

*Supplementary Materials for:*

**Temperature Memory Effect in Amorphous Shape Memory Polymers**

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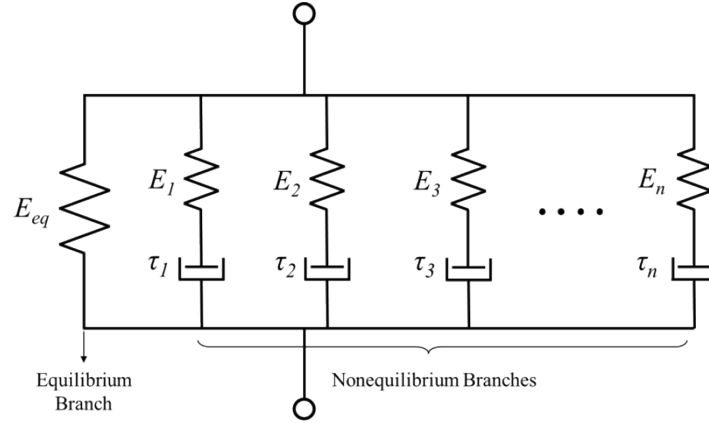
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This file includes:

1. Multi-branch Rheological Model and the Material Parameters Identification

## 1. Multi-branch Rheological Model and the Material Parameters Identification

The thermomechanical properties of a thermo-rheologically simple polymer is described by a multi-branch rheological model, as shown in Figure S1. A detailed description of the multi-branch model for SMP application can be found in Yu et al [1], and its extension for 3D nonlinear model can be found in Westbrook et al [2]. As shown in the figure, the model consists of one equilibrium branch and several thermoviscoelastic nonequilibrium branches (number of  $n$ ) to represent the multiple relaxation processes (for example, Rouse modes) of the polymer.



**Figure. S1.** 1D rheological representation of the used 1D multi-branch model

By applying the Boltzmann's superposition principle [3], the total stress  $\sigma(t)$  of the model is:

$$\sigma(t) = E_{eq} e_m(t) + \sum_{i=1}^n E_i e_i^e \quad (S1a)$$

with

$$e_i^e = \int_0^t \frac{de_m(s)}{ds} \exp\left(-\int_s^t \frac{dt'}{\tau_i(T, t')}\right) ds, \quad (S1b)$$

to describe the elastic strain evolution in each nonequilibrium branches. In Eq. S1a and S1b,  $e_m(t)$  is the total strain.  $E_{eq}$  and  $E_i$  is the elastic modulus in the equilibrium and

nonequilibrium branches, respectively.  $\tau_i(T)$  denotes the temperature dependent relaxation time in the dashpots. Here, it is assumed that the time-temperature shift for each branch follows the same rule. According to the well-established “thermorheological simplicity” principle [3] under a non-isothermal condition, the relaxation times (or viscosity) of each nonequilibrium branch vary as the temperature changes:

$$\tau_i = \alpha_T(T) \tau_i^0, \quad (\text{S2})$$

where  $\alpha_T(T)$  is a time-temperature superposition (TTSP) shifting factor, and  $\tau_i^0$  is the reference relaxation time at the temperature when  $\alpha_T(T) = 1$ .

Following O’Connell and McKenna [4], the method for calculating the temperature influence on the viscoelastic behavior depends on whether the material temperature is above or below a shifting temperature  $T_s$ . For temperatures above and near  $T_s$ , the (Williams-Landel-Ferry) WLF equation is used to calculate  $\alpha_T(T)$  [5]:

$$\log \alpha_T(T) = -\frac{C_1(T - T_r)}{C_2 + (T - T_r)}, \quad (\text{S3a})$$

where  $T_r$  is the WLF reference temperature, and  $C_1$  and  $C_2$  are material constants that depend on the choice of  $T_r$ . In comparison, for temperature below  $T_s$ ,  $\alpha_T(T)$  follows the Arrhenius-Type behavior [6]:

$$\ln \alpha_T(T) = -\frac{AF_c}{k_b} \left( \frac{1}{T} - \frac{1}{T_g} \right), \quad (\text{S3b})$$

where  $A$  is a material constant,  $F_c$  is the configurational energy and  $k_b$  is Boltzmann’s constant.  $T_s$  is the crossing point of two curves representing Eqs. S3a and S3b on a  $\alpha_T(T)$  vs  $T$  plot.

The model parameters were determined by using the  $\tan \delta$  and storage modulus of the SMP obtained in the DMA tests. A detailed model parameters identification method of the multi-branch model can be found in our previous study [1, 7]. As shown in Figure 3, the acrylate SMPs possesses a storage modulus of 3.05 MPa when the temperature is 95°C. This modulus could be considered as the equilibrium modulus of the material as the relaxation time in each nonequilibrium branch at this temperature is minimal. Under the temperature of 5°C, the SMP composite has a storage modulus of 2349.1MPa. So we have  $E_{eq}^0 + \sum_{i=1}^n E_i = 2349.1MPa$ .

By using the nonlinear regression (NLREG) method [8, 9], the parameters for TTSP, the number of nonequilibrium branches, as well as the modulus and relaxation time in each individual branch, can be determined by fitting the model prediction with the experimental DMA results shown in Figure 3 [7]. For the 1D multi-branch linear model, the temperature dependent storage modulus  $E_s(T)$ , loss modulus  $E_l(T)$  and  $\tan \delta(T)$  are respectively expressed as:

$$E_s(T) = E_{eq} + \sum_{i=1}^n \frac{E_i w^2 \tau_i^2}{1 + w^2 \tau_i^2}, \quad (S4a)$$

$$E_l(T) = \sum_{i=1}^n \frac{E_i w \tau_i}{1 + w^2 \tau_i^2}, \quad (S4b)$$

$$\tan \delta(T) = \frac{E_l(T)}{E_s(T)}, \quad (S4c)$$

where  $w$  is the testing frequency. For the convenience of operation, we assume that the relaxation times in the rubbery branch increase in a tenfold sequential. During the NLREG analysis, nonequilibrium branches are gradually added into the model to

improve the prediction, and the branch number is finally determined when the NLREG estimations (shown as dash lines in Figure 3) could capture the experimental storage and  $\tan\delta$  curves within the entire testing temperature range (5° C-95° C). By using the determined TTSP parameters, the shift factors were calculated by using Eq. S3 and plotted in Figure 3b as a function of temperature. The final set of model parameters are then listed in Table 1.

## References

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