supplementary Information 1 (S1)

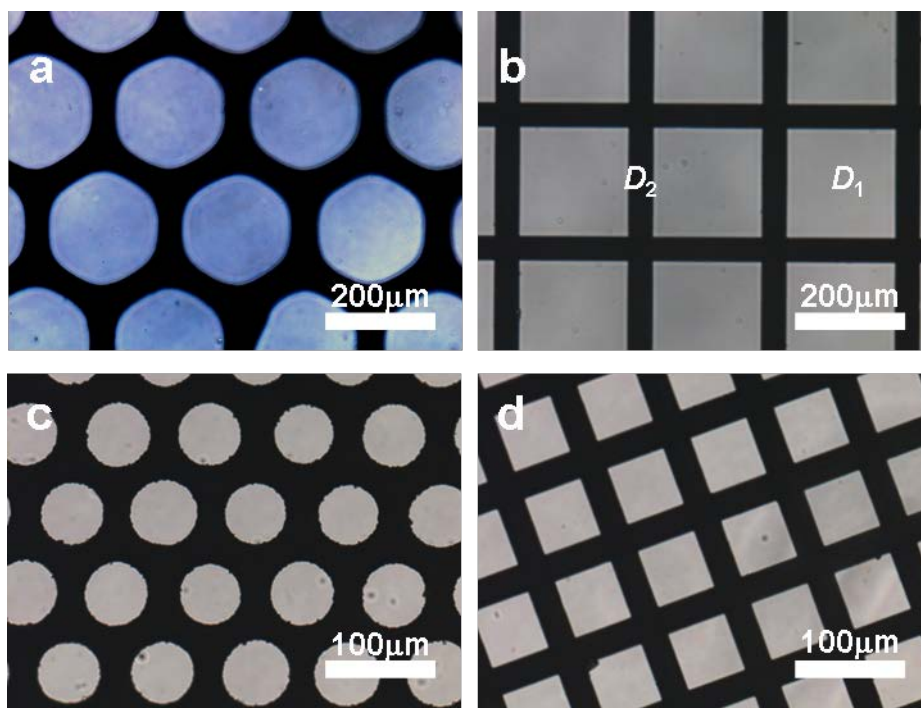


Fig. S1 Optical images of the applied copper grid: (a) H100; (b) S100; (c) H300; (d) S300.

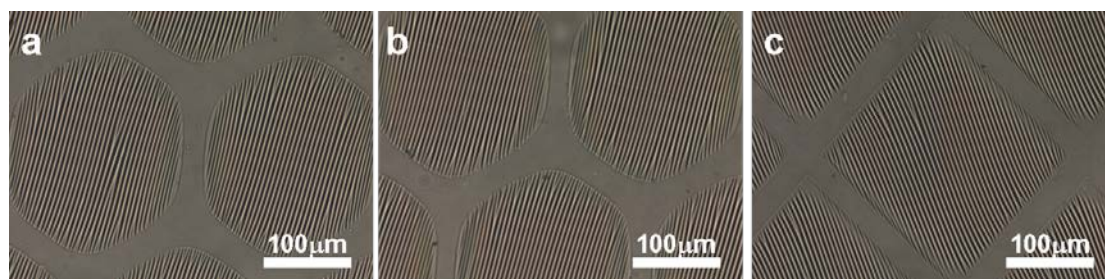


Fig. S2 Optical images of the patterned PDMS wrinkles from the first OPT for 15 min (a), followed by the second OPT for 15 min (b,c). The applied copper grid: (a,b) H100; (c) S100.

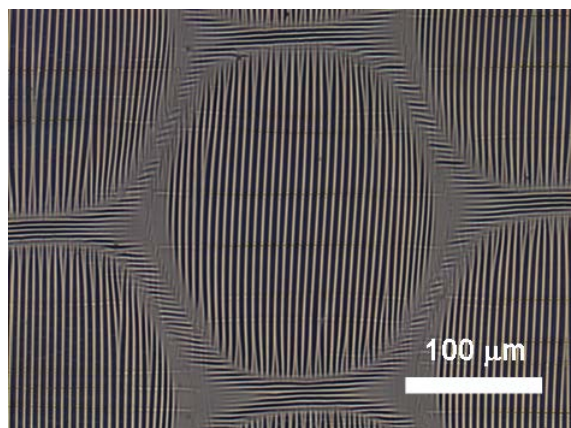


Fig. S3 Optical image of the patterned PDMS wrinkles in the case of $t_1 = 20$ min and $t_2 = 10$ min followed by the post-treatment.

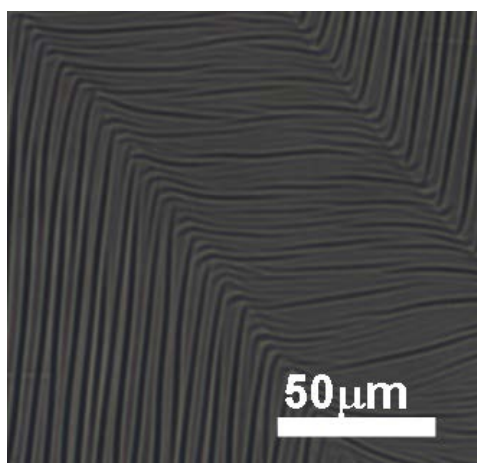


Fig. S4 Optical image of the patterned (Pt/SiO_x)/PDMS wrinkles in the case of $t_1 = t_2 = 10$ min and the ion-sputtered 4 nm-thick Pt film.

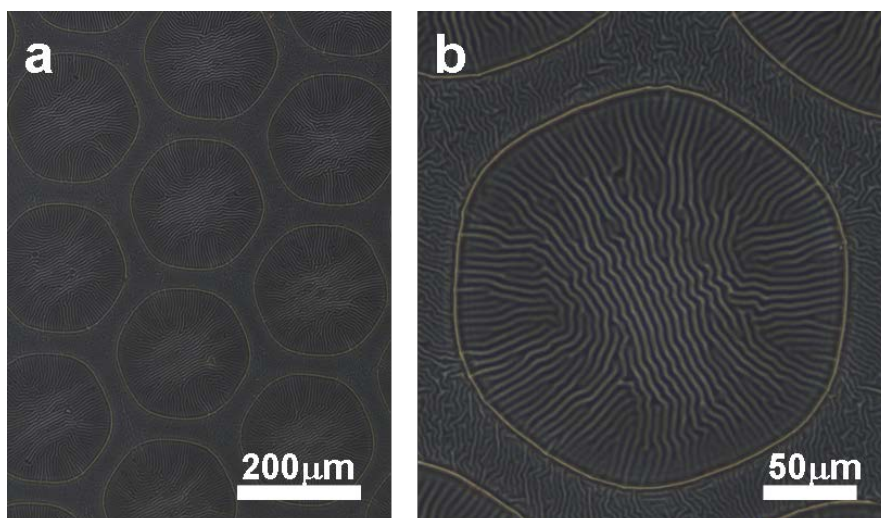


Fig. S5 Optical images of the patterned (Pt/SiO_x)/PDMS wrinkles in the case of $t_1 = t_2 = 10$ min and the ion-sputtered 8 nm-thick Pt film. It is noted that the two times' OPT were both carried out on the strain-free PDMS sheet.

Supplementary Information 2 (S2)

Our analysis below closely follows those given by Koiter for Winkler's foundation,⁴⁶ but here the foundation stiffness is different; instead of Winkler foundation we use the foundation stiffness by Biot⁴⁷ which depends on the wrinkling wavelength and is obtained by assuming the substrate as an elastic half-space.

Assuming the deflection in the following form

$$w(X) = a \sin kX \quad (\text{S-1})$$

Inserting Eq. (S-1) into (1.2) gives

$$P_2[w; N] = \frac{1}{4} [Bk^4 + c - Nk^2] a^2 \quad (\text{S-2})$$

When the axial force N is small, $P_2 > 0$ and the system is stable. At the critical buckling state, $P_2 = 0$ and in this case from Eq. (S-2), we have

$$N_1 = k^2 B + \frac{c}{k^2} = k^2 B + \frac{1}{2k} \bar{E}_s \quad (\text{S-3})$$

The minimum critical load corresponding to

$$k = \left(\frac{\bar{E}_s}{4B} \right)^{\frac{1}{3}} \quad (\text{S-4})$$

is given by

$$N_{\min} = \frac{3}{4} (4\bar{E}_s^2 B)^{\frac{1}{3}} \quad (\text{S-5})$$

From Eqs. (1.4), (S-1) and (S-5), we obtain

$$P_4[w; N_{\min}] = -\frac{5}{1024} \frac{\bar{E}_s^2}{B} a^4 \quad (\text{S-6})$$

Eq. (S-6) indicates that $P_4 < 0$ and the buckled film is very soft.

To further illustrate this point, we explore below the tangent stiffness of the system beyond the bifurcation point. When N is beyond N_{\min} , increment of the second variation P_2 reads

$$\begin{aligned}\Delta P_2[w; N_{\min}] &= \frac{k}{2\pi} \int_0^{2\pi/k} \left[-\frac{1}{2} (N - N_{\min}) a^2 k^2 \cos^2 kX \right] dX \\ &= \frac{1}{4} (1 - \lambda) a^2 N_{\min} k^2\end{aligned}\quad (\text{S-7})$$

where $\lambda = N / N_{\min}$. From Eqs. (S-4), (S-5) and (S-7) we have

$$\Delta P_2[w; N_{\min}] = \frac{3}{16} (1 - \lambda) a^2 \bar{E}_s \left(\frac{\bar{E}_s}{4B} \right)^{1/3} \quad (\text{S-8})$$

Based on Eqs. (S-6) and (S-8), we have the function F defined in the general theory of Koiter in the following form

$$F(a; \lambda) = \frac{3}{16} (1 - \lambda) a^2 \bar{E}_s \left(\frac{\bar{E}_s}{4B} \right)^{1/3} - \frac{5}{1024} \frac{\bar{E}_s^2}{B} a^4 \quad (\text{S-9})$$

The equilibrium condition reads

$$\frac{\partial F}{\partial a} = \bar{E}_s \left(\frac{3}{8} (1 - \lambda) \left(\frac{\bar{E}_s}{4B} \right)^{1/3} a - \frac{5}{256} \frac{\bar{E}_s}{B} a^3 \right) = 0 \quad (\text{S-10})$$

Eq. (S-10) gives

$$a^2 = \frac{24}{5} (1 - \lambda) \left(\frac{\bar{E}_s}{4B} \right)^{-2/3} \quad (\text{S-11})$$

The shorting after the onset of buckling is given by

$$\begin{aligned}-\Delta l &= \int_0^{2\pi/k} \left[\sqrt{1 - w'^2} - 1 \right] dX \approx \int_0^{2\pi/k} \frac{1}{2} w'^2 dX \\ &= \frac{6}{5} (1 - \lambda) \frac{2\pi}{k}\end{aligned}\quad (\text{S-12})$$

Further we obtain

$$\frac{1}{E_t h} = \frac{\partial(-\Delta l / (2\pi / k))}{\partial N} = -\frac{6}{5N_{\min}} \quad (\text{S-13})$$

Eq. (S-13) gives the tangent stiffness of the buckled film as

$$E_t = -\frac{5}{6} \frac{N_{\min}}{h} \quad (\text{S-14})$$

which is negative, indicating that the buckled film is much softer than its ground state, which is consistent with the conclusion drawn from Eq. (S-6).